

Properties and searches of the exotic neutral Higgs bosons in the Georgi-Machacek model

Wei Chiang, Cheng

Center for Mathematics and Theoretical Physics and Department of Physics, National Central University

TSUMURA, KOJI

Department of Physics, Kyoto University

<https://hdl.handle.net/2324/4795978>

出版情報 : Journal of High Energy Physics. 2015 (4), pp.113-, 2015-04-21. Springer Nature
バージョン :

権利関係 : (C)The Authors. Open Access. This article is distributed under the terms of the
Creative Commons Attribution 4.0 International License.



Properties and searches of the exotic neutral Higgs bosons in the Georgi-Machacek model

Cheng-Wei Chiang^{a,b,c,d} and Koji Tsumura^e

^aCenter for Mathematics and Theoretical Physics and Department of Physics,
National Central University, Taoyuan, Taiwan, 32001 R.O.C.

^bInstitute of Physics, Academia Sinica,
Taipei, Taiwan, 11529 R.O.C.

^cPhysics Division, National Center for Theoretical Sciences,
Hsinchu, Taiwan, 30013 R.O.C.

^dKobayashi-Maskawa Institute for the Origin of Particles and the Universe, Nagoya University,
Nagoya, 464-8602 Japan

^eDepartment of Physics, Kyoto University,
Kyoto, 606-8502 Japan

E-mail: chengwei@ncu.edu.tw, ko2@gauge.scphys.kyoto-u.ac.jp

ABSTRACT: The Georgi-Machacek model predicts the existence of four neutral Higgs bosons, one of which can be identified as the 125-GeV Higgs boson. The latest Higgs data favor the parameter space of small mixing angle α between the two custodial singlets of the model. The other two neutral Higgs bosons belong respectively to the custodial triplet and quintet. We study the general decay and production properties of these particles in the small- α scenario. Constraints on the $SU(2)_L$ triplet vacuum expectation value are obtained as a function of the exotic Higgs boson masses using latest ATLAS data of various search channels for additional neutral Higgs bosons.

KEYWORDS: Higgs Physics, Beyond Standard Model

ARXIV EPRINT: [1501.04257](https://arxiv.org/abs/1501.04257)

Contents

1	Introduction	1
2	Georgi-Machacek model	3
3	Branching ratios and signal strengths	6
3.1	The h boson	6
3.2	The H_1^0 boson	7
3.3	The H_3^0 boson	11
3.4	The H_5^0 boson	12
4	Experimental constraints	13
5	Summary	15

1 Introduction

For almost half a century since its introduction to particle physics, the standard model (SM) has survived many precision tests. Over the years, we have attained to knowing sufficient details of its structure and parameters from experimental data. Despite its success in explaining countless empirical observations, the SM cannot be a UV complete theory. From a theoretical point of view, the model is unsatisfactory for various reasons. For example, the Higgs boson mass is found to be at the electroweak scale and its radiative correction poses a fine-tuning issue. There are also unanswered questions such as why fermions of different flavors have such a mass hierarchy and what is the origin of CP violation in the quark sector. Experimentally, we have found the phenomena of neutrino oscillations and dark matter/energy in the Universe that the SM falls short of.

The discovery of the SM-like Higgs boson at the mass of about 125 GeV not only completes the particle spectrum in the SM, but also stimulates studies about detailed properties of the particle and pursuits of an extended Higgs family. Among many new physics models with a larger Higgs sector, the Georgi-Machacek (GM) model [1, 2] has received much attention in view of its intriguing features. The Higgs sector of GM model houses a complex triplet of hypercharge $Y = 1$ and a real triplet of $Y = 0$ under the SM $SU(2)_L \times U(1)_Y$ gauge symmetry, in addition to a SM-like doublet. With vacuum alignment between the complex and real triplets, the model preserves the custodial symmetry at tree level, granting the possibility of a triplet vacuum expectation value (VEV) as large as up to a few tens of GeV. Naturalness associated with divergent one-loop corrections to the electroweak ρ parameter and certain mixings among the Higgs bosons have been studied and found to be similar to the SM Higgs mass [3]. The model predicts the existence of many Higgs

bosons, forming two singlets, one triplet, and one quintet under the custodial symmetry [4]. There have been extensive phenomenological studies of the exotic Higgs bosons in the literature [5–11, 13–15], including their effects in enhancing the strength of phase transition in electroweak baryogenesis [16]. Due to mixing between the doublet and the triplet fields, the coupling between the SM-like Higgs boson and the weak gauge bosons can be greater than the SM value [11, 17–19], which is impossible for models with only extra $SU(2)_L$ singlet and/or doublet fields.¹ Besides, the model has a flavor-universal correction to the Yukawa couplings of the SM-like Higgs boson. Although some of the above-mentioned features are also shared by the septet extension model [20], the latter suffers from an accidental $U(1)$ symmetry that could lead to an undesirable Nambu-Goldstone (NG) boson if it is broken spontaneously. There are also supersymmetric [21, 22] and little Higgs extensions [23, 24] of the GM model.

In this paper, we concentrate on the productions and decays of the neutral Higgs bosons in the GM model at the CERN Large Hadron Collider (LHC). After the discovery of the 125-GeV Higgs boson, it is of great interest to know whether there exists another neutral Higgs boson of higher or even lower mass. Therefore, continued efforts are made to search for such Higgs-like resonances within the full mass range explorable at the LHC, utilizing various major production and decay channels. Such searches offer the exclusion limit in each channel as a function of the Higgs mass. In the GM model, there are totally four neutral Higgs bosons. We use h , H_1^0 to denote respectively the SM-like Higgs boson discovered at the LHC and the other Higgs singlet. Both of them are mixtures of the CP-even custodial singlets. The other two are the CP-odd H_3^0 belonging to the triplet and the CP-even H_5^0 in the quintet. Since h and H_1^0 are mixed states of the two singlets, they involve essentially the same production and decay processes. For definiteness in our analysis, we ignore the possibility of $\varphi \rightarrow \varphi'\varphi'$, with φ denoting any of the other neutral Higgs bosons, because of the uncertainties in the triple Higgs couplings. Due to its odd CP property, the H_3^0 boson does not couple to the weak gauge bosons at tree level. Therefore, it can only be produced via the gluon-gluon fusion (GGF) and top associated production channels. Its decay modes do not include the weak boson pairs. For sufficiently high mass, in particular, it can also decay into hZ . On the other hand, the H_5^0 boson comes purely from the weak isospin triplet fields and does not couple to the SM fermions at tree level. Consequently, it can only be produced through the vector boson fusion (VBF) and Higgs-strahlung production processes, and decay only to the final states of WW , ZZ , and $\gamma\gamma$. It is our objectives in this work to study in detail the above-mentioned features of the neutral Higgs bosons in the GM model and to constrain the model using the available search data.

This paper is organized as follows. Section 2 briefly reviews the GM model and, in particular, provides the couplings of the neutral Higgs bosons with the SM fermions and weak gauge bosons. We show how such couplings vary from their SM values as one adjusts the triplet VEV and the mixing angle α between the two singlets. In view of good agreement between the measured Higgs data and the SM predictions, we will focus on the region of small α . In section 3, we discuss the decay branching ratios and the signal

¹Mixing with an $SU(2)_L$ septet scalar field with $Y = 2$ is another possibility to obtain $\kappa_V > 1$. Such a model also predicts $\rho = 1$ without the need of vacuum alignment [20].

strengths of different production channels for each of the neutral Higgs bosons in the model. Section 4 discusses the constraints on the GM model based upon experimental searches of another SM-like Higgs bosons in various channels. This complements the other constraints from doubly-charged Higgs boson searches and indirectly from B physics and electroweak precision data. We summarize our analysis in section 5.

2 Georgi-Machacek model

The Higgs sector in the GM model comprises an isospin doublet field ϕ with $Y = 1/2$, a complex triplet field χ with $Y = 1$, and a real triplet field ξ with $Y = 0$, where the electric charge Q is given by $Q = T_3 + Y$ with T_3 being the third isospin generator. Organized in an $SU(2)_L \times SU(2)_R$ covariant form, one has:

$$\Phi = \begin{pmatrix} \phi^{0*} & \phi^+ \\ -(\phi^+)^* & \phi^0 \end{pmatrix}, \quad \Delta = \begin{pmatrix} \chi^{0*} & \xi^+ & \chi^{++} \\ -(\chi^+)^* & \xi^0 & \chi^+ \\ (\chi^{++})^* & -(\xi^+)^* & \chi^0 \end{pmatrix}, \quad (2.1)$$

where the phase convention for the component scalar fields is such that $\phi^- = (\phi^+)^*$, $\chi^{--} = (\chi^{++})^*$, $\chi^- = (\chi^+)^*$, $\xi^- = (\xi^+)^*$. Moreover, the neutral components after electroweak symmetry breaking are parameterized as

$$\phi^0 = \frac{1}{\sqrt{2}}(v_\phi + \phi_r + i\phi_i), \quad \chi^0 = v_\chi + \frac{1}{\sqrt{2}}(\chi_r + i\chi_i), \quad \xi^0 = v_\xi + \xi_r, \quad (2.2)$$

where v_ϕ , v_χ and v_ξ denote the VEV's of ϕ , χ and ξ , respectively. The most general Higgs potential consistent with the $SU(2)_L \times SU(2)_R \times U(1)_Y$ symmetry has four dimensionful and five dimensionless parameters. Its explicit form can be found, for example, in ref. [8].

When the VEV's of the two triplet fields are aligned: $v_\chi = v_\xi \equiv v_3$, the original $SU(2)_L \times SU(2)_R$ symmetry in the Higgs potential of the model breaks down to the custodial $SU(2)_V$ symmetry. In this case, the masses of the W and Z bosons have exactly the same form as in the SM: $M_W^2 = g^2 v^2/4$ and $M_Z^2 = g^2 v^2/(4 \cos^2 \theta_W)$, where θ_W is the weak mixing angle and $v^2 \equiv v_\phi^2 + 8v_3^2 = (246 \text{ GeV})^2$. Therefore, the electroweak ρ parameter keeps unity at tree level. Following the convention in ref. [25], we define the ratio of the VEV's as

$$\tan \beta = \frac{v_\phi}{2\sqrt{2}v_3}. \quad (2.3)$$

Note that $\tan \beta$ is the reciprocal of $\tan \theta_H$ used in most other works and goes to infinity in the SM limit.

The component scalar fields can be classified into irreducible representations of $SU(2)_V$ multiplets: the Φ field is decomposed as $\mathbf{2} \otimes \mathbf{2} \rightarrow \mathbf{3} \oplus \mathbf{1}$, and the Δ field as $\mathbf{3} \otimes \mathbf{3} \rightarrow \mathbf{5} \oplus \mathbf{3} \oplus \mathbf{1}$ after the $SU(2)_L \otimes SU(2)_R \rightarrow SU(2)_V$ symmetry breaking. As a result, we have one 5-plet, two 3-plet and two singlet representations under the custodial $SU(2)_V$ symmetry. The 5-plet, denoted by $H_5 = (H_5^{\pm\pm}, H_5^\pm, H_5^0)^T$, arises within the Δ field. The two 3-plet fields mix through the angle β to render a physical CP-odd Higgs 3-plet, denoted by $H_3 = (H_3^\pm, H_3^0)^T$,

Higgs boson	κ_F	κ_V
h	$\cos \alpha / \sin \beta$	$\sin \beta \cos \alpha - \sqrt{8/3} \cos \beta \sin \alpha$
H_1^0	$\sin \alpha / \sin \beta$	$\sin \beta \sin \alpha + \sqrt{8/3} \cos \beta \cos \alpha$
H_3^0	$i\eta_f \cot \beta$	0
H_5^0	0	$\kappa_W = -\cos \beta / \sqrt{3}$ and $\kappa_Z = 2 \cos \beta / \sqrt{3}$

Table 1. Scaling factors of the couplings between neutral Higgs bosons in the GM model with the SM fermions and weak gauge bosons. $\eta_f = +1$ for up-type quarks and -1 for down-type quarks and charged leptons.

and another NG 3-plet, $(G^\pm, G^0)^T$, to become the longitudinal components of the weak gauge bosons. The two CP-even singlet fields further mix by an angle α , determined by the quartic coupling constants in the Higgs potential, to produce the SM-like Higgs boson h and another physical singlet denoted by H_1^0 . Due to the custodial symmetry, Higgs bosons belonging to the same $SU(2)_V$ multiplet are degenerate in mass. Therefore, we will simply use $M_{H_{1,3,5}}$ to denote the masses of $H_{1,3,5}^0$, respectively. For later uses, we write out the physical neutral Higgs fields in terms of the original component scalar fields as follows:

$$\begin{pmatrix} h \\ H_1^0 \\ H_3^0 \end{pmatrix} = \begin{pmatrix} \cos \alpha & -\sin \alpha & 0 \\ \sin \alpha & \cos \alpha & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & \sqrt{\frac{1}{3}} & \sqrt{\frac{2}{3}} \\ 0 & -\sqrt{\frac{2}{3}} & \sqrt{\frac{1}{3}} \end{pmatrix} \begin{pmatrix} \phi_r \\ \xi_r \\ \chi_r \end{pmatrix},$$

$$H_3^0 = \sin \beta \chi_i - \cos \beta \phi_i. \quad (2.4)$$

Here we identify the quantum of the h field as the 125-GeV Higgs boson found at the LHC. Note that, the H_3^0 does not couple to the weak gauge bosons at tree level, whereas the H_5^0 does not couple to the SM fermions.

Denote the couplings of a neutral Higgs boson φ ($= h, H_1^0, H_3^0$ or H_5^0) to a SM fermion pair and a weak gauge boson in the model by $g_{\varphi ff}$ and $g_{\varphi VV}$, respectively. Define the scaling factors κ_F and κ_V as

$$\kappa_F = \frac{g_{\varphi ff}}{g_{hff}^{\text{SM}}}, \quad \text{and} \quad \kappa_V = \frac{g_{\varphi VV}}{g_{hVV}^{\text{SM}}}, \quad (2.5)$$

where g_{hff}^{SM} and g_{hVV}^{SM} refer to the corresponding values in the SM. The scaling factors for the neutral Higgs bosons in the GM model are summarized in table 1. Since H_3 is CP-odd, its coupling with fermions involves a γ_5 factor that is not reflected in the scaling factor given in the table.

As only the two custodial singlets in the model can possibly couple to both SM fermions and weak bosons simultaneously, we plot in figure 1 the contours of different pairs of scaling factors for the two custodial singlets. In figures 1a and 1b, contours are shown on the κ_V – κ_F plane for h and H_1^0 , respectively, for various representative values of v_3 and α . In fact, the blue solid contours in both plots are identical as the scaling factors of h for the phase α are the same as those of H_1^0 for the phase $\alpha + \pi/2$. This also explains that the fixed- α

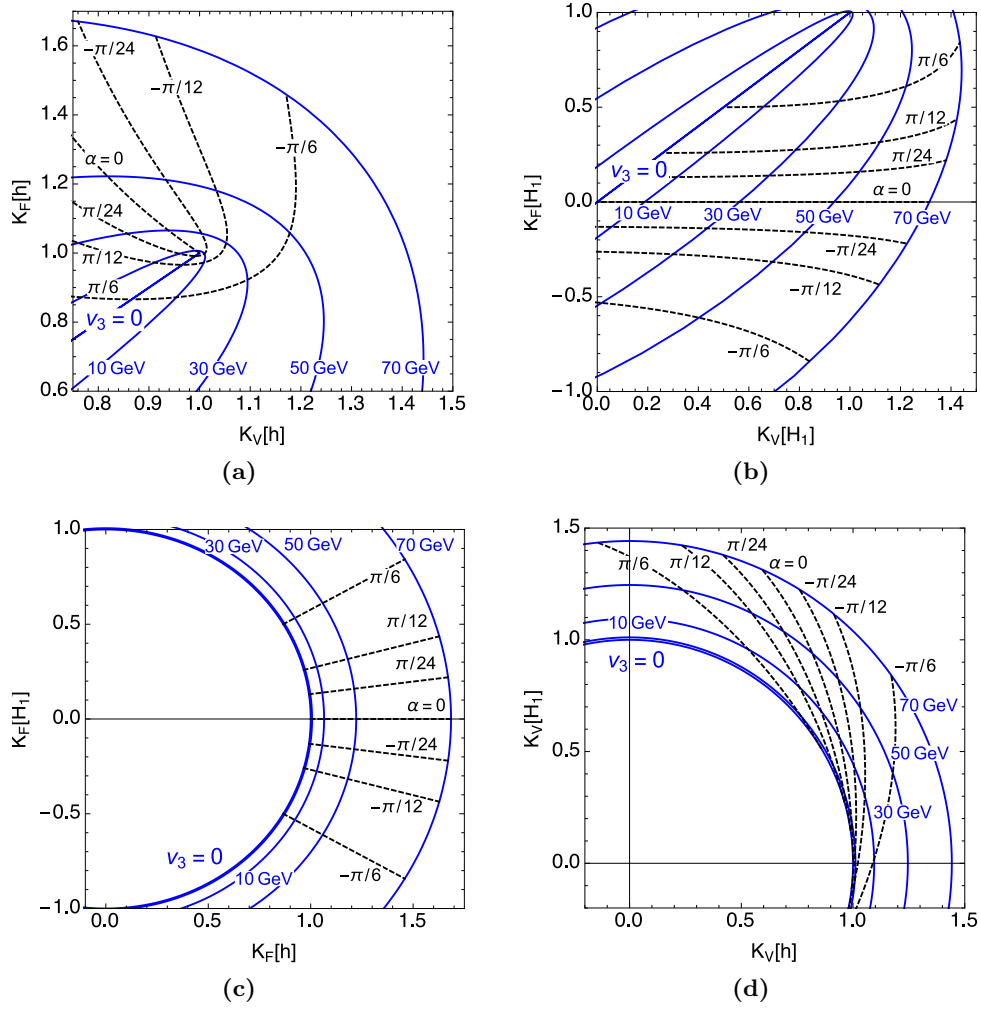


Figure 1. Scaling factors for the two custodial singlets, h and H_1^0 . The blue solid contours are for $v_3 = 0, 10, 30, 50$, and 70 GeV with varying α . The black dotted curves are for $\alpha = -\pi/6, -\pi/12, -\pi/24, 0, \pi/24, \pi/12$, and $\pi/6$ with varying β .

contours (black dotted ones) have the $\alpha \leftrightarrow \alpha + \pi/2$ correspondence between the two plots. Similarly, figure 1c shows contours on the $\kappa_F^h - \kappa_F^{H_1}$ plane, and figure 1d those on the $\kappa_V^h - \kappa_V^{H_1}$ plane. The blue solid contours of fixed v_3 values in the two plots depict a series of circles concentric at the origin, of radii $1/\sin \beta$ and $\sqrt{1 + (5/3)\cos^2 \beta}$, respectively. The only difference is that the black dotted contours of fixed α values are along the radial direction in the former, but spiral out in the latter.

Note that in the limit $\alpha = 0$, meaning that the singlet from the Φ field does not mix with that from the Δ field, the scaling factors $(\kappa_F, \kappa_V) = (1/\sin \beta, \sin \beta)$ for h and $(0, \sqrt{8/3}\cos \beta)$ for H_1^0 . Therefore, the two scaling factors of h are reciprocal to each other with $\kappa_F \geq 1$ and $\kappa_V \leq 1$, as indicated by the $\alpha = 0$ contours in figure 1a. The H_1^0 boson becomes fermiophobic but still allows a nonzero coupling with the weak gauge bosons, depending solely upon the value of v_3 , as shown by the horizontal dotted line in figure 1b. In contrast, the couplings of H_5^0 to the weak bosons are purely proportional to

v_3 , independent of α . H_3^0 also has couplings with the SM fermions as a result of mixing with the custodial triplet in the Φ field.

On the other hand, $\sin \beta \rightarrow 1$ in the limit of vanishing triplet VEV, $v_3 = 0$. In this case, $(\kappa_F, \kappa_V) = (\cos \alpha, \cos \alpha)$ for h and $(\sin \alpha, \sin \alpha)$ for H_1^0 , shown as the two blue diagonal lines in figures 1a and 1b, while both H_3^0 and H_5^0 totally decouple from the SM fermions and weak bosons.

3 Branching ratios and signal strengths

We discuss in this section patterns of the branching ratios and signal strengths for each of the neutral Higgs bosons in the model. In this paper, we define the signal strength of a particular Higgs production channel at LHC as the production rate in the GM model normalized to the SM production rate for a fictitious Higgs with the same mass:

$$\mu_X[\varphi] = \frac{\sigma^{\text{GM}}(pp \rightarrow \varphi) \mathcal{B}^{\text{GM}}(\varphi \rightarrow X)}{\sigma^{\text{SM}}(pp \rightarrow \varphi) \mathcal{B}^{\text{SM}}(\varphi \rightarrow X)}, \quad (3.1)$$

where X is some decay final state of φ , σ and \mathcal{B} denote respectively the cross section and branching ratio, and the narrow width approximation for φ has been assumed.

3.1 The h boson

Since the SM-like Higgs boson h generally has different couplings with SM particles from the SM expectation, the signal strengths of various search channels can be different from unity. Therefore, one could use the measured signal strengths to constrain the parameters α and β , as have been done in ref. [11].² Instead of performing another global fit as in ref. [11], here we want to study the effects of individual production channels. In figure 2, we plot the signal strength contours for the h boson in the model. Plot 2a shows those for the channels of $\gamma\gamma$, $f\bar{f}$, and VV^* via the GGF and top associated (ttH) productions, and plot 2b for the same channels via the VBF and Higgs-strahlung (VH) productions. In both plots, we depict the contours for each signal strength $\mu = 1$ by solid curves, as well as $\mu = 0.8$ and 1.2 by dotted curves. We note in passing that we have ignored the effects of charged Higgs bosons in the $h \rightarrow \gamma\gamma$ decays. This is justified in view of the current overall agreement of signal strengths with the SM expectations.

Because of higher statistics, the signal strengths of $ZZ^*(\rightarrow 4\ell, 2\ell 2\nu)$ and $\gamma\gamma$ channels in the GGF production are measured more accurately than in the VBF production. The latter production channels should have a much better precision at the high-luminosity LHC (HL-LHC). Due to better detector sensitivity, these bosonic channels are superior to the fermionic decay channels ($b\bar{b}$ and $\tau\tau$). The $\tau\tau$ mode in the VBF production has a relatively better sensitivity among the fermionic decays.

As shown in figure 2, the signal strengths of the bosonic channels for non-negligible v_3 have a strong dependence on α . Therefore, one can easily use their current values to exclude the parameter space of $\alpha \gtrsim +0.5$. Moreover, the region for each $0.8 \leq \mu \leq 1.2$

²The effects of the coupling deviations can also be probed by measuring the off-shell width of the SM-like Higgs boson. [12].

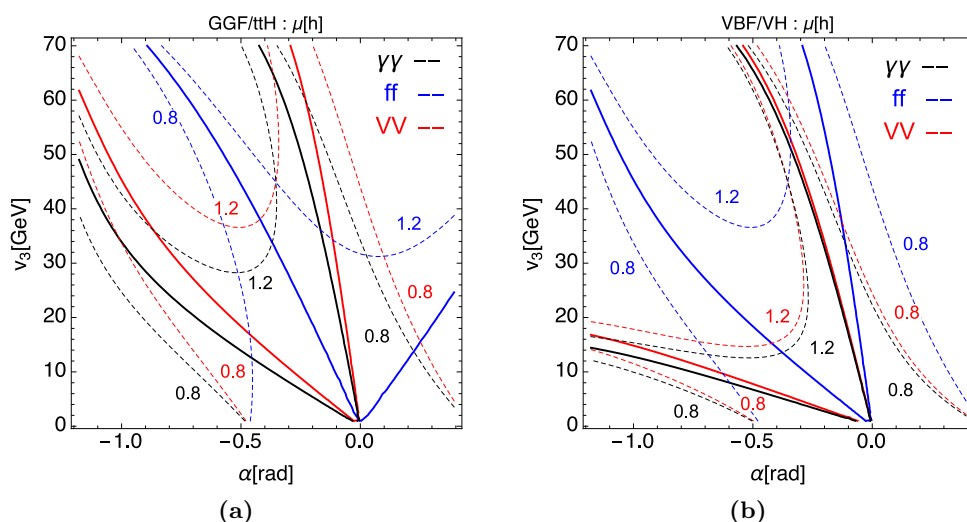


Figure 2. Signal strength contours for the h boson in (a) the GGF/ttH production channels and (b) the VBF/VH production channels. The solid curves correspond to unit signal strengths, and the dotted curves are directly labeled with the signal strength values.

have two branches at large v_3 . Thanks to the different shapes in different channels in both GGF and VBF productions, the region with $\alpha \lesssim -0.5$ in one of the branches can also be excluded. The remaining parameter space can be constrained effectively once we are able to measure the ff decay channels in the GGF process with sufficiently good precision, particularly the $\tau\tau$ channel in the 1-jet category. In addition, if the signal strengths of the $\gamma\gamma$ and VV channels in the VBF and GGF processes can be separately measured to a good accuracy using tighter selection cuts, it is also possible to constrain the remaining parameter space, especially in the large v_3 region.

We note in passing that one can perform a χ^2 fit to obtain a more specific region in the α - v_3 parameter space favoured by the current Higgs data. In this work, however, we try to keep our analysis sufficiently general, except for the choice of small α as argued above, to show the qualitative features of signal strengths for various production and decay channels, thereby learning which ones have more constraining power on the model.

In summary, the current Higgs data point us to the region of small mixing angle α . Therefore, we will focus our attention to such cases in the following analyses. More explicitly, we will consider the examples of $\alpha = -\pi/24$ and $-\pi/12$ with $v_3 = 10, 30$, and 70 GeV.

3.2 The H_1^0 boson

The decay branching ratios of the other custodial singlet, the H_1^0 boson, are shown in figure 3 for $v_3 = 10$ GeV, 30 GeV and 70 GeV. The curves are calculated using the data provided by the LHC Higgs Cross section Working Group [26] with the corresponding correction factors. For the loop-induced processes, leading-order (one-loop) correction factors are evaluated. In each plot, we draw curves for $\alpha = -\pi/24$ (solid) and $-\pi/12$ (dotted). A

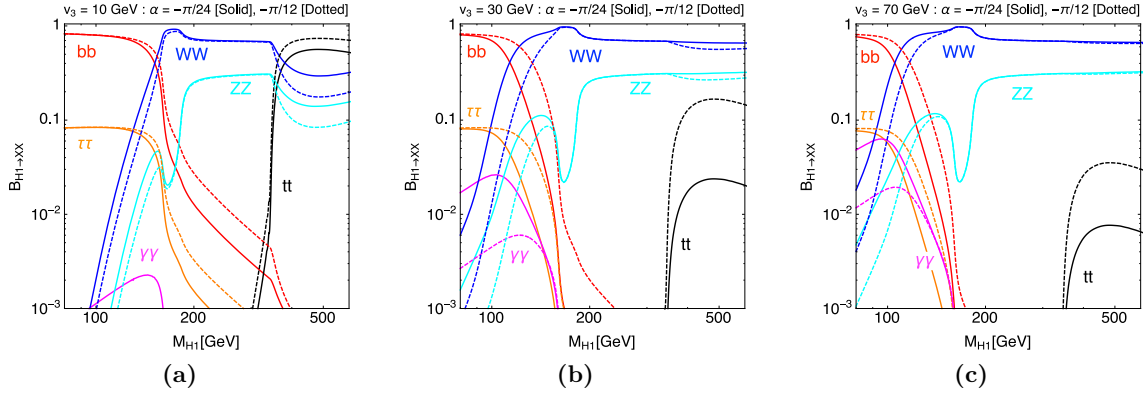


Figure 3. Decay branching ratios of the H_1^0 boson as functions of M_{H_1} . We take (a) $v_3 = 10$ GeV, (b) 30 GeV, and (c) 70 GeV.

comparison between the solid and dotted curves reveal the fermiophobic nature of the H_1^0 boson for small α , as alluded to earlier.

When $v_3 = 10$ GeV or smaller, the dominant decay mode is $b\bar{b}$ when the H_1^0 boson mass $M_{H_1} \lesssim 2M_W$ and switches to the vector boson pairs when $M_{H_1} \gtrsim 2M_W$. This feature is the same as a SM-like Higgs boson with a varying mass. The transition of dominant mode happens at a lower mass for larger v_3 . It is also observed that the decay branching ratio of the $\gamma\gamma$ mode in the low M_{H_1} regime increases with v_3 due to less destructively interfering loop processes. For sufficiently large v_3 , the dominant decay modes are the weak gauge boson pairs. Again, for definiteness, the charged Higgs contributions to $\gamma\gamma$ decay are not included in the calculation. The gg and cc decay modes are not shown in the plots.

We also note that the $H_1^0 \rightarrow hh$ decay mode is not included in the calculation, because the relevant coupling, λ_{hhH_1} , is not determined without explicitly specifying the Higgs sector parameters. Qualitatively, a nonzero λ_{hhH_1} coupling will result in a suppression of the regular search channels when $M_{H_1} \gtrsim 250$ GeV, and may lead to more production of four-body final states. Owing to its small rate, the γZ decay channel is totally neglected here for simplicity.

We now consider major production channels for the H_1^0 boson. We define the signal strength for each channel by normalizing its production rate in the GM to that for the Higgs boson in the SM but with a mass equal to M_{H_1} . Therefore, the signal strengths of the X channel from the GGF/ttH (referred to simply as GGF) and VBF/VH (referred to simply as VBF) productions can be respectively expressed as

$$\mu_X^{\text{GGF}}[H_1] = (\kappa_F^{H_1})^2 \times \frac{\mathcal{B}_X}{\mathcal{B}_X^{\text{SM}}(M_{H_1})} \simeq \frac{(\kappa_F^{H_1})^2 (\kappa_X^{H_1})^2}{(\kappa_V^{H_1})^2 \mathcal{B}_V^{\text{SM}}(M_{H_1}) + (\kappa_F^{H_1})^2 \mathcal{B}_F^{\text{SM}}(M_{H_1})}, \quad (3.2)$$

$$\mu_X^{\text{VBF}}[H_1] = (\kappa_V^{H_1})^2 \times \frac{\mathcal{B}_X}{\mathcal{B}_X^{\text{SM}}(M_{H_1})} \simeq \frac{(\kappa_V^{H_1})^2 (\kappa_X^{H_1})^2}{(\kappa_V^{H_1})^2 \mathcal{B}_V^{\text{SM}}(M_{H_1}) + (\kappa_F^{H_1})^2 \mathcal{B}_F^{\text{SM}}(M_{H_1})}, \quad (3.3)$$

where $\mathcal{B}_V^{\text{SM}}(M_{H_1})$ and $\mathcal{B}_F^{\text{SM}}(M_{H_1})$ denote the inclusive branching ratios of a SM Higgs boson of mass M_{H_1} decaying into vector boson pairs and fermion pairs, respectively, and all the other modes (e.g., multi-particle and $Z\gamma$ final states) are neglected in the last expressions.

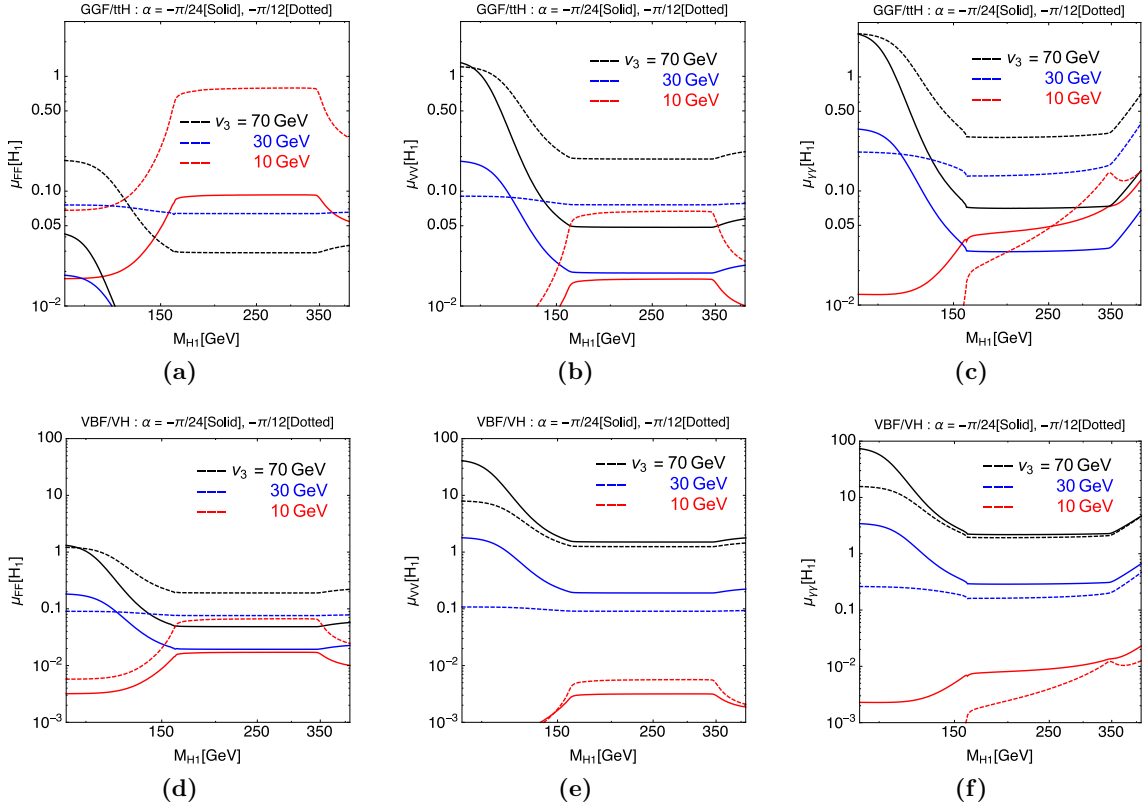


Figure 4. Signal strengths of the H_1^0 boson as functions of M_{H_1} . Plots in the left, middle, and right columns are for the $f\bar{f}$, VV , and $\gamma\gamma$ channels, respectively. Plots in the upper and lower rows are for the GGF/ $t\bar{t}H$ and VBF/ VH productions, respectively.

The signal strengths of $H_1^0 \rightarrow f\bar{f}$, VV , and $\gamma\gamma$ are drawn respectively in the left, middle, and right columns of figure 4. For each channel, we further divide the plots according to the production mechanism to the GGF process (upper row) and the VBF process (lower row).

For small α , such as $-\pi/24$ and $-\pi/12$ in our examples, the scaling factor of Yukawa couplings with H_1^0 is proportional to $\sin \alpha \sim \alpha$. In this scenario, the denominators on the right-hand side of eqs. (3.2) and (3.3) can be dominated by the inclusive bosonic branching ratio. This is the situation when v_3 is sufficiently large, as seen in figure 3. This leads to $\mu_X^{\text{VBF}}[H_1] \sim (\kappa_X^{H_1})^2 / \mathcal{B}_V^{\text{SM}}(M_{H_1})$ for large v_3 . The $1/\mathcal{B}_V^{\text{SM}}(M_{H_1})$ factor is always an enhancement factor that is large when $M_{H_1} \lesssim 2M_W$ and the bosonic decay modes are subdominant, and gets close to unity when $M_{H_1} \gtrsim 2M_W$. This is seen in figure 4 by the decline in the curves for $v_3 = 30$ and 70 GeV. In contrast, the signal strengths through the GGF process $\mu_X^{\text{GGF}}[H_1]$ is scaled from $\mu_X^{\text{VBF}}[H_1]$ by the factor $(\kappa_F^{H_1}/\kappa_V^{H_1})^2$, which is smaller for larger v_3 values. Therefore, one observes this general trend in the plots except when $v_3 = 10$ GeV.

In the $M_{H_1} \gtrsim 2M_W$ region, the predicted signal strengths are relatively flat until the $t\bar{t}$ channel opens up. This reflects the fact that the production rates vary in the same way

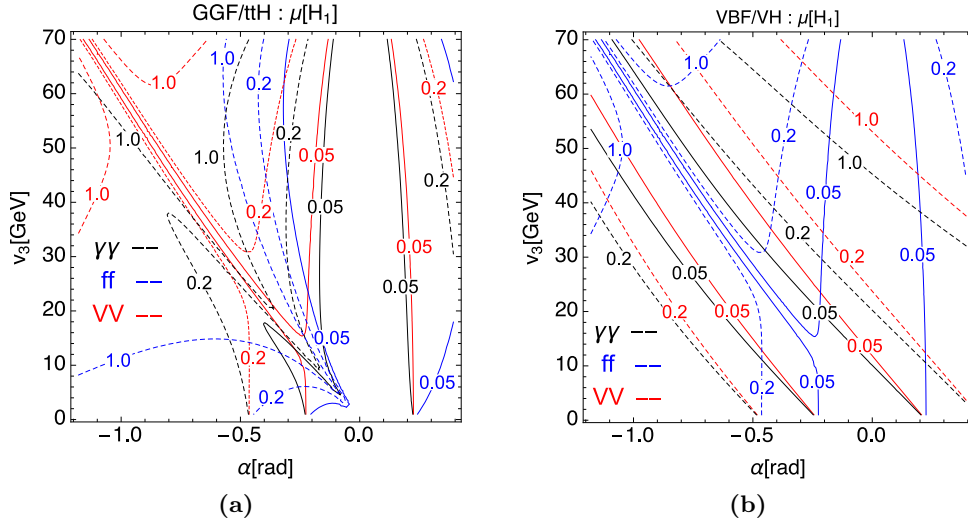


Figure 5. Contours of signal strengths for a high-mass H_1^0 boson in (a) the GGF production and (b) the VBF production shown on the α - v_3 plane.

as the SM Higgs boson in the higher mass region. Finally, it should be noted that the possibility of having the $H_1^0 \rightarrow hh$ decay is not considered here to avoid the uncertainty in the λ_{hhH_1} coupling.

As we showed in figure 3, the $\gamma\gamma$ decay mode is important only in low mass region. In the right column, we find enhancements in the signal strength of $\gamma\gamma$ decay for both the VBF and GGF production channels. This enhancement is caused by the kinematical (off-shell) reduction in the VV decay channels, suppressed fermionic decays by the v_3/v_ϕ factor, and less destructive interference between the W and top loops. The enhancement is substantial for the VBF process, as the GGF process is suppressed by the fermionic scaling factor $\kappa_F^{H_1}$.

Because the H_1^0 decay branching fractions have relatively stable predictions for $M_{H_1} \gtrsim 2M_Z$, let's focus for the moment on this mass regime as the corresponding signal strengths also give stable predictions. The VV decay channels are important once their on-shell decays become open. A significant fraction of H_1^0 decays into VV , ~ 60 (~ 30)% for the $WW(ZZ)$ mode, similarly to the branching ratios of a high-mass SM Higgs boson.

Figure 5 shows the signal strengths for various channels of H_1^0 . Plot 5a involves the channels through the GGF production, and plot 5b those through the VBF production. The signal strengths of fermionic decay modes are approximately symmetric under $\alpha \rightarrow -\alpha$ and vanish for $\alpha = 0$. The singular behaviors along the diagonal line of $v_3 \sim -(\sqrt{3}/8)v\alpha$ is caused by $\kappa_V^{H_1} \sim 0$. Since we are mainly interested in $|\alpha| \lesssim 0.5$, the fermionic decay channels have less restricting power in this part of the parameter space. The bosonic decay modes in the GGF process are also insensitive to parameter variation within the region. Considering the region $\alpha > -0.5$, the ZZ and WW decay channels in the VBF process have a potential to constrain the parameter space with $v_3 \gtrsim 40$ GeV if the signal strength is found to be less than 0.05. The decay branching ratio of $H_1 \rightarrow \gamma\gamma$ is smaller than 10^{-3} , a rate that is difficult to probe in the VBF production.

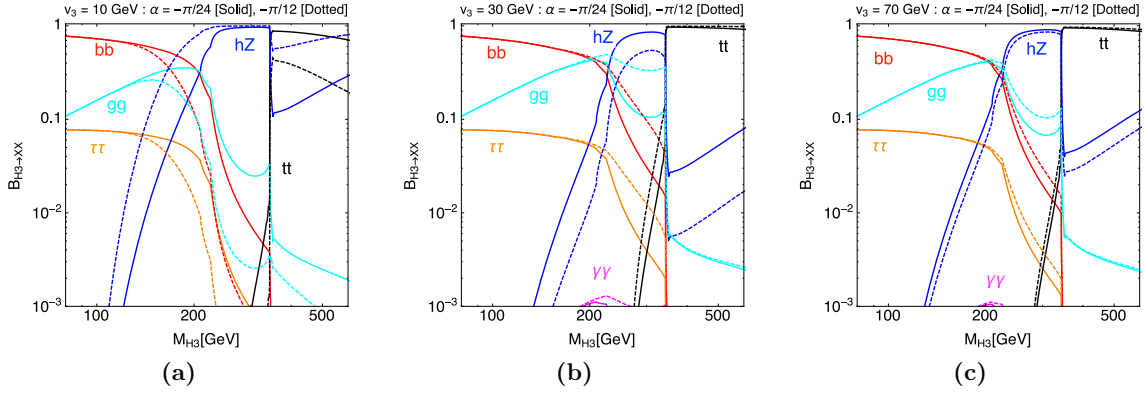


Figure 6. Branching ratios of the H_3^0 boson as functions of M_{H_3} . $v_3 = 10$ GeV, 30 GeV and 70 GeV in plots (a), (b), and (c), respectively.

3.3 The H_3^0 boson

Note that the scaling factor for the couplings of H_3^0 to SM fermions is universally $\pm \cot \beta$, where the sign is determined by the sign of the third isospin number of the fermion. Being a CP-odd particle, its fermion pair decay width is enhanced from that of a CP-even Higgs boson by the factor of $\left(1 - 4M_f^2/M_{H_3}^2\right)^{-1}$. Besides, the H_3^0 boson does not couple to the weak bosons at tree level and therefore has no weak boson pair decay modes. Moreover, its gg and $\gamma\gamma$ decays are mediated only by fermions in the loop. The particle has, however, a tree-level coupling with the SM-like Higgs boson and the Z boson. In fact, the hZ mode becomes dominant once the decay is on mass shell, until the $t\bar{t}$ decay is also on shell. The H_3^0 - h - Z coupling has dependences on both α and β (or v_3). The patterns of branching ratios as functions of the H_3^0 mass M_{H_3} are shown in figure 6.

As shown in figure 6, the $b\bar{b}$ mode remains most dominant in the low-mass region up to the hZ decay threshold. When the $t\bar{t}$ decay is kinematically allowed, it becomes the most dominant mode. The hZ decay mode is important once again in the further higher mass region. Incidentally, the branching ratio curve of the $\gamma\gamma$ mode has exactly the same behavior as the gg curve, but down by the factor of $(Q_t \alpha_{\text{EM}}/\alpha_S)^2 \simeq 0.002$, where α_S and α_{EM} denote respectively the strong and electromagnetic couplings and $Q_t = +2/3$.

Since the H_3^0 does not couple to the weak gauge bosons, one can only make use of the $f\bar{f}$ modes through the GGF production mechanism. In such cases, the signal strengths for the on-shell decays in the Born approximation are

$$\mu_{FF}^{\text{GGF}}[H_3] = (\kappa_F^{H_3})^2 \frac{F_{1/2}^A(M_{H_3})}{F_{1/2}^S(M_{H_3})} \times \frac{\mathcal{B}_X}{\mathcal{B}_X^{\text{SM}}(M_{H_3})} \left(1 - \frac{4M_f^2}{M_{H_3}^2}\right)^{-1}, \quad (3.4)$$

where $F_{1/2}^S(M)$ and $F_{1/2}^A(M)$ are the fermionic loop functions for CP-even and -odd scalar bosons [27], respectively, and $(\kappa_F^{H_3})^2 = \cot^2 \beta$, which goes to zero in the small v_3 limit. Even though the above formula is only approximate, some general properties can be extracted. Since the decays of a SM Higgs boson in the $M_{H_3} \lesssim 2M_W$ region are dominated by the fermion pairs, $\mu_X^{\text{GGF}}[H_3] \sim \cot^2 \beta$. On the other hand, when $M_{H_3} > 2M_W$, the inclusive

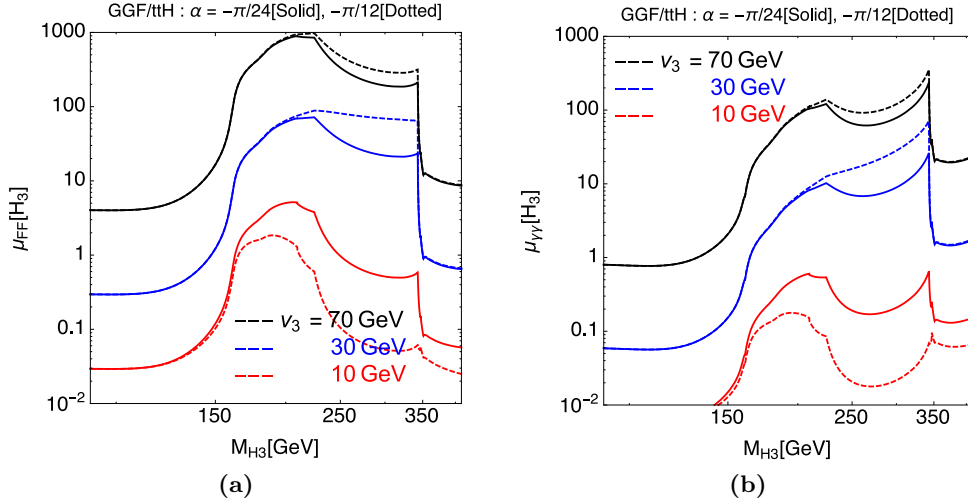


Figure 7. Signal strengths of the H_3^0 boson in the GGF production of (a) the $f\bar{f}$ channels and (b) the $\gamma\gamma$ channel as functions of M_{H_3} .

$\mathcal{B}_F^{\text{SM}}(M_{H_3})$ is very small. Therefore, the signal strengths of the fermionic modes have a significant increase at around $2M_W$. This enhancement is slightly reduced when the hZ mode opens up and further reduced above the $t\bar{t}$ threshold.

In figure 7, we show the signal strengths of the $f\bar{f}$ and $\gamma\gamma$ channels through the GGF production. It is observed that the curves scale with $\cot^2 \beta$. Also, both channels have a significant enhancement in the mass range $2M_W \lesssim M_{H_3} \lesssim M_Z + M_h$. Therefore, a search of these GGF channels in this mass range can put stringent constraints on the model. Regarding the regions of $M_{H_3} \lesssim 2M_W$ and $M_{H_3} \gtrsim 2M_t$, the $f\bar{f}$ channel is smaller than the SM normalization when $v_3 \lesssim 30$ GeV. Although having a signal strength greater than 1 for $v_3 \gtrsim 30$ GeV and $M_{H_3} \gtrsim 200$ GeV, the $\gamma\gamma$ channel is not useful as its production cross section is actually very small.

3.4 The H_5^0 boson

Since the H_5^0 boson does not couple to the SM fermions at tree level, the only search channels are the VV and $\gamma\gamma$ modes via the VBF production mechanism. Note that as given in table 1, both its couplings with the W and Z bosons scale with $\cos \beta$ (or equivalently v_3) without any dependence on α . But the latter is larger than the former by a factor of 2. The pattern of branching ratios is independent of both α and β (or v_3) and plotted as functions of the H_5^0 mass M_{H_5} in figure 8.

The crossing between the curves for the WW and ZZ channels at both low-mass end and when $M_{H_5} \gtrsim 190$ GeV is a result of the non-universal scaling behaviors in the couplings with the weak bosons. The ratio in the high mass region is $WW : ZZ = 1 : 2$ (cf. $2 : 1$ in the SM). This is a nice discriminant for the neutral Higgs boson originated from the custodial quintet. Because of the absence of fermionic decay modes, the fractions of VV and $\gamma\gamma$ decay channels are much larger than in the SM. Furthermore, the $\gamma\gamma$ decay channel can be significant for $M_{H_5} \lesssim 2M_W$ due to the off-shell suppression of VV decay modes.

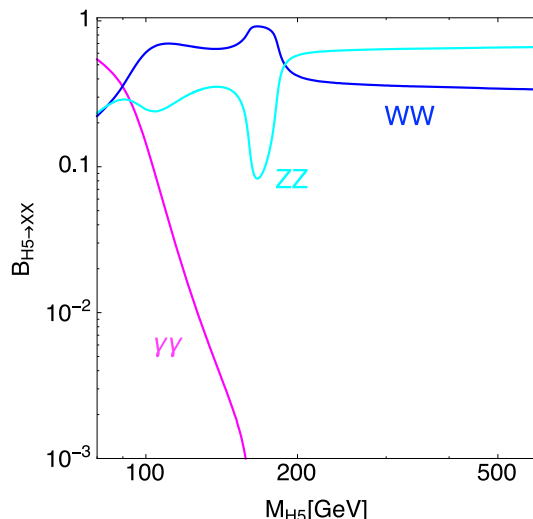


Figure 8. Branching ratios of the H_5^0 boson as functions of M_{H_5} .

Again, contributions from the charged Higgs bosons to the $\gamma\gamma$ decay is neglected because we do not specify the Higgs potential explicitly.

The signal strengths of the ZZ , WW , and $\gamma\gamma$ modes through the VBF mechanism are plotted in figure 9. Note that there is no α dependence in the curves, and the v_3 dependence comes from the production part as the decay branching ratios are independent of both α and β . The VBF cross section given by the Higgs Cross section Working Group [26] is the sum of WW fusion and ZZ fusion cross sections. In order to reflect their respective modifications in the production part, we have evaluated separately the WW fusion and ZZ fusion cross sections using MadGraph5_aMC@NLO [28]. The ratio is found to be about 3 : 1 for a wide range of M_{H_5} . As expected from the prediction of branching ratios, the signal strengths for VV and $\gamma\gamma$ decay channels are both enhanced substantially in the low mass region, $M_{H_5} \lesssim 2M_W$.

4 Experimental constraints

Searches of another SM-like Higgs boson with a different mass have been done at the LHC for various channels. Upper limits at 95% confidence level (CL) are provided as a function of the extra Higgs mass in a wide range, from as small as 60 GeV up to 1 TeV depending on the search channels. We can then take such data as constraints on the searches of the exotic neutral Higgs bosons in the GM model.

In analyzing the experimental data, we find that the ATLAS results are particularly convenient to use because they are all taken from the 8-TeV run with integrated luminosities about 20 fb^{-1} , whereas the CMS results mix measurements in both 7-TeV and 8-TeV runs. Therefore, we will use exclusively the ATLAS results in the following analysis. Up to now, the $H \rightarrow ZZ \rightarrow 4\ell$ mode is analyzed separately in both GGF and VBF+VH production processes [29], and the $H \rightarrow WW \rightarrow e\nu\mu\nu$ mode in both GGF and VBF production processes [30]. The $H \rightarrow \gamma\gamma$ mode is divided into low-mass (65–110 GeV) and high-mass

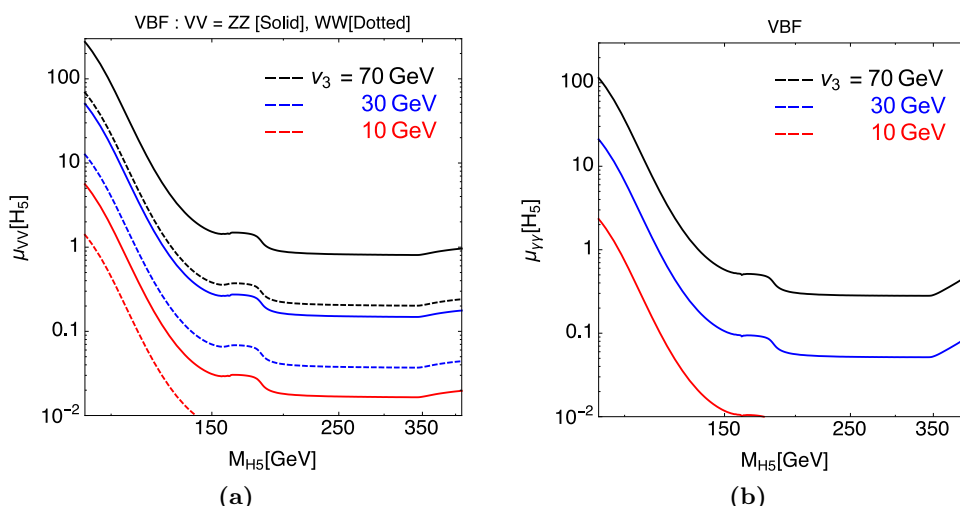


Figure 9. Signal strengths of the H_5^0 boson as functions of M_{H_5} .

(110–600 GeV) regions [31]. Since no similar data analysis using the fermionic modes is available yet, currently we can only extract constraints on the model through the studies of the H_1^0 and H_5^0 bosons. However, as we noted in section 3.3, the fermionic channels in the mass region of $2M_W \lesssim M_{H_3} \lesssim M_Z + M_h$ can easily put stringent constraints on the model parameter space.

Upper limits on the triplet VEV at 95% CL from the ZZ and WW modes are plotted against the H_1^0 mass in figure 10. These curves are obtained by scaling the theory expectation for the GM model and re-extracting the corresponding 95% CL limit on the signal cross section times the branching ratio for each channel given in refs. [29, 30]. It is clear that in both channels, the VBF process has a stronger constraint than the GGF process. This is consistent with the earlier observation that the H_1^0 boson is fermiophobic in the small α scenario considered here. Comparing figure 10b and figure 10d, one notices that the ZZ channel is generally more constraining than the WW channel except for the region $375 \lesssim M_{H_1} \lesssim 450$ GeV, in which the former (latter) has a slightly worse (better) sensitivity experimentally.

In figure 11, we plot the upper limit of the triplet VEV as a function of M_{H_1} from the $\gamma\gamma$ decay mode. In order to compare with the experimental data, we need to have an efficiency factor C_X [31], which is contained in the calculation of the fiducial cross section. The mass dependence of C_X is not presented in ref. [31]; only the range of 0.56 to 0.71 for the C_X factor is mentioned. We here take $C_X = 0.56$ to calculate a *conservative* limit on the triplet VEV. As shown in figure 11, the data are more (less) constraining for smaller α in the lower (higher) mass regime.

Finally, figure 12 shows the upper limit of the triplet VEV as a function of M_{H_5} from the $ZZ \rightarrow 4\ell$ decay mode via the VBF mechanism. Clearly, this constraint from the search of H_5^0 is weaker than those presented in figures 10 and 11, where we focus on the small mixing scenario ($-\pi/12 \leq \alpha \leq 0$). This is related to the fact that the signal strength is mainly enhanced in the low-mass region only. We note in passing that no useful

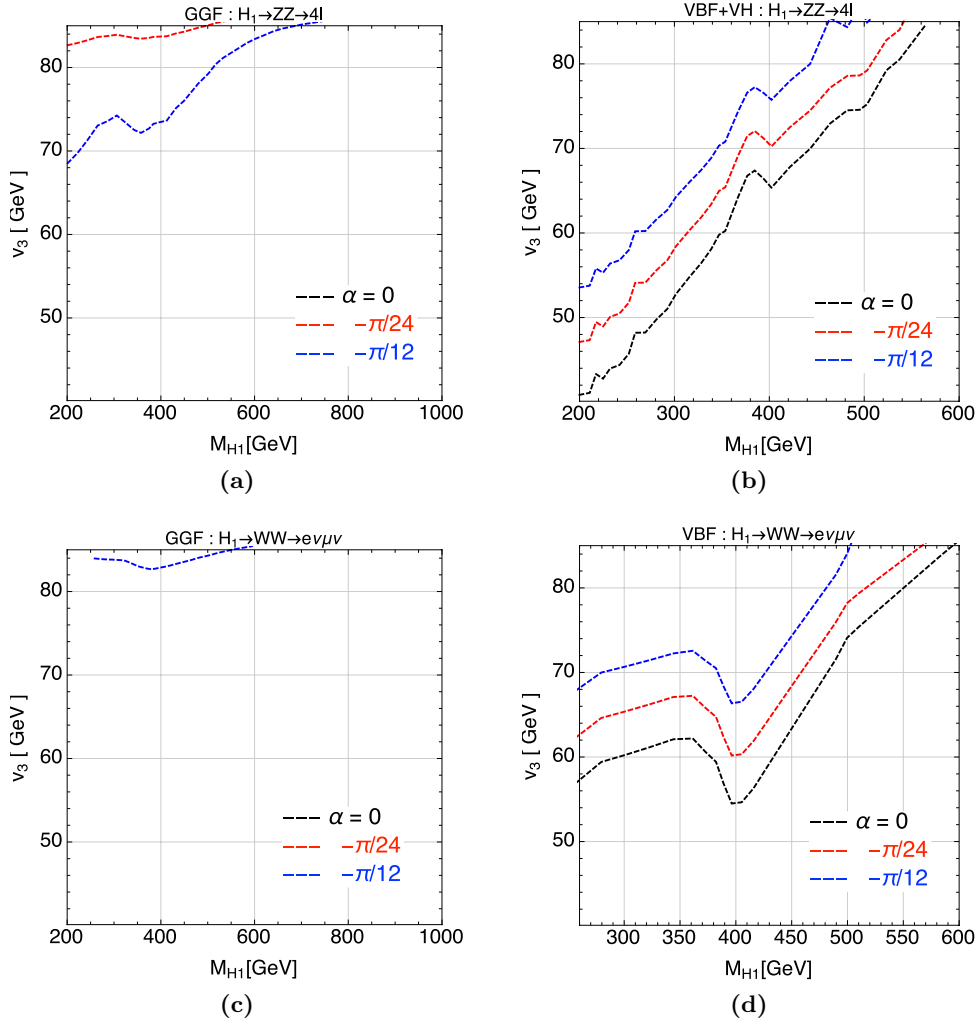


Figure 10. Upper limits on v_3 as a function of M_{H_1} for different channels and production mechanisms. Plot (a) involves the GGF mechanism, plot (b) the VBF+VH mechanism, plot (c) the GGF mechanism, and plot (d) the VBF mechanism. From top to bottom are for the WW and ZZ modes.

constraint can be obtained from the WW mode. Similar to $H_1^0 \rightarrow \gamma\gamma$, one can also extract constraints from the $H_5^0 \rightarrow \gamma\gamma$ mode, as shown in figure 13. The $H_5^0 \rightarrow \gamma\gamma$ channel may pose a stringent bound on v_3 in the small M_{H_5} region. However, there is an uncertainty in the Higgs trilinear couplings omitted in this analysis.

5 Summary

The Georgi-Machacek model predicts the existence of four neutral Higgs bosons, one of which can be identified as the discovered SM-like Higgs boson h and the other three H_1^0 , H_3^0 and H_5^0 belong respectively to custodial singlet, triplet and quintet. In addition to restraining the model parameters using the SM-like Higgs data, one can use the data of searching for additional neutral Higgs bosons at other masses to put constraints. Taking

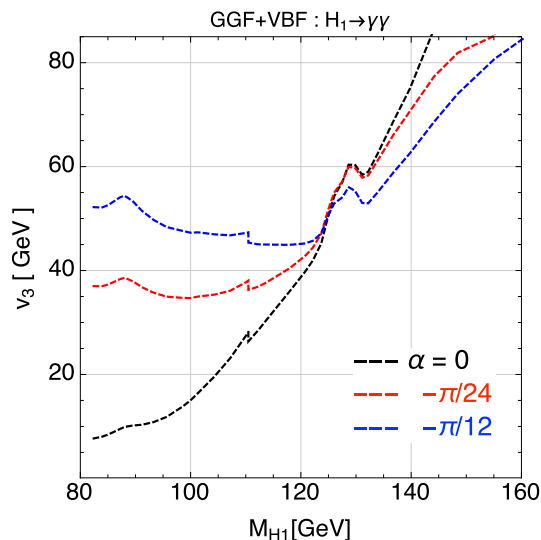


Figure 11. Upper limits on v_3 as a function of M_{H_1} for the $\gamma\gamma$ channels using events produced through the GGF, VBF, and VH mechanisms.

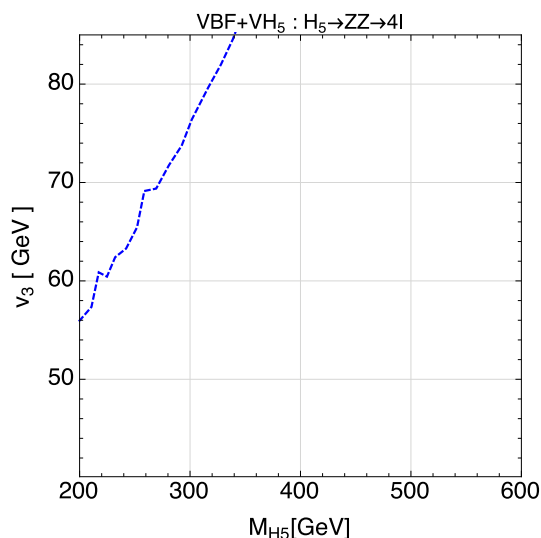


Figure 12. Upper limits on v_3 as a function of M_{H_5} for the $ZZ \rightarrow 4\ell$ channels through the VBF mechanism.

the cue from current SM-like Higgs data, we focus on the parameter space of small mixing angle α and study the decay branching ratios and signal strengths of these Higgs bosons at the CERN LHC. We have examined in this work how they vary with α , the triplet VEV v_3 , and the corresponding Higgs mass.

As a result of mixing with h , H_1^0 share all the decay modes of h . Moreover, in the small α scenario considered in this work, the weak boson decay modes are dominant in the high-mass region for large triplet VEV v_3 . Due to its odd CP property, the H_3^0 does not have the weak boson channels. Its dominant decay mode from the low-mass region up to about 600 GeV is $b\bar{b}$, hZ , and $t\bar{t}$. In the case of H_5^0 , it mainly decays to the ZZ , WW , and $\gamma\gamma$ states. In the high-mass region, the branching ratio of ZZ is about twice that of WW .

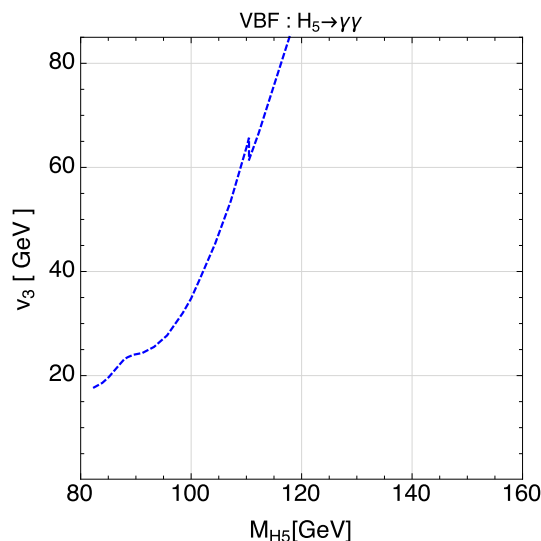


Figure 13. Upper limits on v_3 as a function of M_{H_5} for the $\gamma\gamma$ channels through the VBF mechanism.

Using the latest search data of extra Higgs boson from ATLAS, we have put constraints on v_3 as a function of the Higgs mass for the H_1^0 and H_5^0 bosons. For H_1^0 , we have employed the $ZZ \rightarrow 4\ell$, $WW \rightarrow e\nu\mu\nu$, and $\gamma\gamma$ modes through the gluon-gluon fusion (GGF) and vector boson fusion (VBF) productions. The constraints from the VBF mechanism are stronger than those from the GGF mechanism, and the results have dependence on the α angle. For H_5^0 , we have found that only the $ZZ \rightarrow 4\ell$ mode is useful and renders a weaker constraint than those of H_1^0 . Nevertheless, this constraint is independent of α .

We have found that the fermionic channels of H_3^0 are enhanced a lot in the mass range between $2M_W$ and $2M_t$. Therefore, a search of such channels in this regime can readily discover the particle or put stringent constraints on the model parameters.

Acknowledgments

This research was supported in part by the Ministry of Science and Technology of Taiwan under Grant No. NSC 100-2628-M-008-003-MY4 and in part by the MEXT Grant-in-Aid for Scientific Research on Innovative Areas No. 26104704. C.W.C. would like to thank the hospitality of the Kobayashi-Maskawa Institute at Nagoya University where this work was done during his sabbatical visit. He also thanks Chia-Ming Kuo for detailed explanations of some experimental data.

Open Access. This article is distributed under the terms of the Creative Commons Attribution License ([CC-BY 4.0](https://creativecommons.org/licenses/by/4.0/)), which permits any use, distribution and reproduction in any medium, provided the original author(s) and source are credited.

References

- [1] H. Georgi and M. Machacek, *Doubly Charged Higgs Bosons*, *Nucl. Phys. B* **262** (1985) 463 [[INSPIRE](#)].
- [2] M.S. Chanowitz and M. Golden, *Higgs Boson Triplets With $M_W = M_Z \cos \theta_W$* , *Phys. Lett. B* **165** (1985) 105 [[INSPIRE](#)].
- [3] J.F. Gunion, R. Vega and J. Wudka, *Naturalness problems for $\rho = 1$ and other large one loop effects for a standard model Higgs sector containing triplet fields*, *Phys. Rev. D* **43** (1991) 2322 [[INSPIRE](#)].
- [4] J.F. Gunion, R. Vega and J. Wudka, *Higgs triplets in the standard model*, *Phys. Rev. D* **42** (1990) 1673 [[INSPIRE](#)].
- [5] H.E. Haber and H.E. Logan, *Radiative corrections to the $Zb\bar{b}$ vertex and constraints on extended Higgs sectors*, *Phys. Rev. D* **62** (2000) 015011 [[hep-ph/9909335](#)] [[INSPIRE](#)].
- [6] S. Godfrey and K. Moats, *Exploring Higgs Triplet Models via Vector Boson Scattering at the LHC*, *Phys. Rev. D* **81** (2010) 075026 [[arXiv:1003.3033](#)] [[INSPIRE](#)].
- [7] C.-W. Chiang, T. Nomura and K. Tsumura, *Search for doubly charged Higgs bosons using the same-sign diboson mode at the LHC*, *Phys. Rev. D* **85** (2012) 095023 [[arXiv:1202.2014](#)] [[INSPIRE](#)].
- [8] C.-W. Chiang and K. Yagyu, *Testing the custodial symmetry in the Higgs sector of the Georgi-Machacek model*, *JHEP* **01** (2013) 026 [[arXiv:1211.2658](#)] [[INSPIRE](#)].
- [9] C. Englert, E. Re and M. Spannowsky, *Triplet Higgs boson collider phenomenology after the LHC*, *Phys. Rev. D* **87** (2013) 095014 [[arXiv:1302.6505](#)] [[INSPIRE](#)].
- [10] C. Englert, E. Re and M. Spannowsky, *Pinning down Higgs triplets at the LHC*, *Phys. Rev. D* **88** (2013) 035024 [[arXiv:1306.6228](#)] [[INSPIRE](#)].
- [11] C.-W. Chiang, A.-L. Kuo and K. Yagyu, *Enhancements of weak gauge boson scattering processes at the CERN LHC*, *JHEP* **10** (2013) 072 [[arXiv:1307.7526](#)] [[INSPIRE](#)].
- [12] H.E. Logan, *Hiding a Higgs width enhancement from off-shell $gg (\rightarrow h^*) \rightarrow ZZ$ measurements*, [arXiv:1412.7577](#) [[INSPIRE](#)].
- [13] K. Hartling, K. Kumar and H.E. Logan, *The decoupling limit in the Georgi-Machacek model*, *Phys. Rev. D* **90** (2014) 015007 [[arXiv:1404.2640](#)] [[INSPIRE](#)].
- [14] C.-W. Chiang, S. Kanemura and K. Yagyu, *Novel constraint on the parameter space of the Georgi-Machacek model with current LHC data*, *Phys. Rev. D* **90** (2014) 115025 [[arXiv:1407.5053](#)] [[INSPIRE](#)].
- [15] K. Hartling, K. Kumar and H.E. Logan, *Indirect constraints on the Georgi-Machacek model and implications for Higgs boson couplings*, *Phys. Rev. D* **91** (2015) 015013 [[arXiv:1410.5538](#)] [[INSPIRE](#)].
- [16] C.-W. Chiang and T. Yamada, *Electroweak phase transition in Georgi-Machacek model*, *Phys. Lett. B* **735** (2014) 295 [[arXiv:1404.5182](#)] [[INSPIRE](#)].
- [17] H.E. Logan and M.-A. Roy, *Higgs couplings in a model with triplets*, *Phys. Rev. D* **82** (2010) 115011 [[arXiv:1008.4869](#)] [[INSPIRE](#)].
- [18] A. Falkowski, S. Rychkov and A. Urbano, *What if the Higgs couplings to W and Z bosons are larger than in the Standard Model?*, *JHEP* **04** (2012) 073 [[arXiv:1202.1532](#)] [[INSPIRE](#)].

- [19] S. Chang, C.A. Newby, N. Raj and C. Wanotayaroj, *Revisiting Theories with Enhanced Higgs Couplings to Weak Gauge Bosons*, *Phys. Rev. D* **86** (2012) 095015 [[arXiv:1207.0493](#)] [[INSPIRE](#)].
- [20] J. Hisano and K. Tsumura, *Higgs boson mixes with an SU(2) septet representation*, *Phys. Rev. D* **87** (2013) 053004 [[arXiv:1301.6455](#)] [[INSPIRE](#)].
- [21] L. Cort, M. Garcia and M. Quirós, *Supersymmetric Custodial Triplets*, *Phys. Rev. D* **88** (2013) 075010 [[arXiv:1308.4025](#)] [[INSPIRE](#)].
- [22] M. Garcia-Pepin, S. Gori, M. Quirós, R. Vega, R. Vega-Morales and T.-T. Yu, *Supersymmetric Custodial Higgs Triplets and the Breaking of Universality*, *Phys. Rev. D* **91** (2015) 015016 [[arXiv:1409.5737](#)] [[INSPIRE](#)].
- [23] S. Chang and J.G. Wacker, *Little Higgs and custodial SU(2)*, *Phys. Rev. D* **69** (2004) 035002 [[hep-ph/0303001](#)] [[INSPIRE](#)].
- [24] S. Chang, *A ‘Littlest Higgs’ model with custodial SU(2) symmetry*, *JHEP* **12** (2003) 057 [[hep-ph/0306034](#)] [[INSPIRE](#)].
- [25] S. Kanemura, K. Tsumura, K. Yagyu and H. Yokoya, *Fingerprinting nonminimal Higgs sectors*, *Phys. Rev. D* **90** (2014) 075001 [[arXiv:1406.3294](#)] [[INSPIRE](#)].
- [26] LHC HIGGS CROSS SECTION WORKING GROUP, S. Heinemeyer et al., *Handbook of LHC Higgs Cross Sections: 3. Higgs Properties*, [arXiv:1307.1347](#) [[INSPIRE](#)].
- [27] J.F. Gunion, H.E. Haber, G. Kane and S. Dawson, *The Higgs Hunter’s Guide*, Addison-Wesley, Reading MA U.S.A. (1990) [*Front. Phys.* **80** (2000) 1] [[INSPIRE](#)].
- [28] J. Alwall et al., *The automated computation of tree-level and next-to-leading order differential cross sections and their matching to parton shower simulations*, *JHEP* **07** (2014) 079 [[arXiv:1405.0301](#)] [[INSPIRE](#)].
- [29] ATLAS collaboration, *Measurements of the properties of the Higgs-like boson in the four lepton decay channel with the ATLAS detector using 25 fb⁻¹ of proton-proton collision data*, *ATLAS-CONF-2013-013* (2013) [ATLAS-COM-CONF-2013-018] [[INSPIRE](#)].
- [30] ATLAS collaboration, *Search for a high-mass Higgs boson in the $H \rightarrow WW \rightarrow l\nu l\nu$ decay channel with the ATLAS detector using 21 fb⁻¹ of proton-proton collision data*, *ATLAS-CONF-2013-067* (2013) [ATLAS-COM-CONF-2013-082] [[INSPIRE](#)].
- [31] ATLAS collaboration, *Search for Scalar Diphoton Resonances in the Mass Range 65–600 GeV with the ATLAS Detector in pp Collision Data at $\sqrt{s} = 8$ TeV*, *Phys. Rev. Lett.* **113** (2014) 171801 [[arXiv:1407.6583](#)] [[INSPIRE](#)].