### 九州大学学術情報リポジトリ Kyushu University Institutional Repository

## Periodic Behaviors of Quantum Cellular Automata

Inokuchi, Shuichi Faculty of Mathematics, Kyushu University

Mizoguchi, Yoshihiro Faculty of Mathematics, Kyushu University

Lee, Hyen Yeal School of Electrical and Computer Engineering, Pusan National University

Kawahara, Yasuo Department of Infomatics, Kyushu University

https://hdl.handle.net/2324/8691

出版情報: MHF Preprint Series. 2007-16, 2007-12-03. 九州大学大学院数理学研究院

バージョン: 権利関係:



### MHF Preprint Series

Kyushu University
21st Century COE Program
Development of Dynamic Mathematics with
High Functionality

### Periodic Behaviors of Quantum Cellular Automata

S. Inokuchi, Y. Mizoguchi, H. Y. Lee, Y. Kawahara

MHF 2007-16

(Received December 3, 2007)

Faculty of Mathematics Kyushu University Fukuoka, JAPAN

### Periodic Behaviors of Quantum Cellular Automata

Shuichi Inokuchi \* Yoshihiro Mizoguchi † Hyen Yeal Lee ‡
Yasuo Kawahara §

#### Abstract

In this paper we investigate quantum cellular automata whose global transitions are defined using a global transition function of classical cellular automata. And we prove the periodicity of behaviors of some quantum cellular automata.

Key words: cellular automata, quantization, reversibility, periodicity

#### 1 Introduction

Since classical (discrete) cellular automata (CA, for short) were introduced by J. von Neumann about 60 years ago, CA have been applied to various fields and much research is still reported.

Feynman proposed first the notion of cellular automata on principle of quantum mechanics in 1982[3], and Watrous [10] introduced the first formal model of quantum CA as a kind of quantum computer. He proved that there exists a partitioned quantum CA which can simulate any quantum Turing machine efficiently with constant slowdown. After introduction of the formal model some researches on properties of quantum CA were reported. Dürr and Santha [2] considered the properties between the local function of quantum CA and the unitarity of the global transition function, and proposed an algorithm to decide if a linear quantum CA is unitary. Van Dam [9] investigated quantum CA with circular bounded configurations and proved the existence of a universal quantum CA. Although almost quantum CA investigated till now have infinite cell space, Inokuchi and Mizoguchi [4] dealt with quantum CA with finite cell array. They introduced a notion of quantum CA with cyclic finite cell array and showed a sufficient condition for local transition functions to form quantum

<sup>\*</sup>Faculty of Mathematics, Kyushu University, 6-10-1 Hakozaki, Fukuoka 812-8581, Japan. e-mail:inokuchi@math.kyushu-u.ac.jp

<sup>&</sup>lt;sup>†</sup>Faculty of Mathematics, Kyushu University, 6-10-1 Hakozaki, Fukuoka 812-8581, Japan. <sup>‡</sup>School of Electrical and Computer Engineering, Pusan National University, Busan 609-735. Korea.

<sup>§</sup>Department of Infomatics, Kyushu University, 744 Motooka, Fukuoka 819-0395, Japan.

CA. And they introduced some examples of quantization method of reversible classical CA. And several types of construction method of quantum cellular automata were proposed by Schumacher and Werner [8] and they showed that any quantum cellular automaton is structurally reversible. Nevertheless much research on quantum CA have been published, research on dynamical behaviors of quantum CA have not been reported very much. Inui et al. [6] studied statistical dynamical behaviors of a quantum CA. They calculated the time averaged probability of finding a configuration for cell size 4 exactly in finite quantum CA defined by quantization of classical CA with Wolfram's rule 150, in addition they proved that the time averaged mean density of cells with state 1 is 0.5 for arbitrary cell size.

Because quantum computation operates under the unitary law, the reversibility of both quantum and classical systems has received much attention. A lot of papers concerned with the reversibility of CA have been published and particularly the research by Morita and Harao [7] is widely noticed. They introduced partitioned CA with which we can easily construct reversible CA, and proved that reversible CA are capable of universal computation. Wolfram [11] investigated the reversibility of several models of CA of the elementary CA with infinite cell array and showed that only six CA, whose transition functions are identity function, right-shift function, left-shift function and these complement functions, are reversible. In addition, Inokuchi et al. [5] investigated the reversibility of elementary CA with one dimensional finite cell array and proved that some CA including non-trivial CA are reversible and the other CA are not reversible.

A quantization of classical CA by rotation of classical cells is introduced in [4] and we get quantizable classical CA, i.e. reversible CA, from the results of [5]. In this paper we investigate quantum CA with finite cell array which can be determined by reversible CA and the quantization, and we focus on the periodic behaviors of quantum CA. Bertoni and Carpentieri [1] proved that for any unitary matrix  $\Delta$  and any  $\varepsilon > 0$  there exists  $t \in \mathbb{N}$  such that  $||\Delta^t - id|| \le \varepsilon$ . Hence any quantum CA behave almost periodically. Reversible classical CA with finite cell array behave periodically because CA with finite cell array are finite transition systems. It can be conjectured intuitively that quantum CA, which are determined by reversible CA and rotation of the product of  $\pi$  and a rational number, behave periodically, for example, the quantum CA, which is defined by a classical CA with behaviors of period 5 and rotation of  $\frac{2\pi}{3}$ , have periodic behaviors of period 15. In the following discussion we will prove for some quantum CA to behave periodically.

#### 2 Preliminaries

In this section we define quantum CA and we mention the reversibility of classical CA with one dimensional finite cell array. This definition of quantum CA was introduced by Inokuchi and Mizoguchi[4]. The global transition function of quantum CA is defined by the global transition function of classical CA and a

rotation matrix.

Let  $\mathbb{C}$  be the set of all complex numbers and I a singleton set  $\{*\}$ .

**Definition 1** Let X and Y be finite sets. A (transition) matrix of size  $|Y| \times |X|$  is a mapping

$$\alpha: X \times Y \to \mathbb{C}$$

for  $x \in X$  and  $y \in Y$ .

**Definition 2** For each element  $x \in X$  we define the matrix  $\varepsilon_x : I \times X \to \mathbb{C}$  of size  $|X| \times |I|$  by

$$\varepsilon_x(*,a) = \left\{ \begin{array}{ll} 1 & a = x \\ 0 & a \neq x \end{array} \right.$$

**Example 3** We let  $Q = \{0,1\}$ . Then  $\varepsilon_{00} : I \times Q^2 \to \mathbb{C}$  is

$$\varepsilon_{00} = \begin{pmatrix} \varepsilon_{00}(*, (0, 0)) \\ \varepsilon_{00}(*, (0, 1)) \\ \varepsilon_{00}(*, (1, 0)) \\ \varepsilon_{00}(*, (1, 1)) \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \end{pmatrix}.$$

**Lemma 4** Every matrix  $\rho: I \times X \to \mathbb{C}$  can be uniquely represented as a linear combination

$$\rho = \sum_{x \in X} k_x \varepsilon_x$$

where  $k_x$  is a complex number.

**Definition 5** A matrix  $\alpha: X \times Y \to \mathbb{C}$  is called a quantum matrix (q-matrix) if it satisfies the unitary law

$$\alpha^T \alpha = Id_{|X|}$$

where  $\alpha^T$  is the transposed matrix of alpha and  $Id_{|X|}$  is the identity matrix of size |X|.

**Lemma 6** • Every composite of q-matrices is also a q-matrix.

• A matrix  $M: X \times Y \to \mathbb{C}$  is a q-matrix if and only if  $\mu^T \rho = (M\mu)^T (M\rho)$  for all matrices  $\rho, \mu: I \times X \to \mathbb{C}$ .

Let Q denote the set  $\{0,1\}$  of binary digits and n a positive integer and  $f:Q^3\to Q$  be a local function with rule number R  $(0\le R\le 255)$ , where

$$R = \sum_{abc \in Q^3} 2^{4a+2b+c} f(abc).$$

A classical CA  $CA-R_c(n)$  with cyclic boundary condition has a global transition function  $\delta_{R,c}: Q^n \to Q^n$  defined by

$$\delta_{R,c}(x_1x_2\cdots x_n) = f(x_nx_1x_2)f(x_1x_2x_3)\cdots f(x_{n-1}x_nx_1),$$

and a classical CA  $CA - R_{a-b}(n)$  with fixed boundary condition a-b  $(a,b \in Q)$  has a global transition function  $\delta_{R,a-b}:Q^n\to Q^n$  defined by

$$\delta_{R,a-b}(x_1 x_2 \cdots x_n) = f(ax_1 x_2) f(x_1 x_2 x_3) \cdots f(x_{n-1} x_n b)$$

for  $x_1 x_2 \cdots x_n \in Q^n$ .

**Definition 7** A (rotation) matrix  $\lambda_{\theta}: Q \times Q \to \mathbb{C}$  is defined by

$$\lambda_{\theta} = \begin{pmatrix} \lambda_{\theta}(0,0) & \lambda_{\theta}(0,1) \\ \lambda_{\theta}(1,0) & \lambda_{\theta}(1,1) \end{pmatrix} = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} : Q \times Q \to \mathbb{C}$$
$$\begin{pmatrix} z_{0}' \\ z_{1}' \end{pmatrix} = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} z_{0} \\ z_{1} \end{pmatrix} \qquad (z_{0}, z_{1}, z_{0}', z_{1}' \in \mathbb{C}).$$

Trivially  $\lambda_{\theta}$  satisfies the unitary law, that is,  $\lambda_{\theta}$  is a q-matrix.

**Definition 8** A quantum  $CA\ QCA - [\delta, \theta](n)$  with n cells is a system consisting of a bijection  $\delta: Q^n \to Q^n$  and a q-matrix  $\lambda_\theta: Q \times Q \to \mathbb{C}$ .

$$QCA - [\delta, \theta](n) = (\delta : Q^n \to Q^n, \lambda_\theta : Q \times Q \to \mathbb{C})$$

For a function  $\delta: Q^n \to Q^n$  a matrix  $\Gamma: Q^n \times Q^n \to \mathbb{C}$  is defined by

$$\Gamma(y,x) = \begin{cases} 1 & (y = \delta(x)) \\ 0 & (y \neq \delta(x)). \end{cases}$$

If  $\delta$  is a bijection it is trivial that  $\Gamma$  is a q-matrix. For a quantum CA  $QCA - [\delta, \theta](n)$  its global transition q-matrix  $\Delta : Q^n \times Q^n \to \mathbb{C}$  is defined by

$$\Delta = (\lambda_{\theta} \otimes \lambda_{\theta} \otimes \cdots \otimes \lambda_{\theta}) \Gamma : Q^n \times Q^n \to \mathbb{C},$$

$$\Delta(y,x) = \lambda_{\theta}(y_1,\delta(x)_1) \otimes \lambda_{\theta}(y_2,\delta(x)_2) \otimes \cdots \otimes \lambda_{\theta}(y_n,\delta(x)_n) \qquad (x,y \in Q^n).$$

**Example 9** QCA- $[\delta_{90,0-0},\theta](2)$  is constructed by  $\delta_{90,0-0}:Q^2\to Q^2$  and  $\lambda_\theta:Q\times Q\to\mathbb{C}$  where  $\delta_{90,0-0}$  is the global transition function of CA- $90_{0-0}(2)$ . The global transition q-matrix  $\Delta$  is

$$\Delta = \left( \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \right) \otimes \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \right) \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$
$$= \begin{pmatrix} \cos^2\theta & \cos\theta\sin\theta & \cos\theta\sin\theta & \sin^2\theta \\ -\cos\theta\sin\theta & -\sin^2\theta & \cos^2\theta & \cos\theta\sin\theta \\ -\cos\theta\sin\theta & \cos^2\theta & -\sin^2\theta & \cos\theta\sin\theta \\ \sin^2\theta & -\cos\theta\sin\theta & -\cos\theta\sin\theta & \cos^2\theta \end{pmatrix}$$

Hence the q-matrix  $\Delta \varepsilon_{01}$  after one step transition from initial q-matrix  $\varepsilon_{01}$  is

$$\Delta \varepsilon_{01} = \Delta \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \end{pmatrix} = \begin{pmatrix} \cos \theta \sin \theta \\ -\sin^2 \theta \\ \cos^2 \theta \\ -\cos \theta \sin \theta \end{pmatrix}$$

The global transition function  $\delta: Q^n \to Q^n$  of a classical CA is not always a bijection. Since quantum computations have to satisfy the unitarity law, we can construct a quantum CA  $QCA - [\delta, \theta](n)$  if and only if the transition function  $\delta$  of a classical CA is reversible. The following table is an extract from the results of [5] and shows the reversibility of 1D finite classical CA CA - R(n) with fixed and cyclic boundary conditions.

Rule numbers	Fixed boundary	Cyclic boundary
204, 51	Reversible	Reversible
240,15, 170, 85	Not reversible	Reversible
90, 165	Reversible	Not reversible
	iff $n = 0 \pmod{2}$	
60,195,102,153	Reversible	Not reversible
150,105	Reversible	$\operatorname{Reversible}$
	iff $n \neq 2 \pmod{3}$	iff $n \neq 0 \pmod{3}$
166,180,154,210,	Not reversible	Reversible
89,75,101,45		iff $n = 1 \pmod{2}$
others	Not reversible	Not reversible

We introduce the following term concerned with periodic behaviors of quantum CA.

**Definition 10** A quantum CA  $QCA - [\delta, \theta](n)$  is called "global periodic" if there exists a positive integer p such that

$$\Delta^p = Id_{|Q^n|}.$$

Now we will use new notations of  $\kappa_0 = \cos \theta$ ,  $\kappa_1 = \sin \theta$  and  $\oplus$  defined as follows:

$$(x \oplus y)_i = x_i + y_i \pmod{2} \qquad (x, y \in Q^n)$$

Then using those notations, the (a, b)-th component of  $\lambda_{\theta}$  is represented by

$$\lambda_{\theta}(a,b) = (-1)^{(1-a)b} \kappa_{a \oplus b}(\theta).$$

And we let  $u = 11 \cdots 1 \in Q^n$ . Then for  $x, y \in Q^n$  we define

$$\langle x, y \rangle = \sum_{k=1}^{n} x_k y_k,$$

which is the absolute sum but not the sum modulo 2, and

$$m(x) = \langle x, u \rangle$$
.

It is trivial that  $m(x \oplus u) = n - m(x)$  and  $\langle x, y \rangle + \langle x, z \rangle = \langle x, y \oplus z \rangle \pmod{2}$ .

At the end of this section we prove the following lemma which show the elements of the global transition q-matrix  $\Delta$  of quantum CA.

**Lemma 11** For all configurations  $x, y \in Q^n$  the (y, x)-th component of the q-matrix  $\Delta : Q^n \times Q^n \to \mathbb{C}$  of  $QCA - [\delta, \theta](n)$  is given by

$$\Delta(y,x) = (-1)^{\langle \delta(x), y \oplus u \rangle} \cos^{n-m(\delta(x) \oplus y)} \theta \sin^{m(\delta(x) \oplus y)} \theta.$$

**Proof.** Let  $x, y, z \in Q^n$  and  $x' = \delta(x)$ . Then we have

$$\Delta(y, x) 
= \lambda_{\theta}(y_{1}, x'_{1})\lambda_{\theta}(y_{2}, x'_{2}) \cdots \lambda_{\theta}(y_{n}, x'_{n}) 
= (-1)^{(1-y_{1})x'_{1}} \kappa_{y_{1} \oplus x'_{1}}(\theta) (-1)^{(1-y_{2})x'_{2}} \kappa_{y_{2} \oplus x'_{2}}(\theta) \cdots (-1)^{(1-y_{n})x'_{n}} \kappa_{y_{n} \oplus x'_{n}}(\theta) 
= (-1)^{} \kappa_{y_{1} \oplus x'_{1}}(\theta) \kappa_{y_{2} \oplus x'_{2}}(\theta) \cdots \kappa_{y_{n} \oplus x'_{n}}(\theta) 
= (-1)^{} \cos^{n-m(x' \oplus y)} \theta \sin^{m(x' \oplus y)} \theta$$

### 3 Analysis of the Behaviors of Quantum Cellular Automata

All reversible CA with 2 states, 1d finite cell array were listed in the previous section. We can consider that the behaviors of a quantum CA and its symmetric quantum CA are isomorphic. Therefore in this section we deal with either a reversible CA or its symmetric CA.

We will show the results of computer simulations of quantum CA QCA- $[\delta_{60,0-0}, \frac{\pi}{3}](5)$  (figure 1) and QCA- $[\delta_{60,0-0}, \frac{\pi}{4}](5)$  (figure 2). First we calculated

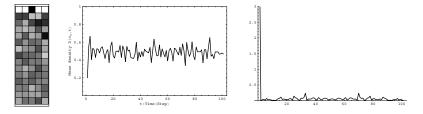


Figure 1: QCA- $[\delta_{60,0-0}, \frac{\pi}{3}](5)$ 

the probability of finding state 1 at each cells where the initial q-matrix is  $\varepsilon_{00100}$ , and the results are presented by darkness in the lefthandside figure. If quantum CA is global periodic we can see it from the lefthandside figure. So we can guess that QCA- $[\delta_{60,0-0}, \frac{\pi}{4}](5)$  is global periodic. In order to examine in detail we calculated the change of the mean density of cells with state 1 (middle figure) and the discrete Fourier transformation of the middle figure (righthandside figure). But from Figure 1 we cannot obtain the result that QCA- $[\delta_{60,0-0}, \frac{\pi}{3}](5)$  is global periodic. From our computer simulation we can guess the following theorem which show the global periodicity of the behaviors of some quantum CA.

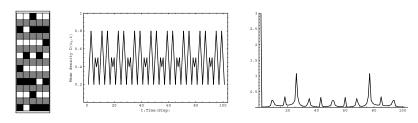


Figure 2: QCA- $[\delta_{60,0-0}, \frac{\pi}{4}](5)$ 

**Theorem 12** Let n be a positive integer. Then for QCA- $[\delta_{R,b.c.}, \theta](n)$  the following table hold where r is a rational number, g.p shows that QCA- $[\delta_{R,b.c.}, \theta](n)$  is global periodic for any n, and  $g.p._1$ ,  $g.p._2$ ,  $g.p._3$  and  $g.p._4$  show that QCA- $[\delta_{R,b.c.}, \theta](n)$  global periodic if n is even, if  $n \neq 0 \pmod{3}$ , if  $n \neq 2 \pmod{3}$  and if n is odd, respectively.

$\theta$	$\frac{\pi}{2}$		$\frac{\pi}{4}$		$r\pi$		others	
b.c.	0-0	c	0-0	c	0-0	c	0-0	c
R=204	g.p.	g.p.	g.p.	g.p.	g.p.	g.p.	?	?
R=51	g.p.	g.p.	g.p.	g.p.	g.p.	g.p.	g.p.	g.p.
R = 240,15,	?	g.p.	?	g.p.	?	g.p.	?	?
170,85								
R = 90,165	g.p1	?	g.p1	?	?	?	?	?
R = 60,195	g.p	?	g.p.	?	?	?	?	?
102,153								
R = 150, 105	g.p2	$g.p{3}$	g.p2	g.p3	?	?	?	?
R = 166, 180,	?	g.p4	?	?	?	?	?	?
154,210,								
89,75,								
101,45								

The following lemma is very useful for proving the theorem in the following discussion.

**Lemma 13** For any  $x \in Q^n$  the equation

$$\sum_{y \in Q^n} (-1)^{\langle x,y \rangle} = \left\{ \begin{array}{ll} 2^n & if \ x = 0 \\ 0 & otherwise \end{array} \right.$$

holds.

Proof.

$$\sum_{y \in Q^n} (-1)^{\langle x, y \rangle} = \sum_{y \in Q^n} (-1)^{\langle x, y \rangle}$$
$$= \sum_{y \in Q^n} (-1)^{x_1 y_1 + \dots + x_n y_n}$$

$$= \prod_{i=1}^{n} (1 + (-1)^{x_i})$$

$$= \begin{cases} 2^n & \text{if } x = 0 \\ 0 & \text{otherwise} \end{cases}$$

#### 3.1 Quantum CA with Rotation Angle $\theta = \frac{\pi}{2}$

In this subsection we discuss the behaviors of quantum CA decided by reversible CA and the  $\frac{\pi}{2}$  rotation.

The rotation q-matrix  $\lambda_{\theta}$  of  $\theta = \frac{\pi}{2}$  is presented by

$$\lambda_{\frac{\pi}{2}} = \left( \begin{array}{cc} 0 & -1 \\ 1 & 0 \end{array} \right)$$

And the negation function  $\neg: Q \to Q$  is easily represented by a unitary matrix

$$\neg = \left(\begin{array}{cc} 0 & 1 \\ 1 & 0 \end{array}\right) : Q \times Q \to \mathbb{C}$$

The negation function  $\neg: Q \to Q$  is extended into a function  $\neg_n: Q^n \to Q^n$  by

$$\neg_n = \neg \otimes \cdots \otimes \neg$$
.

Now we let  $\delta_R: Q^n \to Q^n$  be a global transition function of a reversible classical CA and  $\delta_{R'}: Q^n \to Q^n$  the compliment function of  $\delta_R$ , that is,

$$\delta_{R'}(x) = \neg_n(\delta_R(x))$$

and let  $\Delta: Q^n \times Q^n \to \mathbb{C}$  and  $\Delta': Q^n \times Q^n \to \mathbb{C}$  be global transition q-matrices of  $QCA - [\delta_R, \frac{\pi}{2}](n)$  and  $QCA - [\delta_{R'}, 0](n)$ , respectively. Note that  $\delta_{R'}$  is reversible and R' = 255 - R.

Then we have

$$\begin{array}{lcl} \Delta(z,x) & = & (\lambda_{\frac{\pi}{2}} \otimes \cdots \otimes \lambda_{\frac{\pi}{2}})(z,\delta_R(x)) \\ & = & \left\{ \begin{array}{ll} (-1)^{m(\delta(x))} & (z = \delta_R(x) \oplus u) \\ 0 & (z \neq \delta_R(x) \oplus u) \end{array} \right. \end{array}$$

and  $\delta'(x) = \delta(x) \oplus u$  for  $x, z \in Q^n$ . Hence the equation

$$|\Delta(z, x)| = |\Delta'(z, x)|$$

holds for any  $x, z \in Q^n$ . The quantum CA  $QCA - [\delta_{R'}, 0](n)$  is essentially the classical CA  $CA - R'_{b.c.}(n)$  and  $CA - R'_{b.c.}(n)$  is reversible. Therefore we have the following proposition.

**Proposition 14** For any quantum cellular automata  $QCA - [\delta_R, \frac{\pi}{2}](n)$  defined by the global transition function  $\delta_R$  of reversible classical CA and the  $\frac{\pi}{2}$  rotation q-matrix  $\lambda_{\frac{\pi}{2}}$  there exists a positive integer p such that

$$\Delta^p = Id_{|Q^n|}.$$

#### 3.2 Quantum CA with Rule 204

Trivially the global transition function  $\delta_{204}$  of CA - 204(n) is the identity function  $Id_{Q^n}$  on  $Q^n$ . We let  $\Delta = \lambda_{\theta} \otimes \lambda_{\theta} \otimes \cdots \otimes \lambda_{\theta} : Q^n \times Q^n \to \mathbb{C}$  be the global transition q-matrix of the quantum CA  $QCA - [\delta_{204}, \theta](n)$ . First we will prove the following lemma.

**Lemma 15** For any  $\theta, \theta' \in \mathbb{R}$  and any  $x \in Q^n$  the following hold:

1. 
$$(\lambda_{\theta} \otimes \cdots \otimes \lambda_{\theta})(\lambda_{\theta'} \otimes \cdots \otimes \lambda_{\theta'}) = (\lambda_{\theta+\theta'} \otimes \cdots \otimes \lambda_{\theta+\theta'})$$

2. 
$$\Delta(x,x) = \cos^n \theta$$
,

3. 
$$\Delta(x \oplus u, x) = (-1)^{n-m(x)} \sin^n \theta$$
.

The global transition q-matrix  $\Delta$  is unitary and using the above two lemmas we have

$$\Delta^m(x,x) = \cos^n m\theta.$$

Hence we can get the following proposition:

**Proposition 16** In  $QCA - [\delta_{204}, \theta](n)$  the following equation

$$\Delta^m = Id_{Q^n}$$

holds if there exists m satisfying

$$m = \left\{ \begin{array}{ll} \min \{ \ m' \in \mathbb{N} \mid m' \theta \in \pi \mathbb{Z} \ \} & if \ n \ is \ even, \\ \min \{ \ m' \in \mathbb{N} \mid m' \theta \in 2\pi \mathbb{Z} \ \} & otherwise \end{array} \right.$$

#### 3.3 Quantum CA with Rule 51

The global transition function  $\delta_{51}$  of the CA CA-51(n) is the negation function  $\neg_n$ . The global transition q-matrix of the quantum CA  $QCA-[\delta_{51},\theta](n)$  is given by

$$\Delta = (\lambda_{\theta} \neg) \otimes \cdots \otimes (\lambda_{\theta} \neg) : Q^n \times Q^n \to \mathbb{C}.$$

And we have

$$(\lambda_{\theta} \neg)^{2} = \left( \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \right)^{2}$$

$$= \begin{pmatrix} \sin \theta & \cos \theta \\ \cos \theta & -\sin \theta \end{pmatrix}^{2}$$

$$= Id_{|Q|}.$$

Therefore we can prove the following proposition:

**Proposition 17** In  $QCA - [\delta_{51}, \theta](n)$  the following equation

$$\Delta^2 = Id_{|Q^n|}$$

holds.

Proof.

$$\Delta^{2} = ((\lambda_{\theta} \neg) \otimes \cdots \otimes (\lambda_{\theta} \neg))^{2}$$

$$= (\lambda_{\theta} \neg)^{2} \otimes \cdots \otimes (\lambda_{\theta} \neg)^{2}$$

$$= Id_{|Q|} \otimes \cdots \otimes Id_{|Q|}$$

$$= Id_{|Q^{n}|}$$

#### 3.4 Quantum CA with Rule 240 and 15

Classical CA with the local functions of rule number 240 and 15 are reversible only in the case of cyclic boundary condition. We can prove the following lemmas on quantum CA  $QCA - [\delta_{240,c}, \theta](n)$  and  $QCA - [\delta_{15,c}, \theta](n)$ .

**Lemma 18** Let  $\Gamma: Q^n \times Q^n \to \mathbb{C}$  and  $\Gamma': Q^n \times Q^n \to \mathbb{C}$  be matrices determined by  $\delta_{240,c}: Q^n \to Q^n$  and  $\delta_{15,c}: Q^n \to Q^n$  respectively. For any  $\theta$  the following holds:

1. 
$$\Gamma(\lambda_{\theta} \otimes \cdots \otimes \lambda_{\theta}) = (\lambda_{\theta} \otimes \cdots \otimes \lambda_{\theta})\Gamma$$
,

2. 
$$((\lambda_{\theta} \otimes \cdots \otimes \lambda_{\theta})\Gamma')^2 = \Gamma^2$$

Using the above two lemmas we can get the configuration after 1 or 2 transition steps in  $QCA - [\delta_{240,c}, \theta](n)$  and  $QCA - [\delta_{15,c}, \theta](n)$  from  $QCA - [\delta_{204}, \theta](n)$ . Hence we get the following proposition:

**Proposition 19** Let  $\Delta$ ,  $\Delta'$  and  $\Delta''$  be the global transition q-matrices of  $QCA-[\delta_{204},\theta](n)$ ,  $QCA-[\delta_{240,c},\theta](n)$  and  $QCA-[\delta_{15,c},\theta](n)$ , respectively. If there exists a positive integer p such that  $\Delta^p=Id_{Q^n}$  then the equations

$$(\Delta')^{p'} = Id_{|Q^n|},$$

and

$$(\Delta'')^{p''} = Id_{|Q^n|}$$

hold where p' = lcm(p, n) and p'' = 2p'.

#### 3.5 Quantum CA with Rule 90

Classical CA with the local function of rule number 90 are reversible only in the case of  $n=0 \pmod 2$  and fixed boundary condition. In this subsection we discuss the behaviors of quantum CA decided by  $\delta_{90,0-0}$ ,  $\theta=\frac{\pi}{4}$  and cell size  $n=0 \pmod 2$ .

When n is even, in classical CA CA - 90(n) of 0 - 0 boundary condition it can be checked that  $\delta_{90,0-0}(x) = 0^n$  if and only if  $x = 0^n$  for  $x \in Q^n$ . Hnce we can prove the following lemma.

**Lemma 20** In  $QCA - [\delta_{90,0-0}, \frac{\pi}{4}](n)$  for any  $x \in Q^n$  the equation

$$\Delta^2(x \oplus u, x) = (-1)^{x_1 + x_n}$$

holds where n is even.

**Proposition 21** In  $QCA - [\delta_{90,0-0}, \frac{\pi}{4}](n)$  the equation

$$\Delta^4 = Id_{|Q^n|}$$

holds where n is even.

Proof.

$$\begin{array}{lcl} \Delta^4(z,x) & = & \displaystyle\sum_{y\in Q^n} \Delta^2(z,y) \Delta^2(y,x) \\ \\ & = & \left\{ \begin{array}{ll} \Delta^2(z,y) \Delta^2(y,x) & (if \ y=x\oplus u \ and \ z=y\oplus u) \\ 0 & (otherwise) \end{array} \right. \\ \\ & = & \left\{ \begin{array}{ll} 1 & (z=x) \\ 0 & (z\neq x) \end{array} \right. \end{array}$$

#### 3.6 Quantum CA with Rule 165

Classical CA with the local function of rule number 165 are reversible only in the case of  $n=0 \pmod 2$  and fixed boundary condition. In this subsection we discuss the behaviors of quantum CA decided by  $\delta_{165,0-0}$ ,  $\theta=\frac{\pi}{4}$  and cell size  $n=0 \pmod 2$ .

In classical CA CA - 165(n) of 0 - 0 boundary condition we can check easily that  $\delta_{165,0-0}(x) = 0^n$  if and only if

$$x = \begin{cases} (0110)^k & n = 4k \\ (1100)^k 11 & n = 4k + 2 \end{cases}$$

for  $x \in Q^n$  where n is even and  $k = \lfloor \frac{n}{4} \rfloor$ . Hence we can get the following lemma:

**Lemma 22** Let  $w \in Q^n$  be a configuration such that  $\delta_{165,0-0}(w) = 0^n$ . Then in  $QCA - [\delta_{165,0-0}, \frac{\pi}{4}](n)$ 

$$\Delta^2(z,x) = \begin{cases} (-1)^{x_2 + x_3 + \dots + x_{n-1}} & (z = w \oplus x \oplus u) \\ 0 & (z \neq w \oplus x \oplus u) \end{cases}$$

where n is even.

**Proposition 23** In  $QCA - [\delta_{165,0-0}, \frac{\pi}{4}](n)$  the equation

$$\Delta^4 = Id_{|Q^n|}$$

holds where n is even.

**Proof.** Let  $w \in Q^n$  be a configuration such that  $\delta_{165,0-0}(w) = 0^n$ .

$$\begin{array}{lcl} \Delta^4(z,x) & = & \displaystyle\sum_{y\in Q^n} \Delta^2(z,y)\Delta^2(y,x) \\ \\ & = & \left\{ \begin{array}{ll} \Delta^2(z,y)\Delta^2(y,x) & (if \ y=w\oplus x\oplus u \ and \ z=w\oplus y\oplus u) \\ 0 & (otherwise) \end{array} \right. \\ \\ & = & \left\{ \begin{array}{ll} 1 & (z=x) \\ 0 & (z\neq x) \end{array} \right. \end{array}$$

#### 3.7 Quantum CA with Rule 60

Classical CA with the local function of rule number 60 are reversible in the case of any cell size n and fixed boundary condition. In this subsection we discuss the behaviors of quantum CA decided by  $\delta_{60,0-0}$  and  $\theta = \frac{\pi}{4}$ .

For  $x \in Q^n$  and an integer  $i \ (0 \le i \le n)$  we define  $x(i) \in Q^n$  as follows:

- x(0) = x
- $x(i)_i = x_{n-i+1} + 1 \pmod{2}$ ,
- $x(i)_{i+j} = x_j + x_{n-i+1} + 1 \pmod{2}$  for  $1 \le j \le n-i$ ,
- $x(i)_{i-j} = x_{n-i+1} + x_{n-j+1} \pmod{2}$  for  $1 \le j \le i-1$ .

We can easily check the followings:

- $x(1) = (0, x_1, x_2, \dots, x_{n-1}) \oplus (x_n, x_n, \dots, x_n) \oplus u$
- x(i+1) = x(i)(1)
- X(n)(1) = x

In addition, we can easily check that for any  $x \in Q^n$  the equation  $\delta_{60,0-0}(x) \oplus \delta_{102,0-0}(z \oplus u) = 0^n$  holds if  $z = (0, x_1, x_2, \dots, x_{n-1}) \oplus (x_n, x_n, \dots, x_n) \oplus u$ . Then the following lemma can be proved.

**Lemma 24** In  $QCA - [\delta_{60,0-0}, \frac{\pi}{4}](n)$  the equation

$$\Delta^{2i}(z,x) = \begin{cases} (-1)^{\sum_{j=1}^{i} x(j-1)_n} & (z=x(i)) \\ 0 & (z \neq x(i)) \end{cases}$$

holds for any configuration  $x \in Q^n$  and  $1 \le i \le n$ .

For any  $x \in Q^n$  we have

$$\sum_{j=0}^{n} x(j)_n = 2\sum_{j=1}^{n} x_j + n.$$

Therefore the following proposition holds:

**Proposition 25** In  $QCA - [\delta_{60,0-0}, \frac{\pi}{4}](n)$  the following

$$\Delta^{2n+2} = (-1)^n Id_{|Q^n|}$$

holds.

Proof.

$$\begin{array}{lll} \Delta^{2n+2}(z,x) & = & \displaystyle \sum_{y \in Q^n} \Delta^2(z,y) \Delta^{2n}(y,x) \\ \\ & = & \left\{ \begin{array}{lll} (-1)^{\sum_{j=1}^n x(j-1)_n} (-1)^{y_n} & if \ y = x(n) \ and \ z = y(1) \\ 0 & otherwise \end{array} \right. \\ \\ & = & \left\{ \begin{array}{lll} (-1)^{\sum_{j=0}^n x(j)_n} & (z = x) \\ 0 & (z \neq x) \end{array} \right. \\ \\ & = & \left\{ \begin{array}{lll} (-1)^n & (z = x) \\ 0 & (z \neq x) \end{array} \right. \end{array}$$

#### 3.8 Quantum CA with Rule 195

Classical CA with the local function of rule number 195 are reversible in the case of any cell size n and fixed boundary condition. In this subsection we discuss the behaviors of quantum CA decided by  $\delta_{195,0-0}$  and  $\theta = \frac{\pi}{4}$ .

For  $x \in Q^n$  and an integer i  $(0 \le i \le n)$  we define  $x(i) \in Q^n$  as follows:

- x(0) = x
- $x(i)_i = x_{n-i+1} + ni + 1 \pmod{2}$ ,
- $x(i)_{i+j} = x_j + x_{n-i+1} + i(n+j) + 1 \pmod{2}$ for  $1 \le j \le n-i$ ,
- $x(i)_{i-j} = x_{n-i+1} + x_{n-j+1} + i(n+j) + j(n+1) \pmod{2}$  for  $1 \le j \le i-1$ .

We can easily check the following equations:

- $x(1) = (x_n, x_n, \dots, x_n) \oplus (0, x_1, x_2, \dots, x_{n-1}) \oplus (n-1, n-2, \dots, 0)$
- x(i+1) = x(i)(1)
- x(n)(1) = x

And we can easily check that for  $x \in Q^n$  the equation

$$\delta_{195,0-0}(x) \oplus \delta_{170,0-0}(z \oplus u) \oplus z \oplus u = 0^n$$

holds if

$$z = (x_n, x_n, \dots, x_n) \oplus (0, x_1, x_2, \dots, x_{n-1}) \oplus (n-1, n-2, \dots, 0).$$

Then the following lemma holds:

**Lemma 26** In  $QCA = [\delta_{195,0-0}, \frac{\pi}{4}](n)$  the equation

$$\Delta^{2i}(z,x) = \begin{cases} (-1)^{\sum_{j=1}^{i} \{ \langle x(j-1), u \rangle + nx(j-1)_n + \frac{n(n+1)}{2} \}} & (z=x(i)) \\ 0 & (z \neq x(i)) \end{cases}$$

holds for any configuration  $x \in Q^n$  and  $1 \le i \le n$ .

**Proposition 27** In  $QCA - [\delta_{195,0-0}, \frac{\pi}{4}](n)$  the following

$$\Delta^{2n+2} = \left\{ \begin{array}{ll} -Id_{|Q^n|} & n=4k+3 \\ Id_{|Q^n|} & otherwise \end{array} \right.$$

holds.

**Proof.** We have

$$\begin{split} & \Delta^{2n+2}(z,x) \\ &= \sum_{y \in Q^n} \Delta^2(z,y) \Delta^{2n}(y,x) \\ &= \begin{cases} & (-1)^{\sum_{j=1}^n \{ < x(j-1), u > + nx(j-1)_n + \frac{n(n+1)}{2} \}} (-1)^{\{ < y, u > + ny_n + \frac{n(n+1)}{2} \}} \\ & & if \quad y = x(n) \quad and \quad z = y(1) \end{cases} \\ &= \begin{cases} & (-1)^{\sum_{j=1}^{n+1} \{ < x(j-1), u > + nx(j-1)_n + \frac{n(n+1)}{2} \}} & (z = x) \\ & 0 & (z \neq x) \end{cases} \\ &= \begin{cases} & (-1)^{\left\{ \frac{n^3 + 2n^2 + n}{2} + \frac{4n^4 + 3n^3 + 5n^2}{6} + \frac{7n^4 + 8n^3 + 29n^2 + 4n}{12} \right\}} & (z = x) \\ & 0 & (z \neq x) \end{cases} \end{split}$$

and we can check that

- $\frac{n^3+2n^2+n}{2}+\frac{4n^4+3n^3+5n^2}{6}+\frac{7n^4+8n^3+29n^2+4n}{12}$  is odd if n=4k+3
- $\frac{n^3 + 2n^2 + n}{2} + \frac{4n^4 + 3n^3 + 5n^2}{6} + \frac{7n^4 + 8n^3 + 29n^2 + 4n}{12}$  is even otherwise.

Therefore we have

$$\Delta^{2n+2} = \left\{ \begin{array}{ll} -Id_{|Q^n|} & n=4k+3 \\ Id_{|Q^n|} & otherwise. \end{array} \right.$$

3.9 Quantum CA with Rule 150

Classical CA with the local function of rule number 150 are reversible in the following cases:

•  $n \neq 2 \pmod{3}$  and fixed boundary condition

•  $n \neq 0 \pmod{3}$  and cyclic boundary condition

where n is cell size. In this subsection we discuss the behaviors of each quantum CA decided as follows:

- $\delta_{150,c}$ , any  $\theta$  and cell size n=4
- $\delta_{150,c}$ ,  $\theta = \frac{\pi}{4}$  and cell size  $n \neq 0 \pmod{3}$
- $\delta_{150,0-0}$ ,  $\theta = \frac{\pi}{4}$  and cell size  $n \neq 2 \pmod{3}$

**3.9.1** 
$$QCA - [\delta_{150,c}, \theta](4)$$

We set

$$g(n) = \frac{1}{2} \left( \sum_{i=1}^{n} \cos(4i - 2)\theta \right) \sin 2\theta$$

$$A1(n) = \begin{pmatrix} \cos^{2}(2n\theta) & -g(n) & -g(n) & 0 \\ g(n) & \cos^{2}(2n\theta) & 0 & -g(n) \\ g(n) & 0 & \cos^{2}(2n\theta) & -g(n) \\ 0 & g(n) & g(n) & \cos^{2}(2n\theta) \end{pmatrix},$$

$$A2(n) = \begin{pmatrix} g(n) & 0 & 0 & g(n) \\ 0 & g(n) & -g(n) & 0 \\ 0 & -g(n) & g(n) & 0 \\ g(n) & 0 & 0 & g(n) \end{pmatrix},$$

$$A3(n) = \begin{pmatrix} 0 & -g(n) & -g(n) & \sin^{2}(2n\theta) \\ g(n) & 0 & -\sin^{2}(2n\theta) & -g(n) \\ g(n) & -\sin^{2}(2n\theta) & 0 & -g(n) \\ \sin^{2}(2n\theta) & g(n) & g(n) & 0 \end{pmatrix},$$

$$B1(n) = \begin{pmatrix} A1(n) & -A2(n) \\ A2(n) & A1(n) \end{pmatrix},$$

$$B2(n) = \begin{pmatrix} A2(n) & -A3(n) \\ A3(n) & A2(n) \end{pmatrix}, and$$

$$C(n) = \begin{pmatrix} B1(n) & -B2(n) \\ B2(n) & B1(n) \end{pmatrix}.$$

Then we can show the following lemma:

**Lemma 28** In  $QCA - [\delta_{150,c}, \theta](4)$  the following equation

$$\Delta^{2k} = C(k)$$

holds for any positive integer k.

Therefore considering positive integer m such that  $\cos^2(2m\theta) = 1$  we can get the following proposition

**Proposition 29** In  $QCA - [\delta_{150,c}, \theta](4)$  the following equation

$$\Delta^{2m} = Id_{|Q^n|}$$

holds if there exists m satisfying  $m = \min\{ m' \in \mathbb{N} \mid 2m'\theta \in \pi\mathbb{Z} \}.$ 

**3.9.2** 
$$QCA - [\delta_{150,c}, \frac{\pi}{4}](n)$$

In the classical CA CA-150 of cyclic boundary condition we can easily check that  $\delta_{150,c}(x)=0^n$  if and only if  $x=0^n$  for  $x\in Q^n$  where  $n\neq 0\pmod 3$ . Using the fact and lemma 13 we can prove the following lemma:

**Lemma 30** In  $QCA - [\delta_{150,c}, \frac{\pi}{4}](n)$  the equation

$$\Delta^2(z,x) = \left\{ \begin{array}{cc} (-1)^{< x,u>} & (z=x \oplus u) \\ 0 & (z \neq x \oplus u) \end{array} \right.$$

holds for all  $x, z \in Q^n$  where  $n \neq 0 \pmod{3}$ .

The following proposition can be got immediately from the above lemma.

**Proposition 31** In  $QCA - [\delta_{150,c}, \frac{\pi}{4}](n)$  the equation

$$\Delta^4 = (-1)^n Id_{|Q^n|}$$

holds where  $n \neq 0 \pmod{3}$ .

**Proof.** For any  $x, z \in Q^n$  we have

$$\begin{array}{lcl} \Delta^4(z,x) & = & \displaystyle\sum_{y\in Q^n} \Delta^2(z,y) \Delta^2(y,x) \\ \\ & = & \left\{ \begin{array}{ll} \Delta^2(x,x\oplus u) \Delta^2(x\oplus u,x) & (x=z) \\ 0 & (x\neq z) \end{array} \right. \\ \\ & = & \left\{ \begin{array}{ll} (-1)^n & (x=z) \\ 0 & (x\neq z) \end{array} \right. \end{array}$$

**3.9.3**  $QCA - [\delta_{150,0-0}, \frac{\pi}{4}](n)$ 

In the classical CA CA - 150 of 0 - 0 boundary condition we can easily check that  $\delta_{150,0-0}(x) = 0^n$  if and only if  $x = 0^n$  for  $x \in Q^n$  where  $n \neq 2 \pmod{3}$ . Then we can prove the following lemma:

**Lemma 32** In  $QCA = [\delta_{150,0-0}, \frac{\pi}{4}](n)$  the equation

$$\Delta^2(z,x) = \left\{ \begin{array}{cc} (-1)^{x_2+x_3+\cdots+x_{n-1}} & (z=x\oplus u) \\ 0 & (z\neq x\oplus u) \end{array} \right.$$

holds for all  $x, z \in Q^n$  where  $n \neq 2 \pmod{3}$ .

Using this lemma we can get the following proposition:

**Proposition 33** In  $QCA - [\delta_{150,0-0}, \frac{\pi}{4}](n)$ 

$$\Delta^4 = (-1)^n Id_{|Q^n|}$$

holds where  $n \neq 2 \pmod{3}$ .

**Proof.** For any  $x, z \in Q^n$  we have

$$\begin{array}{lcl} \Delta^4(z,x) & = & \displaystyle \sum_{y \in Q^n} \Delta^2(z,y) \Delta^2(y,x) \\ \\ & = & \left\{ \begin{array}{ll} \Delta^2(x,x \oplus u) \Delta^2(x \oplus u,x) & (x=z) \\ 0 & (x \neq z) \end{array} \right. \\ \\ & = & \left\{ \begin{array}{ll} (-1)^{x_2+x_3+\cdots+x_{n-1}} (-1)^{(x_2+1)+(x_3+1)+\cdots+(x_{n-1}+1)} & (x=z) \\ 0 & (x \neq z) \end{array} \right. \\ \\ & = & \left\{ \begin{array}{ll} (-1)^{n-2} & (x=z) \\ 0 & (x \neq z) \end{array} \right. \end{array}$$

#### 3.10 Quantum CA with Rule 105

Classical CA with the local function of rule number 105 are reversible in the following cases:

- $n \neq 2 \pmod{3}$  and fixed boundary condition
- $n \neq 0 \pmod{3}$  and cyclic boundary condition

where n is cell size. In this subsection we discuss the behaviors of each quantum CA decided by as follows:

- $\delta_{105,c}$ ,  $\theta = \frac{\pi}{4}$  and cell size  $n \neq 0 \pmod{3}$
- $\delta_{105,0-0}$ ,  $\theta = \frac{\pi}{4}$  and cell size  $n \neq 2 \pmod{3}$

### **3.10.1** $QCA - [\delta_{105,c}, \frac{\pi}{4}](n)$

In the classical CA CA - 105 of cyclic boundary condition we can easily check that  $\delta_{105,c}(x) = 0^n$  if and only if  $x = 1^n$  for  $x \in Q^n$  where  $n \neq 0 \pmod{3}$ . Then we can prove the following lemma:

**Proposition 34** In  $QCA - [\delta_{105,c}, \frac{\pi}{4}](n)$  the equation

$$\Delta^2 = Id_{|Q^n|}$$

holds where  $n \neq 0 \pmod{3}$ .

Proof.

$$\begin{array}{lcl} \Delta^2(z,x) & = & \displaystyle \sum_{y \in Q^n} \Delta(z,y) \Delta(y,x) \\ \\ & = & \displaystyle \sum_{y \in Q^n} \left\{ \frac{(-1)^{<\delta(y),z \oplus u>}}{(\sqrt{2})^n} \times \frac{(-1)^{<\delta(x),y \oplus u>}}{(\sqrt{2})^n} \right\} \\ \\ & = & \displaystyle \frac{(-1)^{n+3m(x)+m(z \oplus u)}}{2^n} \displaystyle \sum_{y \in Q^n} (-1)^{<\delta(x \oplus z \oplus u),y>} \\ \\ & = & \displaystyle \left\{ \begin{array}{ll} 1 & (z=x) \\ 0 & (z \neq x) \end{array} \right. \end{array}$$

**3.10.2**  $QCA - [\delta_{105,0-0}, \frac{\pi}{4}](n)$ 

In the classical CA CA-105 of 0-0 boundary condition we can easily check that  $\delta_{105,0-0}(x)=0^n$  if and only if

$$x = \begin{cases} (010)^k & if \quad n = 3k \\ (100)^k 1 & if \quad n = 3k + 1. \end{cases}$$

for  $x \in \mathbb{Q}^n$  where  $n \neq 2 \pmod{3}$ . Then we can prove the following lemma:

**Lemma 35** Let  $w \in Q^n$  be a configuration such that  $\delta_{105,0-0}(w) = 0$ . Then in  $QCA - [\delta_{105,0-0}, \frac{\pi}{4}](n)$  the equation

$$\Delta^{2}(z,x) = \begin{cases} (-1)^{x_{1}+x_{n}} & (z = w \oplus x \oplus u) \\ 0 & (z \neq w \oplus x \oplus u) \end{cases}$$

holds for any  $x, z \in Q^n$  where  $n \neq 2 \pmod{3}$ .

**Proposition 36** In  $QCA - [\delta_{105,0-0}, \frac{\pi}{4}](n)$  the equation

$$\Delta^4 = Id_{|Q^n|}$$

holds where  $n \neq 2 \pmod{3}$ .

Proof.

$$\begin{array}{lcl} \Delta^4(z,x) & = & \displaystyle\sum_{y\in Q^n} \Delta^2(z,y)\Delta^2(y,x) \\ \\ & = & \left\{ \begin{array}{ll} \Delta^2(z,y)\Delta^2(y,x) & (if\ y=w\oplus x\oplus u\ and\ z=w\oplus y\oplus u) \\ 0 & (otherwise) \end{array} \right. \\ \\ & = & \left\{ \begin{array}{ll} 1 & (z=x) \\ 0 & (z\neq x) \end{array} \right. \end{array}$$

#### 4 Conclusion

In this paper we treated quantum CA whose global transition function is defined by the global transition function of classical CA and rotation of cells. A quantum cell can be represented by any point on the sphere and a transition in quantum CA is considered as a movement from a point to another point on the sphere. While trivially classical CA with finite cell array behave in finite space, we think that generally quantum CA (including quantum CA quantizated classical finite CA) behave infinitely for any initial configuration. In this paper we proved that some quantum CA is global periodic. Global periodic quantum CA have infinite variety of configurations and finite behaviors from any initial configuration. That is, in global periodic quantum CA for any initial configuration we can get the same configuration after some step transitions. By this property we may apply global periodic quantum CA to cryptography theory or quantum communication theory by using rule number and rotation angle as keys.

#### References

- [1] Bertoni, A., and Carpentieri, M. (2001). Regular Languages Accepted by Quantum Automata. *Inf. Comput.*, 165, 174–182.
- [2] Dürr, C., and Santha, M. (2002). A Decision Procedure for unitary linear quantum cellular automata. SIAM J. Comput., 31(4), 1076–1089.
- [3] Feynman, R. (1982). Simulationg Physics with Computers. *International Journal of Theoretical Physics* Vol. 21, No. 6&7, pp. 467–488.
- [4] Inokuchi, S., and Mizoguchi, Y. (2005). Generalized Partitioned Quantum Cellular Automata and Quantization of Classical CA. *Int. Journ. of Unconventional Computing*, 1, 149–160.
- [5] Inokuchi, S., Honda, K., Lee, H., Sato, T., Mizoguchi, Y., and Kawahara, Y. (2005). On reversible cellular automata with finite cell array. *Proceedings* of Fourth International Conference on Unconventional Computation, LNCS 3699.
- [6] Inui, N., Inokuchi, S., Mizoguchi, Y., and Konno, N. (2005). Statistical properties of a quantum cellular automaton. *Physical Review A*, 72, 032323.
- [7] Morita, K., and Harao, M. (1989). Computation universality of onedimensional reversible (injective) cellular automata. *Transactions of the IE-ICE*, E 72, 758-762.
- [8] Schumacher, B. and Werner, R. (2004). Reversible Quantum Cellular Automata. arXiv:quant-ph/0405174
- [9] van Dam, W. (1996). A Universal Quantum Cellular Automaton. Proceedings of PhysComp96, 323–331.

- [10] Watrous, J. (1995). On One-Dimensional Quantum Cellularn Automata. Proceedings of the 36th Annual Symposium on Foundations of Computer Science, 528–537.
- [11] Wolfram, S. (2002). A New Kind of Science, Wolfram Media, Inc.

### Appendix A Proofs of Lemmas

**Proof.** (of lemma 15)

1. For any  $\theta, \theta' \in \mathbb{R}$  we can easily check  $\lambda_{\theta} \lambda_{\theta'} = \lambda_{\theta + \theta'}$ . Hence we have

$$(\lambda_{\theta} \otimes \cdots \otimes \lambda_{\theta})(\lambda_{\theta'} \otimes \cdots \otimes \lambda_{\theta'})$$

$$= ((\lambda_{\theta} \lambda_{\theta'}) \otimes \cdots \otimes (\lambda_{\theta} \lambda_{\theta'}))$$

$$= (\lambda_{\theta+\theta'} \otimes \cdots \otimes \lambda_{\theta+\theta'}).$$

2.

$$\Delta(x,x) = (\lambda_{\theta} \otimes \cdots \otimes \lambda_{\theta})(x,x)$$

$$= \lambda_{\theta}(x_1, x_1) \times \cdots \times \lambda_{\theta}(x_n, x_n)$$

$$= \cos^n \theta.$$

3.

$$\Delta(x \oplus u, x) = (\lambda_{\theta} \otimes \cdots \otimes \lambda_{\theta})(x \oplus u, x) 
= \lambda_{\theta} (1 - x_{1}, x_{1}) \times \cdots \times \lambda_{\theta} (1 - x_{n}, x_{n}) 
= \sin^{m(x)} \theta(-\sin \theta)^{n - m(x)} 
= (-1)^{n - m(x)} \sin^{n} \theta.$$

**Proof.** (of lemma 18)

1. For any  $x, y \in Q^n$  we have

$$((\lambda_{\theta} \otimes \cdots \otimes \lambda_{\theta})\Gamma)(y, x) = (\lambda_{\theta} \otimes \cdots \otimes \lambda_{\theta})(y, x_{n}x_{1}x_{2} \cdots x_{n-1})$$

$$= \lambda_{\theta}(y_{1}, x_{n}) \times \lambda_{\theta}(y_{2}, x_{1}) \times \lambda_{\theta}(y_{3}, x_{2}) \times \cdots \times \lambda_{\theta}(y_{n}, x_{n-1})$$

$$= \lambda_{\theta}(y_{2}, x_{1}) \times \lambda_{\theta}(y_{3}, x_{2}) \times \cdots \times \lambda_{\theta}(y_{n}, x_{n-1}) \times \lambda_{\theta}(y_{1}, x_{n})$$

$$= ((\lambda_{\theta} \otimes \cdots \otimes \lambda_{\theta}))(y_{2}y_{3} \cdots y_{n}y_{1}, x)$$

$$= (\Gamma(\lambda_{\theta} \otimes \cdots \otimes \lambda_{\theta}))(y, x).$$

2. The equations  $\neg_n\Gamma=\Gamma\neg_n$  and  $\Gamma'=\neg_n\Gamma$  can be checked easily. Hence we have

$$((\lambda_{\theta} \otimes \cdots \otimes \lambda_{\theta})\Gamma')^{2}$$

$$= ((\lambda_{\theta} \otimes \cdots \otimes \lambda_{\theta}) \neg_{n} \Gamma)^{2}$$

$$= (\lambda_{\theta} \otimes \cdots \otimes \lambda_{\theta}) \neg_{n} (\lambda_{\theta} \otimes \cdots \otimes \lambda_{\theta}) \Gamma \neg_{n} \Gamma$$

$$= ((\lambda_{\theta} \otimes \cdots \otimes \lambda_{\theta}) \neg_{n})^{2} \Gamma^{2}$$

$$= \Gamma^{2}$$

**Proof.** (of lemma 20)

$$\Delta^{2}(z,x) = \sum_{y \in Q^{n}} \Delta(z,y)\Delta(y,x)$$

$$= \sum_{y \in Q^{n}} \left\{ \frac{(-1)^{<\delta_{90,0-0}(y),z \oplus u>}}{(\sqrt{2})^{n}} \times \frac{(-1)^{<\delta_{90,0-0}(x),y \oplus u>}}{(\sqrt{2})^{n}} \right\}$$

$$= \frac{(-1)^{<\delta_{90,0-0}(x),u>}}{2^{n}} \sum_{y \in Q^{n}} (-1)^{<\delta_{90,0-0}(x \oplus z \oplus u),y>}$$

$$= \begin{cases} (-1)^{<\delta_{90,0-0}(x),u>} & (x \oplus z \oplus u = 0^{n}) \\ 0 & (x \oplus z \oplus u \neq 0^{n}) \end{cases}$$

$$= \begin{cases} (-1)^{x_{1}+x_{n}} & (z = x \oplus u) \\ 0 & (z \neq x \oplus u) \end{cases}$$

**Proof.** (of lemma 22)

$$\begin{split} \Delta^2(z,x) &= \sum_{y \in Q^n} \Delta(z,y) \Delta(y,x) \\ &= \sum_{y \in Q^n} \left\{ \frac{(-1)^{<\delta_{165,0-0}(y),z \oplus u>}}{(\sqrt{2})^n} \times \frac{(-1)^{<\delta_{165,0-0}(x),y \oplus u>}}{(\sqrt{2})^n} \right\} \\ &= \frac{(-1)^{<\delta_{165,0-0}(x) \oplus z \oplus u,u>}}{2^n} \sum_{y \in Q^n} (-1)^{<\delta_{165,0-0}(x \oplus z \oplus u),y>} \\ &= \left\{ \begin{array}{c} (-1)^{<\delta_{165,0-0}(x) \oplus z \oplus u,u>} & (x \oplus z \oplus u = w) \\ 0 & (x \oplus z \oplus u \neq w) \\ 0 & (z \neq w \oplus x \oplus u) \end{array} \right. \end{split}$$

**Proof.** (of lemma 24)

We prove it by induction on i. First, we have

$$\begin{split} \Delta^2(z,x) &= \sum_{y \in Q^n} \Delta(z,y) \Delta(y,x) \\ &= \sum_{y \in Q^n} \left\{ \frac{(-1)^{<\delta_{60,0-0}(y),z \oplus u>}}{(\sqrt{2})^n} \times \frac{(-1)^{<\delta_{60,0-0}(x),y \oplus u>}}{(\sqrt{2})^n} \right\} \\ &= \frac{(-1)^{<\delta_{60,0-0}(x),u>}}{2^n} \sum_{y \in Q^n} (-1)^{<\delta_{60,0-0}(x) \oplus \delta_{102,0-0}(z \oplus u),y>} \end{split}$$

$$= \begin{cases} (-1)^{<\delta_{60,0-0}(x),u>} & (z=x(1)) \\ 0 & (z\neq x(1)) \end{cases}$$
$$= \begin{cases} (-1)^{x(0)_n} & (z=x(1)) \\ 0 & (z\neq x(1)) \end{cases}.$$

And we assume that the equation holds in the case of  $i \leq k$ . Then we have

$$\begin{array}{lcl} \Delta^{2k+2}(z,x) & = & \displaystyle \sum_{y \in Q^n} \Delta^2(z,y) \Delta^{2k}(y,x) \\ \\ & = & \left\{ \begin{array}{ll} (-1)^{\sum_{j=1}^k x(j-1)_n} (-1)^{y_n} & if \ y = x(k) \ and \ z = y(1) \\ 0 & otherwise \end{array} \right. \\ \\ & = & \left\{ \begin{array}{ll} (-1)^{\sum_{j=1}^{k+1} x(j-1)_n} & (z = x(k+1)) \\ 0 & (z \neq x(k+1)) \end{array} \right. \end{array}$$

That is, in the case of i = k + 1 the equation holds.

#### **Proof.** (of lemma 26)

We prove it by induction on i. First, we have

$$\begin{split} \Delta^2(z,x) &= \sum_{y \in Q^n} \Delta(z,y) \Delta(y,x) \\ &= \sum_{y \in Q^n} \left\{ \frac{(-1)^{<\delta_{195,0-0}(y),z \oplus u>}}{(\sqrt{2})^n} \times \frac{(-1)^{<\delta_{195,0-0}(x),y \oplus u>}}{(\sqrt{2})^n} \right\} \\ &= \frac{(-1)^{<\delta_{195,0-0}(x) \oplus z \oplus u,u>}}{2^n} \sum_{y \in Q^n} (-1)^{<\delta_{195,0-0}(x) \oplus \delta_{170,0-0}(z \oplus u) \oplus z \oplus u,y>} \\ &= \left\{ \begin{array}{c} (-1)^{<\delta_{195,0-0}(x) \oplus z \oplus u,u>} & (z = x(1)) \\ 0 & (z \neq x(1)) \end{array} \right. \\ &= \left. \begin{cases} (-1)^{ + nx_n + \frac{n(n+1)}{2}} & (z = x(1)) \\ 0 & (z \neq x(1)) \end{array} \right. \end{split}$$

And we assume that the equation holds in the case of  $i \leq k$ . Then we have

$$\begin{split} & \Delta^{2k+2}(z,x) \\ &= \sum_{y \in Q^n} \Delta^2(z,y) \Delta^{2k}(y,x) \\ &= \begin{cases} & (-1)^{\sum_{j=1}^k \{ < x(j-1),u > + nx(j-1)_n + \frac{n(n+1)}{2} \} + < y(0),u > + ny(0)_n + \frac{n(n+1)}{2} \}} \\ & & if \ y = x(k) \ and \ z = y(1) \end{cases} \\ &= \begin{cases} & (-1)^{\sum_{j=1}^{k+1} \{ < x(j-1),u > + nx(j-1)_n + \frac{n(n+1)}{2} \}} \\ & 0 \end{cases} & (z = x(k+1)) \\ &0 & (z \neq x(k+1)) \end{cases} \end{split}$$

**Proof.** (of lemma 30)

$$\begin{split} \Delta^2(z,x) &= \sum_{y \in Q^n} \Delta(z,y) \Delta(y,x) \\ &= \sum_{y \in Q^n} \left\{ \frac{(-1)^{<\delta(y),z \oplus u>}}{(\sqrt{2})^n} \times \frac{(-1)^{<\delta(x),y \oplus u>}}{(\sqrt{2})^n} \right\} \\ &= \frac{(-1)^{< x,u>}}{2^n} \sum_{y \in Q^n} (-1)^{<\delta(x \oplus z \oplus u),y>} \\ &= \left\{ \begin{array}{cc} (-1)^{< x,u>} & (x \oplus z = u) \\ 0 & (x \oplus z \neq u) \end{array} \right. \end{split}$$

**Proof.** (of lemma 32)

$$\begin{split} \Delta^2(z,x) &= \sum_{y \in Q^n} \Delta(z,y) \Delta(y,x) \\ &= \sum_{y \in Q^n} \left\{ \frac{(-1)^{<\delta(y),z \oplus u>}}{(\sqrt{2})^n} \times \frac{(-1)^{<\delta(x),y \oplus u>}}{(\sqrt{2})^n} \right\} \\ &= \frac{(-1)^{x_2 + x_3 + \dots + x_{n-1}}}{2^n} \sum_{y \in Q^n} (-1)^{<\delta(x \oplus z \oplus u),y>} \\ &= \left\{ \begin{array}{cc} (-1)^{x_2 + x_3 + \dots + x_{n-1}} & (z = x \oplus u) \\ 0 & (z \neq x \oplus u) \end{array} \right. \end{split}$$

**Proof.** (of lemma 35)

$$\begin{split} \Delta^2(z,x) &= \sum_{y \in Q^n} \Delta(z,y) \Delta(y,x) \\ &= \sum_{y \in Q^n} \left\{ \frac{(-1)^{<\delta(y),z \oplus u>}}{(\sqrt{2})^n} \times \frac{(-1)^{<\delta(x),y \oplus u>}}{(\sqrt{2})^n} \right\} \\ &= \frac{(-1)^{<\delta(x) \oplus z \oplus u,u>}}{2^n} \sum_{y \in Q^n} (-1)^{<\delta(x \oplus z \oplus u),y>} \\ &= \left\{ \begin{array}{cc} (-1)^{<\delta(x) \oplus z \oplus u,u>} & (x \oplus z \oplus u = w) \\ 0 & (x \oplus z \oplus u \neq w) \end{array} \right. \\ &= \left\{ \begin{array}{cc} (-1)^{x_1+x_n} & (z = w \oplus x \oplus u) \\ 0 & (z \neq w \oplus x \oplus u) \end{array} \right. \end{split}$$

### Appendix B Computer Simulations

In this appendix we will show three kinds of results of computer simulations for each  $QCA - [\delta_{R,bc}, \lambda_{\theta}](n)$ . We let  $c_2 = \varepsilon_{0^{n-2}10}$  and  $c_4 = \varepsilon_{0^{n-3}100}$ . The first result (lefthandside figure) shows the probability of finding state 1 at each cells by darkness. The second result (middle figure) shows the change of the mean density of cells with state 1 and the thrid result (righthandside figure) shows the discrete Fourier transiformation of the second result.

### Appendix B.1 Behaviors of $QCA - [\delta_{60,0-0}, \lambda_{\theta}](5)$

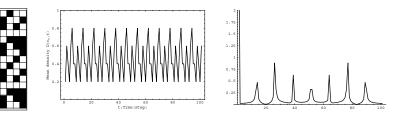


Figure 3:  $\theta = \frac{\pi}{2}$ 

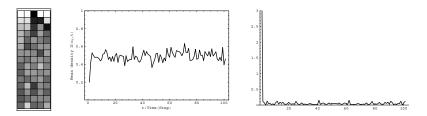


Figure 4:  $\theta = 0.35764$ 

## Appendix B.2 Behaviors of $QCA - [\delta_{204,c}, \theta](5)$

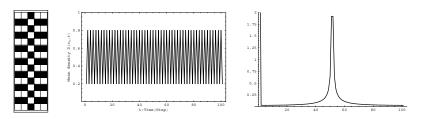


Figure 5:  $\theta = \frac{\pi}{2}$ 

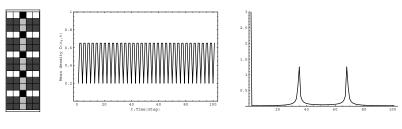


Figure 6:  $\theta = \frac{\pi}{3}$ 

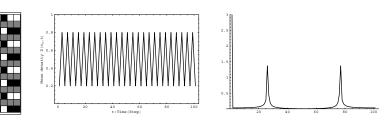


Figure 7:  $\theta = \frac{\pi}{4}$ 

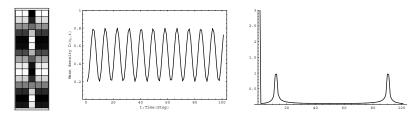


Figure 8:  $\theta = 0.35764$ 

## Appendix B.3 Behaviors of $QCA - [\delta_{51,c}, \lambda_{\theta}](5)$

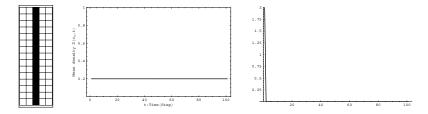


Figure 9:  $\theta = \frac{\pi}{2}$ 

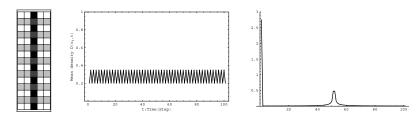


Figure 10:  $\theta = \frac{\pi}{3}$ 

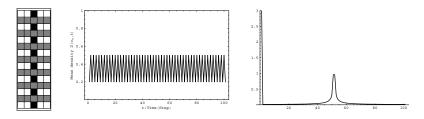


Figure 11:  $\theta = \frac{\pi}{4}$ 

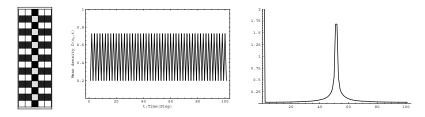


Figure 12:  $\theta = 0.35764$ 

### Appendix B.4 Behaviors of $QCA - [\delta_{240,c}, \lambda_{\theta}](5)$

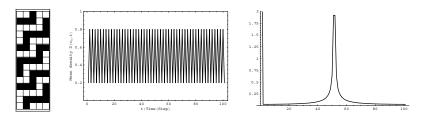


Figure 13:  $\theta = \frac{\pi}{2}$ 

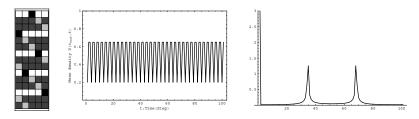


Figure 14:  $\theta = \frac{\pi}{3}$ 

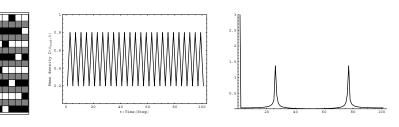


Figure 15:  $\theta = \frac{\pi}{4}$ 

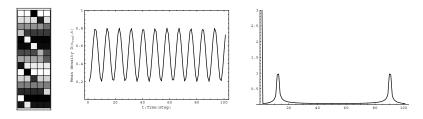


Figure 16:  $\theta = 0.35764$ 

### Appendix B.5 Behaviors of $QCA - [\delta_{90,0-0}, \lambda_{\theta}](4)$

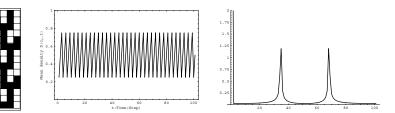


Figure 17:  $\theta = \frac{\pi}{2}$ 

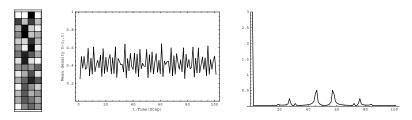


Figure 18:  $\theta = \frac{\pi}{3}$ 

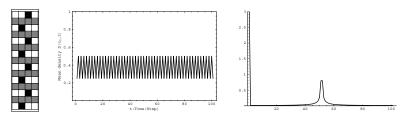


Figure 19:  $\theta = \frac{\pi}{4}$ 

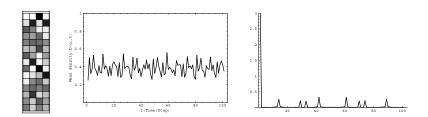


Figure 20:  $\theta = 0.35764$ 

## Appendix B.6 Behaviors of $QCA - [\delta_{165,0-0}, \lambda_{\theta}](4)$

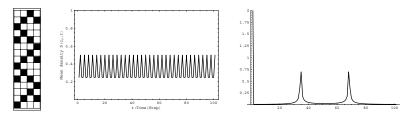


Figure 21:  $\theta = \frac{\pi}{2}$ 

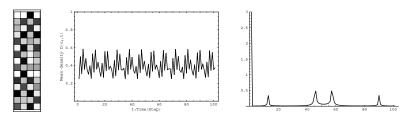


Figure 22:  $\theta = \frac{\pi}{3}$ 

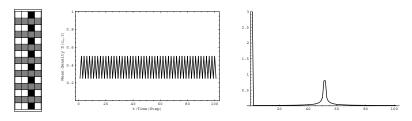


Figure 23:  $\theta = \frac{\pi}{4}$ 

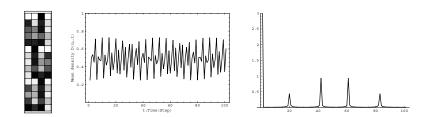


Figure 24:  $\theta = 0.35764$ 

## Appendix B.7 Behaviors of $QCA - [\delta_{195,0-0}, \lambda_{\theta}](5)$

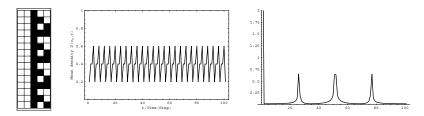


Figure 25:  $\theta = \frac{\pi}{2}$ 

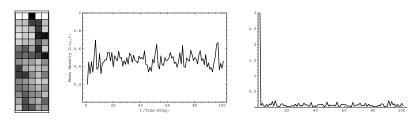


Figure 26:  $\theta = \frac{\pi}{3}$ 

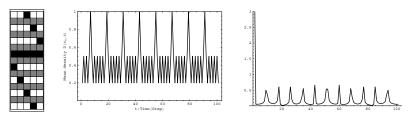


Figure 27:  $\theta = \frac{\pi}{4}$ 

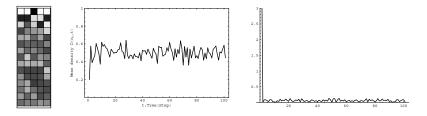


Figure 28:  $\theta = 0.35764$ 

### Appendix B.8 Behaviors of $QCA - [\delta_{150,0-0}, \lambda_{\theta}](4)$

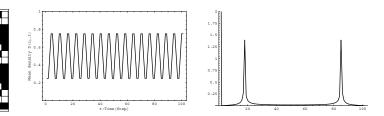


Figure 29:  $\theta = \frac{\pi}{2}$ 

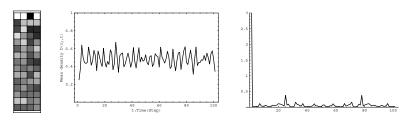


Figure 30:  $\theta = \frac{\pi}{3}$ 

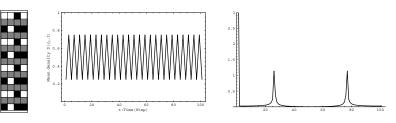


Figure 31:  $\theta = \frac{\pi}{4}$ 

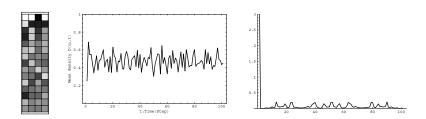


Figure 32:  $\theta = 0.35764$ 

## Appendix B.9 Behaviors of $QCA - [\delta_{150,c}, \lambda_{\theta}](5)$

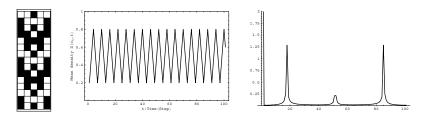


Figure 33:  $\theta = \frac{\pi}{2}$ 

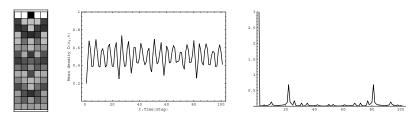


Figure 34:  $\theta = \frac{\pi}{3}$ 

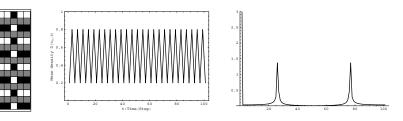


Figure 35:  $\theta = \frac{\pi}{4}$ 

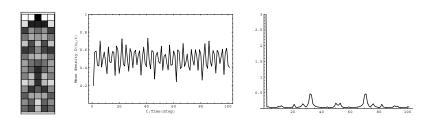


Figure 36:  $\theta = 0.35764$ 

## Appendix B.10 Behaviors of $QCA - [\delta_{105,0-0}, \lambda_{\theta}](4)$

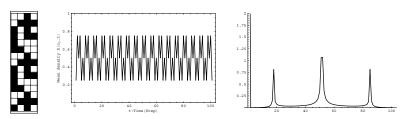


Figure 37:  $\theta = \frac{\pi}{2}$ 

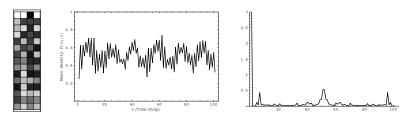


Figure 38:  $\theta = \frac{\pi}{3}$ 

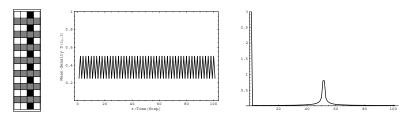


Figure 39:  $\theta = \frac{\pi}{4}$ 

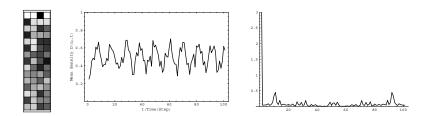


Figure 40:  $\theta = 0.35764$ 

### Appendix B.11 Behaviors of $QCA - [\delta_{105,c}, \lambda_{\theta}](5)$

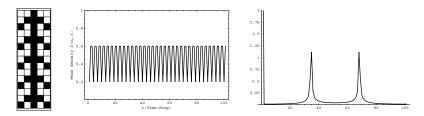


Figure 41:  $\theta = \frac{\pi}{2}$ 

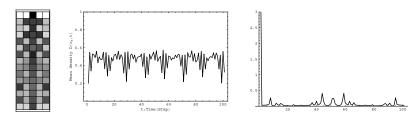


Figure 42:  $\theta = \frac{\pi}{3}$ 

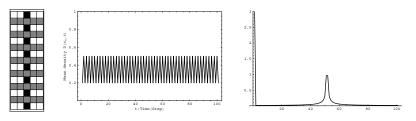


Figure 43:  $\theta = \frac{\pi}{4}$ 

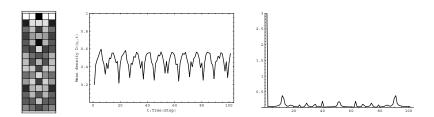


Figure 44:  $\theta = 0.35764$ 

### List of MHF Preprint Series, Kyushu University

# 21st Century COE Program Development of Dynamic Mathematics with High Functionality

MHF2005-1	Hideki KOSAKI						
	Matrix trace inec	qualities	related t	to uncer	tainty	princip]	le

- MHF2005-2 Masahisa TABATA

  Discrepancy between theory and real computation on the stability of some finite element schemes
- MHF2005-3 Yuko ARAKI & Sadanori KONISHI
  Functional regression modeling via regularized basis expansions and model selection
- MHF2005-4 Yuko ARAKI & Sadanori KONISHI Functional discriminant analysis via regularized basis expansions
- MHF2005-5 Kenji KAJIWARA, Tetsu MASUDA, Masatoshi NOUMI, Yasuhiro OHTA & Yasuhiko YAMADA
  Point configurations, Cremona transformations and the elliptic difference Painlevé equations
- MHF2005-6 Kenji KAJIWARA, Tetsu MASUDA, Masatoshi NOUMI, Yasuhiro OHTA & Yasuhiko YAMADA

  Construction of hypergeometric solutions to the q Painlevé equations
- MHF2005-7 Hiroki MASUDA Simple estimators for non-linear Markovian trend from sampled data: I. ergodic cases
- MHF2005-8 Hiroki MASUDA & Nakahiro YOSHIDA Edgeworth expansion for a class of Ornstein-Uhlenbeck-based models
- MHF2005-9 Masayuki UCHIDA Approximate martingale estimating functions under small perturbations of dynamical systems
- MHF2005-10 Ryo MATSUZAKI & Masayuki UCHIDA One-step estimators for diffusion processes with small dispersion parameters from discrete observations
- MHF2005-11 Junichi MATSUKUBO, Ryo MATSUZAKI & Masayuki UCHIDA Estimation for a discretely observed small diffusion process with a linear drift
- MHF2005-12 Masayuki UCHIDA & Nakahiro YOSHIDA AIC for ergodic diffusion processes from discrete observations

- MHF2005-13 Hiromichi GOTO & Kenji KAJIWARA Generating function related to the Okamoto polynomials for the Painlevé IV equation
- MHF2005-14 Masato KIMURA & Shin-ichi NAGATA Precise asymptotic behaviour of the first eigenvalue of Sturm-Liouville problems with large drift
- MHF2005-15 Daisuke TAGAMI & Masahisa TABATA

  Numerical computations of a melting glass convection in the furnace
- MHF2005-16 Raimundas VIDŪNAS

  Normalized Leonard pairs and Askey-Wilson relations
- MHF2005-17 Raimundas VIDŪNAS
  Askey-Wilson relations and Leonard pairs
- MHF2005-18 Kenji KAJIWARA & Atsushi MUKAIHIRA Soliton solutions for the non-autonomous discrete-time Toda lattice equation
- MHF2005-19 Yuu HARIYA Construction of Gibbs measures for 1-dimensional continuum fields
- MHF2005-20 Yuu HARIYA Integration by parts formulae for the Wiener measure restricted to subsets in  $\mathbb{R}^d$
- MHF2005-21 Yuu HARIYA A time-change approach to Kotani's extension of Yor's formula
- MHF2005-22 Tadahisa FUNAKI, Yuu HARIYA & Mark YOR Wiener integrals for centered powers of Bessel processes, I
- MHF2005-23 Masahisa TABATA & Satoshi KAIZU Finite element schemes for two-fluids flow problems
- MHF2005-24 Ken-ichi MARUNO & Yasuhiro OHTA

  Determinant form of dark soliton solutions of the discrete nonlinear Schrödinger equation
- MHF2005-25 Alexander V. KITAEV & Raimundas VIDŪNAS Quadratic transformations of the sixth Painlevé equation
- MHF2005-26 Toru FUJII & Sadanori KONISHI Nonlinear regression modeling via regularized wavelets and smoothing parameter selection
- MHF2005-27 Shuichi INOKUCHI, Kazumasa HONDA, Hyen Yeal LEE, Tatsuro SATO, Yoshihiro MIZOGUCHI & Yasuo KAWAHARA
  On reversible cellular automata with finite cell array

- MHF2005-28 Toru KOMATSU

  Cyclic cubic field with explicit Artin symbols
- MHF2005-29 Mitsuhiro T. NAKAO, Kouji HASHIMOTO & Kaori NAGATOU A computational approach to constructive a priori and a posteriori error estimates for finite element approximations of bi-harmonic problems
- MHF2005-30 Kaori NAGATOU, Kouji HASHIMOTO & Mitsuhiro T. NAKAO Numerical verification of stationary solutions for Navier-Stokes problems
- MHF2005-31 Hidefumi KAWASAKI
  A duality theorem for a three-phase partition problem
- MHF2005-32 Hidefumi KAWASAKI
  A duality theorem based on triangles separating three convex sets
- MHF2005-33 Takeaki FUCHIKAMI & Hidefumi KAWASAKI An explicit formula of the Shapley value for a cooperative game induced from the conjugate point
- MHF2005-34 Hideki MURAKAWA
  A regularization of a reaction-diffusion system approximation to the two-phase Stefan problem
- MHF2006-1 Masahisa TABATA Numerical simulation of Rayleigh-Taylor problems by an energy-stable finite element scheme
- MHF2006-2 Ken-ichi MARUNO & G R W QUISPEL Construction of integrals of higher-order mappings
- MHF2006-3 Setsuo TANIGUCHI On the Jacobi field approach to stochastic oscillatory integrals with quadratic phase function
- MHF2006-4 Kouji HASHIMOTO, Kaori NAGATOU & Mitsuhiro T. NAKAO A computational approach to constructive a priori error estimate for finite element approximations of bi-harmonic problems in nonconvex polygonal domains
- MHF2006-5 Hidefumi KAWASAKI A duality theory based on triangular cylinders separating three convex sets in  $\mathbb{R}^n$
- MHF2006-6 Raimundas VID $\bar{\text{U}}$ NAS Uniform convergence of hypergeometric series
- MHF2006-7 Yuji KODAMA & Ken-ichi MARUNO N-Soliton solutions to the DKP equation and Weyl group actions

- MHF2006-8 Toru KOMATSU
  Potentially generic polynomial
- MHF2006-9 Toru KOMATSU

  Generic sextic polynomial related to the subfield problem of a cubic polynomial
- MHF2006-10 Shu TEZUKA & Anargyros PAPAGEORGIOU

  Exact cubature for a class of functions of maximum effective dimension
- MHF2006-11 Shu TEZUKA On high-discrepancy sequences
- MHF2006-12 Raimundas VIDŪNAS Detecting persistent regimes in the North Atlantic Oscillation time series
- MHF2006-13 Toru KOMATSU

  Tamely Eisenstein field with prime power discriminant
- MHF2006-14 Nalini JOSHI, Kenji KAJIWARA & Marta MAZZOCCO Generating function associated with the Hankel determinant formula for the solutions of the Painlevé IV equation
- MHF2006-15 Raimundas VIDŪNAS

  Darboux evaluations of algebraic Gauss hypergeometric functions
- MHF2006-16 Masato KIMURA & Isao WAKANO New mathematical approach to the energy release rate in crack extension
- MHF2006-17 Toru KOMATSU

  Arithmetic of the splitting field of Alexander polynomial
- MHF2006-18 Hiroki MASUDA Likelihood estimation of stable Lévy processes from discrete data
- MHF2006-19 Hiroshi KAWABI & Michael RÖCKNER
  Essential self-adjointness of Dirichlet operators on a path space with Gibbs measures via an SPDE approach
- MHF2006-20 Masahisa TABATA

  Energy stable finite element schemes and their applications to two-fluid flow problems
- MHF2006-21 Yuzuru INAHAMA & Hiroshi KAWABI
  Asymptotic expansions for the Laplace approximations for Itô functionals of Brownian rough paths
- MHF2006-22 Yoshiyuki KAGEI Resolvent estimates for the linearized compressible Navier-Stokes equation in an infinite layer

- MHF2006-23 Yoshiyuki KAGEI
  - Asymptotic behavior of the semigroup associated with the linearized compressible Navier-Stokes equation in an infinite layer
- MHF2006-24 Akihiro MIKODA, Shuichi INOKUCHI, Yoshihiro MIZOGUCHI & Mitsuhiko FUJIO
- The number of orbits of box-ball systems MHF2006-25 Toru FUJII & Sadanori KONISHI
  - Multi-class logistic discrimination via wavelet-based functionalization and model selection criteria
- MHF2006-26 Taro HAMAMOTO, Kenji KAJIWARA & Nicholas S. WITTE Hypergeometric solutions to the q-Painlevé equation of type  $(A_1 + A'_1)^{(1)}$
- MHF2006-27 Hiroshi KAWABI & Tomohiro MIYOKAWA

  The Littlewood-Paley-Stein inequality for diffusion processes on general metric spaces
- MHF2006-28 Hiroki MASUDA
  Notes on estimating inverse-Gaussian and gamma subordinators under highfrequency sampling
- MHF2006-29 Setsuo TANIGUCHI

  The heat semigroup and kernel associated with certain non-commutative harmonic oscillators
- MHF2006-30 Setsuo TANIGUCHI Stochastic analysis and the KdV equation
- MHF2006-31 Masato KIMURA, Hideki KOMURA, Masayasu MIMURA, Hidenori MIYOSHI, Takeshi TAKAISHI & Daishin UEYAMA

  Quantitative study of adaptive mesh FEM with localization index of pattern
- MHF2007-1 Taro HAMAMOTO & Kenji KAJIWARA Hypergeometric solutions to the q-Painlevé equation of type  $A_4^{(1)}$
- MHF2007-2 Kouji HASHIMOTO, Kenta KOBAYASHI & Mitsuhiro T. NAKAO Verified numerical computation of solutions for the stationary Navier-Stokes equation in nonconvex polygonal domains
- MHF2007-3 Kenji KAJIWARA, Marta MAZZOCCO & Yasuhiro OHTA A remark on the Hankel determinant formula for solutions of the Toda equation
- MHF2007-4 Jun-ichi SATO & Hidefumi KAWASAKI Discrete fixed point theorems and their application to Nash equilibrium
- MHF2007-5 Mitsuhiro T. NAKAO & Kouji HASHIMOTO Constructive error estimates of finite element approximations for non-coercive elliptic problems and its applications

- MHF2007-6 Kouji HASHIMOTO
  - A preconditioned method for saddle point problems
- MHF2007-7 Christopher MALON, Seiichi UCHIDA & Masakazu SUZUKI Mathematical symbol recognition with support vector machines
- MHF2007-8 Kenta KOBAYASHI

On the global uniqueness of Stokes' wave of extreme form

MHF2007-9 Kenta KOBAYASHI

A constructive a priori error estimation for finite element discretizations in a non-convex domain using singular functions

MHF2007-10 Myoungnyoun KIM, Mitsuhiro T. NAKAO, Yoshitaka WATANABE & Takaaki NISHIDA

A numerical verification method of bifurcating solutions for 3-dimensional Rayleigh-Bénard problems

MHF2007-11 Yoshiyuki KAGEI

Large time behavior of solutions to the compressible Navier-Stokes equation in an infinite layer

- MHF2007-12 Takashi YANAGAWA, Satoshi AOKI and Tetsuji OHYAMA Human finger vein images are diverse and its patterns are useful for personal identification
- MHF2007-13 Masahisa TABATA

Finite element schemes based on energy-stable approximation for two-fluid flow problems with surface tension

MHF2007-14 Mitsuhiro T. NAKAO & Takehiko KINOSHITA

Some remarks on the behaviour of the finite element solution in nonsmooth domains

MHF2007-15 Yoshiyuki KAGEI & Takumi NUKUMIZU

Asymptotic behavior of solutions to the compressible Navier-Stokes equation in a cylindrical domain

MHF2007-16 Shuichi INOKUCHI, Yoshihiro MIZOGUCHI, Hyen Yeal LEE & Yasuo KAWA-HARA

Periodic Behaviors of Quantum Cellular Automata