九州大学学術情報リポジトリ Kyushu University Institutional Repository

Prediction for progression of transverse cracking in CFRP cross-ply laminates using Monte Carlo method

Onodera, Sota

Department of Aerospace Engineering, Tohoku University: Graduate Student

Nagumo, Yoshiko

Department of Aerospace Engineering, Tohoku University: Assistant Professor

Okabe, Tomonaga

Department of Aerospace Engineering, Tohoku University: Professor

https://hdl.handle.net/2324/4403408

出版情報: Advanced Composite Materials. 26 (5), pp.477-491, 2017-05-22. Taylor and Francis

バージョン:

権利関係:



原稿種類:研究論文

題目: モンテカルロ法による CFRP 直交積層板のトランスバースクラック進展解析

(Prediction for progression of transverse cracking in CFRP cross-ply laminates using Monte Carlo method)

氏名および会員資格:

小野寺壮太・学生会員 (Sota ONODERA) *1 南雲佳子・正会員 (Yoshiko NAGUMO) *2

岡部朋永·正会員 (Tomonaga OKABE) *3

勤務先および職名:

*1~*3 東北大学大学院工学研究科航空宇宙工学専攻(980-8579 宮城県仙台市青葉区荒巻字青葉 6-6-01)

Department of Aerospace Engineering, Tohoku University

- *1 大学院生 Graduate Student
- *2 助教 Assistant Professor
- *3 教授 Professor

連絡先:南雲佳子

〒980-8579 宮城県仙台市青葉区荒巻字青葉 6-6-01 東北大学大学院 工学研究科 航空宇宙工学専攻

TEL: 022-795-6932 FAX: 022-795-6983

E-mail: nagumo@plum.mech.tohoku.ac.jp

使用ソフト: Microsoft Word 2016

使用機種: Microsoft Windows 7 Professional

Abstract

This study predicted transverse cracking progression in laminates including 90° plies. The refined stress field (RSF) model, which takes into account thermal residual strain for plies including transverse cracks is formulated, and the energy release rate associated with transverse cracking is calculated using this model. For comparison, the energy release rate based on the continuum damage mechanics (CDM) model is formulated. Next, transverse cracking progression in CFRP cross-ply laminates including 90° plies is predicted based on both stress and energy criteria using the Monte Carlo method. The results indicated that the RSF model and the CDM model proposed in this study can predict the experiment results for the relationship between transverse crack density and ply strain in 90° ply. The models presented in this paper can be applied to an arbitrary laminate including 90° plies.

(135 words)

Key words: CFRP, composite laminate, transverse crack, energy release rate, Monte Carlo method

Prediction for progression of transverse cracking in CFRP cross-ply laminates using Monte Carlo method Sota Onodera, Yoshiko Nagumo, and Tomonaga Okabe

1 INTRODUCTION

Recently, fiber-reinforced plastics (FRPs) have been used in aerospace fields where reducing weight is important. Practically, laminates made by stacking unidirectional fiber -reinforced laminae, which have predefined mechanical properties, are commonly used because unidirectional fiber -reinforced plastic is highly anisotropic. Several kinds of laminate exist. For example, cross-ply laminate has fiber angles of 0° and 90° to the loading direction, angle-ply laminate has fiber angles of $\pm \theta^{\circ}$ to the loading direction, and quasi-isotropic laminate has isotropic elastic moduli in the plane. To utilize these laminates of arbitrary lay-up configurations, it is essential to determine the fracture process of the laminate. Fiber-reinforced laminate exhibits specific fracture processes, such as transverse cracks (cracks in a direction parallel to the fiber), delamination, and breakage of the fiber. In these fracture processes, transverse cracking occurs in the earliest stage. The initiation o f transverse cracking causes stress redistribution, resulting in delamination and breakage of

the fiber due to stress concentration at the crack tip.

Therefore, it is important to clarify the mechanical behavior of laminate, including transverse cracking.

Various studies have focused on laminate including transverse cracks. $^{(1)-29)}$ Pagano et al. $^{(1)}$ formulated an analytical model that predicts the crack-initiation laminate stress of $[0/90_n/0]$ (n=1,2,3,4) cross-ply laminates assuming steady-state cracking, which is the condition under which a transverse crack propagates indefinitely under constant thermomechanical loading. They demonstrated that their analytical results can precisely predict the experiment results for crack-initiated laminate stress. They also compared their results with those obtained by McCartney $^{(15,16)}$ and Dvorak et al. $^{(14)}$

Wang et al.²⁾ calculated the energy release rate using finite-element analysis assuming that effective flaws in the 90° ply propagate into transverse cracks. Also, they simulated transverse crack density as a function of applied laminate stress using the Monte Carlo method. In their analysis, statistical distributions of effective flaw length and flaw spacing were assumed.

Gudmundson and Zang⁷⁾ proposed a sophisticated model that predicts the thermo-elastic properties of a laminate including transverse cracks. In this model, the average crack opening displacement of a transverse crack is assumed to be that of an isotropic elastic medium, and the incremental strain of the ply in which the transverse crack is initiated is calculated. They embedded this model into the classical laminate theory, so that the relationship between transverse crack density of laminates arbitrary lay-up configurations and their mechanical properties could be estimated analytically. The model formulated by Gudmundson and Zang can calculate the average stress distribution in each ply but this model cannot calculate the local stress distribution in each ply. Kobayashi et al. 8, 9) derived the energy release rate for transverse cracks using the m o d e l formulated b y Gudmundson and Zang, and predicted the formation of transverse cracks in quasi-isotropic laminates. Kobayashi et al. also found that the local stress distribution in each ply could not be determined using the model of Gudmundson and Zang, although the average stress distribution in each ply could be derived.

Okabe et al. 11) formulated the continuum damage mechanics model for predicting the stiffness reduction of composite laminates including transverse cracks as a function of crack density. This model can predict the stiffness of laminates damaged due to transverse cracks (or

surface crack) based on only the mechanical properties of a ply and the lay-up configurations. Furthermore, this model can precisely predict the FEA results and experiment results for the elastic moduli of laminates of arbitrary lay-up configuration (e.g., cross-ply, angle-ply, and quasi-isotropic) including transverse cracks. This model can calculate the local stress distribution in a ply including transverse cracking as a function of crack density.

The present study formulates at model for progression of transverse cracking in laminates including 90° ply. First, the local stress distribution in a ply including transverse cracking is formulated using the refined stress field (RSF) model, which extends the model formulated by Okabe et al. 11) to take into account residual thermal strain. Next, the energy release rate associated with transverse cracking is calculated using the RSF model. For comparison, the energy release rate based on the continuum damage mechanics (CDM) model is formulated. Furthermore, transverse cracking progression in CFRP cross-ply laminates including 90° plies is predicted based on both stress and energy criteria using the Monte Carlo method. Finally, the crackinitiated laminate stress and the relationship between transverse crack density and ply strain in 90° ply are calculated using the RSF or the CDM model. The validity

of these models is estimated by comparing the results obtained from these models to the experimental and analytical results in previous studies. 1), 10) The results for cross-ply laminates are calculated in this study; however, the models presented here can be applied to an arbitrary laminate including 90° plies.

2 MECHANICAL MODEL

2.1 Local Stress Distribution in a Ply Including Transverse Cracking

Local stress distribution in a ply including transverse cracking was formulated using the RSF model, which extends the model formulated by Okabe et al. 11) to take into account residual thermal strain. When formulating the RSF model, it was assumed that the ply is thin, and that the damage was due mainly to transverse cracks. Based on these assumptions, the bending of the ply and damage due to delamination were ignored. Figure 1 (a) illustrates a ply including transverse cracks. The 1-axis is parallel to the fiber, and the 2- and 3-axes are normal to the fiber. Figure 1 (b) illustrates a representative volume element (RVE) including transverse cracking on both sides, which is part of the ply in Fig. 1 (a). It should be noted that the coordinate system differs from that in Fig. 1 (a). The

coordinates x, y in Fig. 1 (b) correspond to the 3-, 2-axes in Fig. 1 (a), respectively. The transverse crack is assumed to have a tunnel-like crack surface that is symmetrical about the y-axis. The transverse crack does not propagate into the neighboring ply. The crack spacing is 2l and the thickness of the ply is $t_{\rm ply} = 2t$. The ply strain ε_y^p along the y-axis (or 2-axis) is applied to the RVE (or ply). Due to the symmetric RVE, the region of definition is limited to $0 \le x \le t$ and $0 \le y \le l$.

A plane strain in the isotropic plane (x-y plane) in Fig. 1 (b) is assumed. The displacements in directions x and y are defined as u and v. Considering thermal residual strain, the strain-stress relationship is expressed as

$$\varepsilon_x^M = \frac{\partial u^M}{\partial x} = \varepsilon_x - \alpha_2 \Delta T = C_1 \sigma_x - C_2 \sigma_y, \tag{1}$$

$$\varepsilon_y^M = \frac{\partial v^M}{\partial v} = \varepsilon_y - \alpha_2 \Delta T = C_1 \sigma_y - C_2 \sigma_x, \tag{2}$$

$$\gamma_{xy} = \frac{\sigma_{xy}}{G_{23}} = \frac{\partial v^M}{\partial x} + \frac{\partial u^M}{\partial y} \approx \frac{\partial v^M}{\partial x}.$$
(3)

Constants \mathcal{C}_1 and \mathcal{C}_2 are defined as

$$C_1 = \frac{1 - \nu_{21}\nu_{12}}{E_2}, C_2 = \frac{\nu_{23} + \nu_{21}\nu_{12}}{E_2}, \tag{4}$$

where σ_x and σ_y are stresses along the x- and y-axes, ε_x and ε_y are strains considering the influence of thermal residual strain along the x- and y-axes, ε_x^M and ε_y^M are strains not including the influence of thermal residual

x- and y-axes, u^M and v^M t h e strains along displacements not including the influence of thermal residual strains along the x- and y-axes, γ_{xy} is the engineering shear strain in the x-y plane, σ_{xy} is the shear stress in the x-y plane, and ΔT is temperature change T T_{sf} from the stress free temperature T_{sf} to the testing temperature T. E_1 is the longitudinal Young's modulus along 1-axis, E_2 is the transverse Young's modulus along the 2-axis, G_{23} is the out-of-plane shear modulus in plane 2-3, v_{12} is in-plane Poisson's ratio in plane 1-2, v_{23} is outof-plane Poisson's ratio in plane 2-3, and α_2 is the transverse thermal expansion coefficient along the 2-axis. For engineering shear strain γ_{xy} in Eq. (3), the gradient in the y direction of displacement u is assumed to be quite small. The relationship between $arepsilon_x^M$ and $arepsilon_y^M$ is expressed as

$$\varepsilon_x^M = a \varepsilon_y^M, \tag{5}$$

where a is a proportionality constant defined to satisfy the equilibrium equation of stress. From Eqs. (1), (2), and (5), σ_x and σ_y are defined as

$$\sigma_{x} = \frac{aC_1 + C_2}{C_1^2 - C_2^2} \frac{\partial v^M}{\partial y},\tag{6}$$

$$\sigma_{y} = \frac{C_1 + aC_2}{C_1^2 - C_2^2} \frac{\partial v^M}{\partial y}.$$
 (7)

Additionally, using Eq. (3), σ_{xy} is expressed by

$$\sigma_{xy} = G_{23} \frac{\partial v^M}{\partial x}.$$
 (8)

The equilibrium equation of stress is expressed as

$$\frac{\partial \sigma_x}{\partial x} + \frac{\partial \sigma_{xy}}{\partial y} = 0, \tag{9}$$

$$\frac{\partial \sigma_x}{\partial x} + \frac{\partial \sigma_{xy}}{\partial y} = 0, \tag{9}$$

$$\frac{\partial \sigma_{xy}}{\partial x} + \frac{\partial \sigma_y}{\partial y} = 0. \tag{10}$$

 $\partial^2 v^M / \partial x \partial y \neq 0$ is postulated and Eqs. (6)-(8) substituted into Eqs. (9) and (10), the proportionality constant a and the Laplace equation for displacement v^M are obtained as

$$a = -\{C_2 + G_{23}(C_1^2 - C_2^2)\}/C_1,\tag{11}$$

$$\frac{\partial^2 v^M}{\partial x^2} + \lambda^2 \frac{\partial^2 v^M}{\partial y^2} = 0; \lambda = \sqrt{\frac{C_1 + aC_2}{G_{23}(C_1^2 - C_2^2)}}.$$
 (12)

To determine displacement v^{M} , the boundary conditions of the Laplace equation in Eq. (12) are given by

$$v^M = 0$$
 on $y = 0$, (13)

$$\frac{\partial v^M}{\partial y} = 0 \quad \text{on} \quad y = l, \tag{14}$$

$$\frac{\partial v^M}{\partial x} = 0 \quad \text{on} \quad x = 0, \tag{15}$$

$$v^{M} = \{ \varepsilon_{y}^{p} + (\alpha_{c} - \alpha_{2}) \Delta T \} y \text{ on } x = t,$$
 (16)

where α_{c} is the thermal expansion coefficient of laminate along the 0° direction. At y = 0, no displacement v is considered in Eq. (13). Equation (14) was determined from Eq. (7) considering stress $\sigma_y = 0$ at the crack surface (y=l). Equation (15) was defined based on Eq. (8), considering shear stress $\sigma_{xy} = 0$ on the center line of the ply (x=0). Additionally, the displacement distribution presented as Eq. (16) in the interface was assumed according to the previous study.²⁵⁾ Therefore, the neighboring ply is thought to be deformed uniformly by ε_y^p , regardless of the transverse cracks. Using separation of variables and assigning the boundary conditions presented above to Eq. (12), the solution v^M of the Laplace equation that satisfies the boundary conditions in Eqs. (13) to (16) can be expressed as

$$v^{M} = \frac{8l}{\pi^{2}} \left(\sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{(2n-1)^{2}} \frac{\cosh[(2n-1)\pi\lambda x/(2l)]}{\cosh[(2n-1)\pi\lambda t/(2l)]} \sin\left[\frac{(2n-1)\pi y}{2l}\right] \right) (\varepsilon_{y}^{p} + \varepsilon_{th}), \tag{17}$$

where ε_{th} denotes $(\alpha_c - \alpha_2)\Delta T$. When Eq. (17) is substituted into Eqs. (6)-(8), the local stress distribution in RVE is derived as

$$\sigma_x(x,y) = \frac{4}{\pi} \frac{aC_1 + C_2}{C_1^2 - C_2^2} f(x,y) \left(\varepsilon_y^p + \varepsilon_{th}\right),\tag{18}$$

$$\sigma_{y}(x,y) = \frac{4}{\pi} \frac{C_1 + aC_2}{C_1^2 - C_2^2} f(x,y) \left(\varepsilon_y^p + \varepsilon_{th}\right),\tag{19}$$

$$\sigma_{xy}(x,y) = -\frac{4}{\pi} G_{23} \lambda g(x,y) \left(\varepsilon_y^p + \varepsilon_{th} \right). \tag{20}$$

In the above equations, the functions f(x,y) and g(x,y) are expressed by

$$f(x,y) = \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{(2n-1)^2} \frac{\cosh[(2n-1)\pi\lambda x/(2l)]}{\cosh[(2n-1)\pi\lambda t/(2l)]} \cos\left[\frac{(2n-1)\pi y}{2l}\right],\tag{21}$$

$$g(x,y) = \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{(2n-1)^2} \frac{\sinh[(2n-1)\pi\lambda x/(2l)]}{\cosh[(2n-1)\pi\lambda t/(2l)]} \sin\left[\frac{(2n-1)\pi y}{2l}\right]. \tag{22}$$

The compatibility condition for the strain is not satisfied because of the assumption of Eqs. (3) and (5). Therefore, the proposed stress distribution is analytical solution, but they are not exact solution.

2.2 Formulation of Energy Release Rate Using the RSF Model

Here, only transverse cracking in 90° ply under uniaxial tension is considered. The energy release rate associated with transverse cracking in 90° ply is formulated using the RSF model. Considering that new transverse cracking propagates between two pre-existing transverse cracks (Fig.2), the energy release rate G associated with transverse cracking is defined as

$$G = \frac{U(l) - [U(l_1) + U(l_2)]}{2t},\tag{23}$$

where U(l) is a strain energy stored in the RVE of crack spacing 2l (Fig. 1 (b)). The strain energy U(l) is expressed as

$$U(l) = \frac{1}{2} \int_{V} (\sigma_{x} \varepsilon_{x} + \sigma_{y} \varepsilon_{y} + \sigma_{xy} \gamma_{xy}) dV, \qquad (24)$$

where the integration area V is defined as the inside of the RVE (Fig. 1 (b)). Using Eqs. (1)-(3) and (18)-(20), Eq. (24) can be described as

$$U(l) = \int_0^t \left\{ \int_0^l \left(A\{f(x,y)\}^2 \left(\varepsilon_y^p + \varepsilon_{th} \right)^2 + B\{g(x,y)\}^2 \left(\varepsilon_y^p + \varepsilon_{th} \right)^2 + Cf(x,y) \left(\varepsilon_y^p + \varepsilon_{th} \right) \right) dy \right\} dx.$$

$$(25)$$

The constants A, B, and C are calculated as

$$A = \frac{32(1+a^2)C_1 + 2aC_2}{C_1^2 - C_2^2},$$
 (26)

$$B = \frac{32}{\pi^2} G_{23} \lambda^2, \tag{27}$$

$$C = \frac{8}{\pi} \alpha_2 \Delta T \frac{1+a}{C_1 - C_2}.$$
 (28)

Substituting Eqs. (25)-(28) into Eq. (23), the energy release rate G can be calculated using the RSF model. Equation (25) is calculated numerically using Gauss-Legendre quadrature because it is difficult to solve this equation analytically.

2.3 Formulation of Energy Release Rate Using the CDM Model

For comparison with the energy release rate using RSF model, the energy release rate based on the CDM model is also formulated. In the CDM model, strain energy U(l) stored in the RVE of crack spacing 2l (Fig. 1 (b)) is expressed as

$$U(l) = \frac{1}{2} \cdot E_2 (1 - d_2(l)) (\varepsilon_y^p)^2 \cdot 4tl, \tag{29}$$

where $d_2(l)$ is the damage variable in the 2-axis (direction normal to the fiber) as shown in Fig. 1 (a). According to

Okabe et al.¹¹⁾, by applying the Gudmundson-Zang model⁷⁾ the damage variable $d_2(l)$ is defined as

$$d_2(l) = \frac{\pi}{2} \frac{tE_1(1 - \nu_{12}\nu_{21})}{E_1 - E_2\nu_{12}^2} \frac{1}{l} \sum_{j=1}^{10} \frac{a_j}{(1 + t/l)^j},$$
 (30)

where the constant sequence a_j is given in Table 1. Substituting Eqs. (29) and (30) into (23), the energy release rate G can be calculated using the CDM model.

3 ANALYSIS OF THE TRANSVERSE CRACKING PROGRESSION USING THE MONTE CARLO METHOD

Transverse cracking progression based on both stress and energy criteria was predicted using the Monte Carlo method. First, the 90° ply of length L is discretized into n elements (Fig. 3). The static transverse strength of each element ply is affected by local defects. It is assumed that transverse strength F_i of element i (i=1,2,..., n) follows Weibull distribution R_i , which reflects the random material property to resist crack formation:

$$R_i = 1 - e^{-\frac{\Delta L}{L_0} \left(\frac{F_i}{\sigma_0}\right)^m},\tag{31}$$

where L_0 is the gauge length, ΔL is the element spacing, σ_0 is the transverse strength of ply, and m is the shape parameter. R_i is determined by assigning pseudo-random numbers to each element, and the strength F_i of each element is calculated using Eq. (31). The pseudo-random

numbers are generated by Mathematica 9.0.³⁰⁾ Shape parameter m is a fitting parameter in this study. Next, it is assumed that a micro crack occur at the ply strain ε_y^p in each element satisfying the following stress criteria

$$\sigma_{\mathcal{V}}(0, y) > F_i, \tag{32}$$

where $\sigma_y(0,y)$ is stress in the y-direction at x=0 for the RSF model, given by Eq. (19).

It is then assumed that one of the micro cracks that occurred at the ply strain ε_y^p propagates into a transverse crack. The element with the highest energy release rate associated with transverse cracking among the elements with micro cracks is searched. When this element satisfies the energy criterion

$$G > G_{c_{\ell}} \tag{33}$$

the micro crack in this element propagates into a transverse crack. Here, G_c is the critical energy release rate. Energy release rate G is calculated using the expression formulated by the RSF model or the CDM model. The energy release rate G is calculated assuming that there are virtual transverse cracks at either end of the 90° ply of length L. The transverse crack density is calculated by dividing the number of transverse cracks with the exception of the virtual transverse cracks by the length L.

The progression of transverse cracking in laminates

including 90° plies is analyzed according to the stress and energy criteria. Figure 4 presents the analysis flowchart.

4 RESULTS AND DISCUSSION

4.1 Crack-Initiated Laminate Stress

Before predicting transverse cracking progression, the crack-initiated laminate stress is obtained using the formulated mechanical model and making a comparison with the experimental and analytical results obtained by Pagano¹⁾ et al. The laminate material is IM7/5250-4. Table 1 lists its properties. The out-of-plane shear modulus G_{23} is estimated using the following expression

$$G_{23} = \frac{E_2}{2(1+\nu_{23})}. (34)$$

Cross-ply $[0/90_n/0]$ (n=1,2,3,4) laminates with transverse cracks in the 90° plies are considered. In this analysis, it is assumed that the cracks are equally spaced $(l_1=l_2=l/2)$, and transverse crack density is defined as $\rho=1/l$. The crack-initiated laminate stress is determined using the following procedure.

(1) The energy release rate G is calculated as a function of transverse crack density ρ at ply strain ε_y^p using the RSF model or the CDM model. The transverse crack density ρ_c and ply strain $(\varepsilon_y^p)_c$ are evaluated when the

maximum value of the energy release rate is equal to the critical energy release rate G_c .

(2) The effective compliance \overline{c} of the laminate at crack density ρ_c is obtained from the model to predict the stiffness reduction of laminates including transverse cracks, formulated by Okabe $et\ al.$, $^{11)}$ and the stress σ_c loaded in the laminate along the 0° direction at $(\varepsilon_y^p)_c$ is obtained using the constitutive law. Stress σ_c is defined as crack-initiated laminate stress.

Figure 5 plots the relationship between the crack-initiated laminate stress σ_c calculated using the above method and the thickness of 90° plies. Figure 5 also presents the experiment and analytical results obtained by Pagano¹⁾ when the transverse crack penetrated the width direction of the specimen. The crack-initiated laminate stress obtained using the RSF model is the largest in the analytical models. The reason is that change in the strain distribution at the interfaces of ply with transverse cracks due to the occurrence of transverse cracks in the RSF model is not considered because the strain at the interfaces of the ply with transverse cracks defined by Eq. (16) is constant. Comparison of the RSF model with the experiment results indicates that the thicker the 90° plies, the closer these

values are. The crack-initiated laminate stress obtained using the CDM model is the smallest in the analytical models. It can be confirmed that the RSF model gives the upper limit of crack-initiated laminate stress, and the CDM model gives the lower limit.

4.2 Prediction of Transverse Cracking Progression

Transverse cracking progression is predicted using the mechanical model formulated in section 2 and the analysis method described in section 3, and compared with the experiment results of Ogihara et al. 10) The cross-ply laminates of T800H/3900-2 and G40-800/5260 are analyzed. The material properties of T800H/3900-2 and G40-800/5260 used in the calculation are listed in Table 3. The out-ofplane shear modulus G_{23} is estimated using Eq. (34). The shape parameter m is determined by fitting the relationship between transverse crack density and ply strain for the $[0/90]_s$ laminates obtained by Ogihara et al. 10 and the prediction result for the [0/90]_s laminates obtained using the RSF model. The length L of the analysis region is 50mm, and there are 8000 elements n. The element spacing ΔL is calculated by dividing the length L by the number of element n, and ΔL is 6.25×10^{-3} mm. Transverse crack density is calculated by dividing the number of transverse

cracks by the length L.

First, the analysis results for the [0/90]_s T800H/3900-2 laminate are considered. Figure 6 presents the transverse ply for distribution in the 90° t h e $[0/90]_{s}$ crack T800H/3900-2 laminate along the longitudinal length L using the RSF model under $\varepsilon_{\nu}^{p}=1.08\%$ or $\varepsilon_{\nu}^{p}=1.2\%$. In Fig. 6 (a), there are 24 micro cracks and 12 transverse cracks. In Fig. 6 (b), there are 33 micro cracks and 28 transverse cracks. The transverse cracks are not equally spaced, and it is confirmed that the transverse cracks are distributed almost evenly as the ply strain increases. Figure 7 plots the relationship between the transverse crack density in the 90° plies and the ply strain calculated using the Monte Carlo method described in section 3. For comparison, the experiment results of Ogihara et al. 10) are presented. The results obtained by the RSF model and the CDM model are in good agreement with the experimental results. The first cracking strain obtained using the RSF model exceeds that obtained using the CDM model. This result is consistent with the result in Fig. 7. The behavior of the RSF model and that of the CDM model are almost identical where the transverse crack density is high. Transverse crack density is defined as $\rho = 1/l$ assuming the cracks are equally spaced $(l_1=l_2=l/2)$. The relationship between energy release rate $\it G$ and transverse crack density ρ at ply strain $\varepsilon_y^p=1.0\%$ is plotted in Fig. 8. The vales of the energy release rate of the RSF model and the CDM model are almost identical when ρ is large. Therefore, the behaviors of the RSF model and the CDM model indicated in Fig. 7 almost agree when the transverse crack density is high.

Next, the analysis results for the G40-800/5260 laminate are presented. Figures 9 and 10 plot the relationship between transverse crack density and ply strain calculated using the Monte Carlo method for the $[0/90]_s$ and $[0/90_2]_s$ laminates. For comparison, the experiment results of Ogihara et al. 10) are presented in Figs. 9 and 10. As in Fig. 7, the RSF model and the CDM model behaviors are almost identical in the range where transverse crack density is high. In Fig. 9, the RSF model gives the upper limit of crack-initiated laminate stress, and the CDM model gives the lower limit. The experiment results of the crackinitiated laminate stress exceed the values of the RSF model and the CDM model in Fig. 10. However, the results obtained using the RSF model and the CDM model agree with the experiment results, regardless of the thickness of the 90° plies, confirming the validity of this analysis. Furthermore, when comparing Figs. 9 and 10, the results of the first cracking strain obtained using the RSF

model become closer to the experiment values as the 90° plies become thicker. This is consistent with the result in Fig. 5.

This study calculates the results for cross-ply laminates. However, the presented models can be applied to an arbitrary laminate including 90° plies. Furthermore, by combining the present model and the model for predicting stiffness reduction of laminates including transverse cracks formulated by Okabe *et al.* ¹¹⁾, it is possible to predict the stress-strain behavior of the laminate.

5 CONCLUSIONS

This study investigated transverse cracking progression in laminates including 90° plies. The RSF model, which takes into account the thermal residual strain for ply including transverse cracks, is formulated, and the energy release associated with transverse cracking is rate calculated using this model. For comparison, the energy release rate based on the CDM model is formulated. Next, transverse cracking progression in CFRPcross-ply laminates including 90° plies based on both stress and energy criteria is predicted using the Monte Carlo method. The results indicated that the RSF model and CDM model proposed in this study can predict the experiment results

for the relationship between transverse crack density and ply strain in 90° ply obtained by Ogihara $et\ al.^{10}$ In this study, the results for cross-ply laminates are calculated; however, the presented models can be applied to an arbitrary laminate including 90° ply.

ACKNOWLEDGEMENTS

T. O. appreciates the support of the Cross-ministerial Strategic Innovation Promotion Program and the Ministry of Education, Cluture, Sports, Science and Technology of Japan under Grant-in-Aid Scientific Research (C) No. 15K06597. This work was partly supported by Toray Industries, Inc. and the New Energy and Industrial Technology Development Organization (NEDO) No. 15102315-0, P15006, and 141001611-d.

REFERENCES

- N.J. Pagano, G.A. Schoeppner & R. Kim: Comp. Sci.
 Tech., 58,11 (1998), 1811-1825.
- A. S. D. Wang, P. C. Chou & S.C. Lei : J. Comp. Mater.,
 18, 3 (1984), 239-254.
- 3) J.M. Berthelot & J.-F. Le Corre: Comp. Sci. Tech., 60, 14 (2000), 2659-2669.
- 4) J.M. Berthelot, Appl. Mech. Rev., 56, 1 (2003), 111-147.

- 5) J.M. Berthelot, P. Leblond, A. El Mahi & J. -F. Le Corre: Comp. Part A, 27, 10 (1996), 989-1001.
- 6) N. Takeda, & S. Ogihara: Comp. Sci. Tech., **52**, 2 (1994), 183-195.
- 7) P. Gudmundson & W. Zang: Int. J. Solid Structures, 30, 23 (1993), 3211-3231.
- 8) S. Kobayashi, S. Ogihara, & N. Takeda: Adv. Comp.

 Mater., 9, 4 (2000), 363-375.
- 9) S. Kobayashi, S. Ogihara & N. Takeda: Adv. Comp.

 Mater., 9, 3 (2000), 241-251.
- 10) S. Ogihara, S. Kobayashi, N. Takeda, A. Kobayashi: Comp. Sci. Tech., 61, 8 (2001), 1049-1055.
- 11) T. Okabe, S. Onodera, Y. Kumagai & Y. Nagumo: Int.J. Damage mech. (in press).
- 12) G. J. Dvorak, N. Laws & M. Hejazi: J. Comp. Mater.,19, 3 (1985), 216-234.
- 13) G. J. Dvorak & N. Laws: J. Comp. Mater., 21, 4 (1987), 309-329.
- 14) G. J. Dvorak & N. Laws: Eng. Frac. Mech., 25, 5-6(1986), 763-770.
- 15) McCartney, L. N. : J. Mech. Phys. Solids, 40, 1 (1992), 27-68.
- 16) McCartney, L. N.: Stress transfer mechanics for ply cracks in general symmetric laminates, NPI, Report

- CMMT(A)50, National Physics Laboratory, Teddington, UK, 1996.
- 17) J. A. Nairn: J. Comp. Mater., 23, 11 (1989), 1106-1129.
- 18) J. A. Nairn & S. Hu: Int. J. Fract., 57, 1 (1992), 1-24.
- 19) Y. Huang, J. Varna & R. Talreja: Comp. Sci. Tech.,95, (2014), 100-106.
- S. Ogihara, N. Takeda, S. Kobayashi & A. Kobayashi :
 J. Soc. Mat. Sci., Japan, 47, 1 (1998), 68-72.
- 21) N. Takeda, H. Niizuma, S. Ogihara & A. Kobayashi:

 Materials system, Japan, 14, (1995), 73-78.
- 22) E. V. Iarve, M.R. Gurvich, D.H. Mollenhauer, C.A. Rose & C.G. Dávila: *Int. J. Numer. Meth. Eng.*, 88, 8 (2011), 749-773.
- 23) E.V. Iarve, S. Chellapilla & D. H. Mollenhauer: ICCM 15 Conference Proceedings, Durbon, S. Africa June 27-July 2, (2005).
- 24) M. R. Gurvich: Comp. Sci. Tech., **59**, 11 (1999), 1701-1711.
- 25) J. W. Lee, D. H. Allen and C. E. Harris: J. Comp. Mater., 23, 12 (1989), 1273-1291.
- 26) S. Murakami : Continuum Damage Mechanics: A
 Continuum Mechanics Approach to the Analysis of

Damage and Fracture. Dordrecht Heidelberg London New York: Springer, (2012).

- 27) Y. Kanagawa, S. Murakami, Y. Liu, Q. Bai & K. Tanaka: J. Soc. Mat. Sci., Japan 45, 2 (1996), 206-211.
- 28) Z.C. Xia, R.R. Carr & J.W. Hutchinson: Acta Metall.

 Mater., 41, 8 (1993), 2365-2376.
- Z. Hashin: Eng. Frac. Mech., 25, 5-6 (1986), 771-778.
- 30) Wolfram Research, Inc., Mathematica, Version 9.0, Champaign, IL (2012).

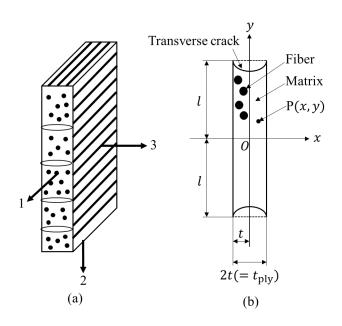


Fig. 1. Schematic of transverse crack model. (a) Ply including transverse crack. (b) Representative volume element.

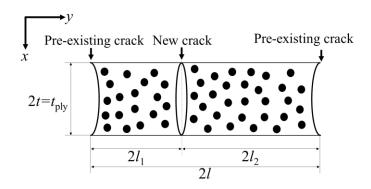


Fig. 2. Formation of a new crack between two preexisting cracks.

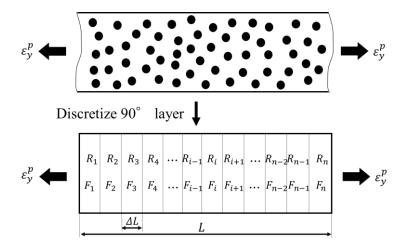


Fig. 3. Discretization of 90° plies for distributed static transverse strength.

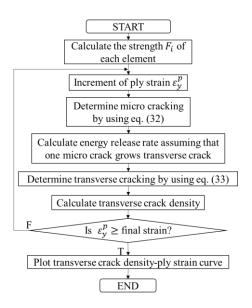


Fig. 4. Flow chart of the procedure implemented for analyzing the progress of transverse cracking.

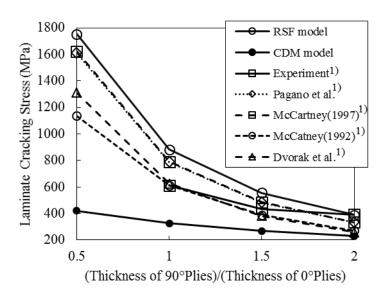


Fig. 5. Laminate cracking stress as a function of thickness of 90° plies / thickness of 0° plies.

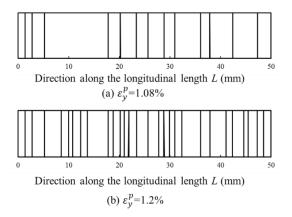


Fig. 6. Prediction of transverse crack distribution in the 90° ply for $[0/90]_s$ T800H/3900-2 laminate along the longitudinal length L using RSF model. Each vertical line represents one transverse crack. (a) Under $\varepsilon_y^p = 1.08\%$, (b) Under $\varepsilon_y^p = 1.2\%$.

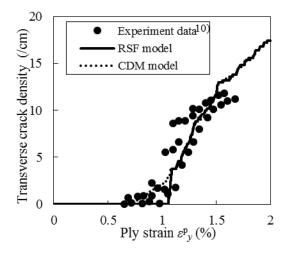


Fig. 7. Transverse crack density in the 90° ply for $[0/90]_s\ T800H/3900-2\ laminate\ as\ a\ function\ of\ ply\ strain.$

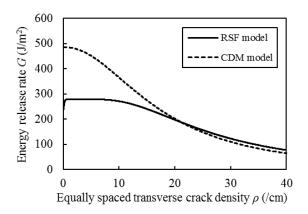


Fig. 8. Energy release rate associated with transverse cracking in 90° ply as a function of equally spaced transverse crack density for $[0/90]_s$ T800H/3900-2 laminate $(\varepsilon_y^p = 1.0\%)$.

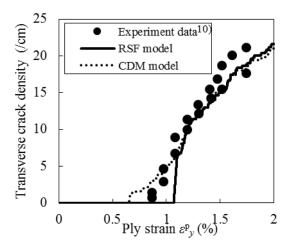


Fig. 9. Transverse crack density in 90° ply for $[0/90]_s$ G40-800/5260 laminate as a function of ply strain.

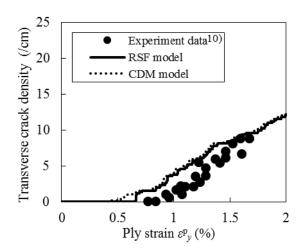


Fig. 10. Transverse crack density in 90° ply for $[0/90_2]_s$ G40-800/5260 laminate as a function of ply strain.

Table 1. Numerical parameter a_j used in Eq. $(30)^{7}$.

j	a_j	j	a_j
1	0.63666	6	-33.09444
2	0.51806	7	74.32002
3	0.51695	8	-103.06411
4	-1.04897	9	73.60337
5	8.95572	10	-20.34326

Table 2. Material properties of IM7/5250-4.1)

Longitudinal Young's Modulus E ₁ (GPa)	165.475
Transverse Young's Modulus E ₂ (GPa)	10.342
In-Plane Poisson's Ratio $ u_{12}$	0.31
Out-of-Plane Poisson's Ratio $ u_{23}$	0.56
In-Plane Shear Modulus G ₁₂ (GPa)	5.7922
Longitudinal Thermal Expansion	0.45
Coefficient α_1 (10 ⁻⁶ / $^{\circ}$ C)	
Transverse Thermal Expansion	24.66
Coefficient α_2 (10 ⁻⁶ / $^{\circ}$ C)	
Stress-Free Temperature T_{sf} (°C)	180
Ambient Testing Temperature T ($^{\circ}\mathbb{C}$)	2 4
Ply thickness t_0 (mm)	0.127
Critical Energy Release Rate G_c (J/m 2)	2 2 5

Table 3. Material properties of T800H/3900-2 and G40- $800/5260.^{8,9})$

	T800H/3900-2	G 4 0 - 8 0 0 / 5 2 6 0
Longitudinal Young's Modulus E_1 (GPa)	143.0	152.0
Transverse Young's Modulus E ₂ (GPa)	7.99	10.0
In-Plane Poisson's Ratio v_{12}	0.35	0.33
Out-of-Plane Poisson's Ratio $ u_{23}$	0.49	0.49
In-Plane Shear Modulus G_{12} (GPa)	3.96	6.94
Longitudinal Thermal Expansion	-1.5	-0.6
Coefficient α_1 (10 ⁻⁶ /°C)		
Transverse Thermal Expansion	34.3	36.0
Coefficient α_2 (10 ⁻⁶ / $^{\circ}$ C)		
Stress-Free Temperature T_{sf} (°C)	165	195
Ambient Testing Temperature T ($^{\circ}\mathbb{C}$)	2 5	2 5
Ply thickness $t_{\rm ply}$ (mm)	0.18	0.14
Gauge length L_0 (mm)	8 0	8 0
Critical Average Stress σ_0 (MPa)	7 2	7 5
Distribution Shape Constant m	10.5	7.5
Critical Energy Release Rate G_c (J/m 2)	3 0 0	3 3 0