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# Variational problem of anisotropic energy and its discretization 

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## Doctoral Dissertation

# Variational problem of anisotropic energy and its discretization 

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## Preface

The classical isoperimetric problem asks the minimizer of the area functional among all closed surfaces enclosing the same volume. The answer is a sphere and this gives a mathematical model of a soap bubble since its surface tension is isotropic, i.e., independent of directions. On the other hand, the shape of a small crystal is the minimizer of the total surface free energy associated to the media subject to a volume constraint condition. In order to model these objects, J. W. Gibbs (1839-1903) introduced an anisotropic surface energy which is the integral of an energy density that depends on the surface normal ([43], [44]). There exists the unique minimizer of the anisotropic energy among all hypersurfaces enclosing the same volume and the minimizer is called the Wulff shape [40]. Moreover, equilibrium hypersurfaces under a volume-constraint condition are characterized by constant anisotropic mean curvature (CAMC) hypersurfaces [24].

In Chapter 1, we prove non-uniqueness results of CAMC hypersurfaces by giving examples (Theorem 1.1.1, Theorem 1.1.2), i.e., there exists an energy such that there exists a CAMC hypersurface in $\mathbb{R}^{n+1}$ which is not the Wulff shape. Our results show that Alexandrov-type theorem and Hopf-type theorem for CAMC hypersurfaces cannot hold in general if we don't assume the convexity of the energy density function. Moreover, we will show that there exist non-trivial self-similar shrinking solutions for the anisotropic mean curvature flow as an application (Theorem 1.1.4).

In Chapter 2, we will develop a theory of discrete curves and surfaces based on the (anisotropic) isoperimetric problem. Here, we mainly focus on discrete objects themselves, not approximation or convergence problems. Such a research field called discrete differential geometry is an active research field because of the interaction among the theory, the algorithm and the visualization on the computer, e.g. [10], [21], [34], [36], [37], ,[38]. Surprisingly, a basic differential geometric treatment of discrete planar curves by using variational methods is not well understood. Therefore we first treat the planar curve case, and our basic tool here is a "discrete curvature vector" on a vertex extracted from the first variation formula. We show that this vector can be used effectively for an unifed interpretation well-known discrete curvature notions, the second variation and stability criteria. Since there is no universal definition of the discrete curvature and our main interest is CAMC curves and surfaces, we define discrete CAMC curves and surfaces without defining the curvature itself. Our definition is inspired by the definition in [37] and we visualize them on the software JavaView.

## Chapter 1

## Non-uniqueness of closed embedded non-smooth hypersurfaces with constant anisotropic mean curvature


#### Abstract

In this chapter, we will show that there exist an anisotropic energy density function such that there exist closed embedded equilibrium surfaces with genus zero in $\mathbb{R}^{3}$ each of which is not (any homothety and translation of) the Wulff shape. We also give nontrivial self-similar shrinking solutions of anisotropic mean curvature flow. These results are generalized to hypersurfaces in $\mathbb{R}^{n+1}$. The contents of this chapter is contained in [20].


### 1.1 Introduction

In this chapter, we consider the geometry of equilibrium surfaces for the isoperimetric problem with respect to the anisotropic energy. Let $\gamma: S^{n} \rightarrow \mathbb{R}_{>0}$ be a positivevalued continuous function on the unit sphere $S^{n}=\left\{\nu \in \mathbb{R}^{n+1} \mid\|\nu\|=1\right\}$ in the $(n+1)$-dimensional Euclidean space $\mathbb{R}^{n+1}$ and this function gives an energy density. Sometimes we assume that $\gamma$ is defined on an open set in $S^{n}$. When we study variational problems of such an energy, it is natural to consider not only smooth surfaces but also surfaces with edges since the energy minimizer may have some singularities as shown in the figure 1.1. And so let $M$ be an $n$-dimensional oriented connected compact $C^{\infty}$ manifold and $X: M=\cup_{i} M_{i} \rightarrow \mathbb{R}^{n+1}$ be a piecewise- $C^{2}$ immersion (the definition of piecewise- $C^{2}$ weak immersion will be given in the section 1.2). Let $\mathcal{S}(X)$ be the set of singularities of $X$ and $\nu=\nu_{X}: M \backslash \mathcal{S}(X) \rightarrow S^{n}$ be the unit normal vector field along $X$. We define the anisotropic energy $\mathscr{F}_{\gamma}(X)$ of $X$ by

$$
\mathscr{F}_{\gamma}(X):=\sum_{i} \int_{M_{i}} \gamma(\nu) d A=\int_{M} \gamma(\nu) d A,
$$



Figure 1.1: Figures of the energy minimizer ( $n=2$ case)
where $d A$ is the $n$-dimensional volume form of $M$ induced by $X$. If the integrand $\gamma$ is defined on an open set $\Omega \subset S^{n}$, we only consider immersions that the image of the unit normal is contained in $\Omega$. In this case, the immersion is said to be compatible with the energy density $\gamma$.

If $\gamma \equiv 1, \mathscr{F}_{\gamma}(X)$ is the usual $n$-dimensional volume of the hypersurface $X$. For any positive number $V>0$, there exists a unique (up to translations in $\mathbb{R}^{n+1}$ ) minimizer $W(V)$ of the energy $\mathscr{F}_{\gamma}$ among all closed hypersurfaces in $\mathbb{R}^{n+1}$ enclosing the same ( $n+$ 1)-dimensional volume $V([40])$. Here a closed hypersurface means that the boundary (having tangent space almost everywhere) of a set of positive Lebesgue measure. The minimizer $W\left(V_{0}\right)$ for $V_{0}=(n+1)^{-1} \int_{S^{n}} \gamma(\nu) d S^{n}$ is called the Wulff shape for $\gamma$ and we denote it by $W_{\gamma}$. The Wulff shape $W_{\gamma}$ is the boundary of the convex set

$$
\bigcap_{\nu \in S^{n}}\left\{x \in \mathbb{R}^{n+1} \mid\langle x, \nu\rangle \leq \gamma(\nu)\right\}
$$

When $\gamma \equiv 1, W_{\gamma}$ is the unit sphere.
A piecewise- $C^{2}$ weak immersion $X: M \rightarrow \mathbb{R}^{n+1}$ is a critical point of $\mathcal{F}_{\gamma}$ for all variations that preserve the enclosed $(n+1)$-dimensional volume (we will call such a variation a volume-preserving variation) if and only if the anisotropic mean curvature of $X$ is constant and $X$ has a certain condition on its singular points (cf. [24]). Here the anisotropic mean curvature $\Lambda$ of $X$ is defined at each regular point of $X$ as $\Lambda:=$ $(1 / n)\left(-\operatorname{div}_{M} D \gamma+n H \gamma\right)$, where $D \gamma$ is the gradient of $\gamma$ on $S^{n}$ and $H$ is the mean curvature of $X$ (cf. [39], [25], [24]). We call such $X$ a CAMC (constant anisotropic mean curvature) hypersurface. When $\gamma \equiv 1, \Lambda$ coincides with $H$ and CAMC hypersurfaces are CMC (constant mean curvature) hypersurfaces (of $C^{\omega}$ class).

A natural question here is
"is any closed CAMC hypersurface $X$ homothetic to the Wulff shape?"
The answer is not affirmative even in the case where $\gamma \equiv 1$ ([42], [22], [23]). However, it is expected that, if $X$ satisfies one of the following additional conditions (I)-(III), the image of $X$ is a homothety of the Wulff shape :
(I) $n=2$ and the genus $g(M)$ of $M$ is 0 . (II) $X$ is an embedding. (III) $X$ is stable. If we assume that the Wulff shape $W_{\gamma}$ is a smooth strictly convex hypersurface, any closed CAMC hypersurface $X$ is also smooth and the above expectation is correct, which was proved by [2], [4], [18] for $\gamma \equiv 1$, and by [16], [31], [26] and [3] for general

| Energy density | $n=2$ and $g(M)=0$ | $X$ is embedded | $X$ is stable |
| :--- | :---: | :---: | :---: |
| $(1) \gamma \equiv 1$ (isotropic) | true $([18])$ | true $([2])$ | true $([4])$ |
| $(2) \gamma \in C^{\infty}$ and strictly convex | true $([26])$ | true $([16])$ | true $([31])$ |
| $(3) \gamma \in C^{2}$ and convex | $?$ | $?$ | true $([24])$ |
| $(4) \gamma \in C^{2}$ | false $([20])$ | false $([20])$ | $?$ |

Table 1.1: Summarization of known results and our results
$\gamma$ satisfying a strong convexity condition (Table 1.1 summarizes the results). Here a CAMC hypersurface is said to be stable if the second variation of the energy $\mathcal{F}_{\gamma}$ for any volume-preserving variation is nonnegative. If $\gamma$ has less regularity or less convexity, the Wulff shape and CAMC hypersurfaces can have "edges". In the planar case, Frank Morgan [30] proved that every closed equilibrium curve for the anisotropic energy is the (possibly multiple of) Wulff shape for any continuous $\gamma$ satisfying the convexity condition.

However, the situation is not the same for non-convex $\gamma$. In this chapter, we will prove the following non-uniqueness results for CAMC hypersurfaces:

Theorem 1.1.1 ([20]). There exists a $C^{\infty}$ function $\gamma: S^{n} \rightarrow \mathbb{R}_{>0}$ which is not a convex integrand such that there exist closed embedded CAMC hypersurfaces in $\mathbb{R}^{n+1}$ for $\gamma$ each of which is not (any homothety or translation of) the Wulff shape.

Theorem 1.1.2 ([20]). There exists a $C^{\infty}$ function $\gamma: S^{2} \rightarrow \mathbb{R}_{>0}$ which is not a convex integrand such that there exist closed embedded CAMC surfaces in $\mathbb{R}^{3}$ with genus zero for $\gamma$ each of which is not (any homothety or translation of) the Wulff shape.

This shows that the Alexandrov-type theorem for embedded CAMC hypersurfaces and the Hopf-type theorem for genus zero surfaces cannot hold unless we assume the convexity of $\gamma$. These results are proved by giving examples (see section 1.4, 1.5, 1.6).

The same examples will be applied to the anisotropic mean curvature flow. In order to give the precise statement, we recall that the Cahn-Hoffman map $\xi_{\gamma}$ for $\gamma$ is the mapping $\xi_{\gamma}: S^{n} \rightarrow \mathbb{R}^{n+1}$ defined as

$$
\xi_{\gamma}(\nu)=\left.D \gamma\right|_{\nu}+\gamma(\nu) \nu, \quad \nu \in S^{n}
$$

here the tangent space $T_{\nu}\left(S^{n}\right)$ of $S^{n}$ at $\nu \in S^{n}$ is naturally identified with the $n$ dimensional linear subspace of $\mathbb{R}^{n+1}$. Let $X_{t}: M \rightarrow \mathbb{R}^{n+1}$ be one-parameter family of piecewise- $C^{2}$ weak $\underset{\sim}{\sim}$ mersions with anisotropic mean curvature $\Lambda_{t}$. Assume that the Cahn-Hoffman field $\widetilde{\xi}_{t}$ along $X_{t}$ is defined on $M$. If $X_{t}$ satisfies

$$
\partial X_{t} / \partial t=\Lambda_{t} \widetilde{\xi}_{t}
$$

it is called an anisotropic mean curvature flow, which diminishes the anisotropic energy if $\Lambda_{t} \not \equiv 0$. By a simple observation we show the following.

Theorem 1.1.3 ([20]). Let $c$ be a positive constant. Set

$$
X_{t}:=\sqrt{2(c-t)} \xi_{\gamma}, \quad t \leq c
$$

Then $X_{t}$ is a self-similar shrinking solution of the anisotropic mean curvature flow for $\gamma$, that is
(1) $\partial X_{t} / \partial t=\Lambda_{t} \widetilde{\xi}_{t}$, and
(2) $X_{t}$ is homothetic to $\xi_{\gamma}$ and it shrinks as $t$ increases.

By using this result and by giving examples, we prove the following result (see section 1.4, 1.5, 1.7).

Theorem 1.1.4 ([20]). There exists a $C^{\infty}$ function $\gamma: S^{n} \rightarrow \mathbb{R}_{>0}$ which is not a convex integrand such that there exist closed embedded self-similar shrinking solutions in $\mathbb{R}^{n+1}$ for $\gamma$ each of which is homeomorphic to $S^{n}$ and is not (any homothety or translation of) the Wulff shape.

In contrast with this result, the round sphere is the only closed embedded self-similar shrinking solution of mean curvature flow in $\mathbb{R}^{3}$ with genus zero ([7]).

Finally we give two conjetures about the uniqueness problems studied in Theorems 1.1.1, 1.1.2, and 1.1.3.

Conjecture 1. Assume that $\gamma: S^{n} \rightarrow \mathbb{R}_{>0}$ is of $C^{2}$. Let $X: M \rightarrow \mathbb{R}^{n+1}$ be a closed CAMC hypersurface. We assume that the $j$-th anisotropic mean curvature of $X$ for $\gamma$ is integrable for $j=1, \ldots, n$. Then, if $X$ satisfies at least one of the conditions (I), (II), (III) above, $X(M)$ is a subset of a homothety of the image $\xi_{\gamma}\left(S^{n}\right)$ of the Cahn-Hoffman $\operatorname{map} \xi_{\gamma}$.

Conjecture 2. Assume that $\gamma: S^{2} \rightarrow \mathbb{R}_{>0}$ is of $C^{2}$. Then any closed embedded selfsimilar shrinking solution of the anisotropic mean curvature flow for $\gamma$ in $\mathbb{R}^{3}$ with genus zero is a subset of a homothety of $\xi_{\gamma}\left(S^{2}\right)$.

### 1.2 Preliminaries

### 1.2.1 Piecewise- $C^{k}$ weak immersion and its anisotropic energy

First we recall the definition of a piecewise- $C^{r}$ weak immersion, $(r \in \mathbb{N})$, defined in [24]. Let $M=\bigcup_{i=1}^{k} M_{i}$ be an $n$-dimensional oriented compact connected $C^{\infty}$ manifold, where each $M_{i}$ is an $n$-dimensional connected compact submanifold of $M$ with boundary, and $M_{i} \cap M_{j}=\partial M_{i} \cap \partial M_{j},(i, j \in\{1, \cdots, k\}, i \neq j)$. We call a map $X: M \rightarrow \mathbb{R}^{n+1}$ a piecewise- $C^{r}$ weak immersion (or a piecewise- $C^{r}$ weakly immersed hypersurface) if $X$ satisfies the following conditions (A1), (A2), and (A3) $(i \in\{1, \cdots, k\})$.
(A1) $X$ is continuous, and each $X_{i}:=\left.X\right|_{M_{i}}: M_{i} \rightarrow \mathbb{R}^{n+1}$ is of $C^{r}$.
(A2) The restriction $\left.X\right|_{\text {Int } M_{i}}$ of $X$ to the interior $\operatorname{Int} M_{i}$ of $M_{i}$ is a $C^{r}$-immersion.
(A3) The unit normal vector field $\nu_{i}: \operatorname{Int} M_{i} \rightarrow S^{n}$ along $\left.X_{i}\right|_{\operatorname{Int} M_{i}}$ can be extended to a $C^{r-1}$-mapping $\nu_{i}: M_{i} \rightarrow S^{n}$. Here, if $\left(u^{1}, \cdots, u^{n}\right)$ is a local coordinate system in $M_{i},\left\{\nu_{i}, \partial / \partial u^{1}, \ldots, \partial / \partial u^{n}\right\}$ gives the canonical orientation in $\mathbb{R}^{n+1}$.

The anisotropic energy of a piecewise- $C^{1}$ weak immersion $X: M \rightarrow \mathbb{R}^{n+1}$ is defined as follows. Assume that $\gamma: S^{n} \rightarrow \mathbb{R}_{\geq 0}$ is a nonnegative continuous function. Let $\nu: M \backslash \mathcal{S}(X) \rightarrow S^{n}$ be the unit normal vector field along $\left.X\right|_{M \backslash \mathcal{S}(X)}$. The anisotropic energy $\mathscr{F}_{\gamma}(X)$ of $X$ is defined as

$$
\begin{equation*}
\mathscr{F}_{\gamma}(X):=\int_{M} \gamma(\nu) d A:=\sum_{i=1}^{k} \int_{M_{i}} \gamma\left(\nu_{i}\right) d A . \tag{2.1}
\end{equation*}
$$

If $\gamma \equiv 1, \mathscr{F}_{\gamma}(X)$ is the usual $n$-dimensional volume of the hypersurface $X$ (that is the $n$-dimensional volume of $M$ with the metric induced by $X$ ).

### 1.2.2 Wulff shapes and convex functions

Next we discuss about the Wulff shape and convex functions briefly (see [24] for details). As we remarked in the preface, for an energy density $\gamma: S^{n} \rightarrow \mathbb{R}_{>0}$ the minimizer of the anisotropic energy $\mathscr{F}_{\gamma}$ among all hypersurfaces in $\mathbb{R}^{n+1}$ enclosing the same volume is (up to translation and homothety) the Wulff shape $W_{\gamma}$ for $\gamma$. Moreover the Wulff shape can be written as

$$
W_{\gamma}:=\partial \bigcap_{\nu \in S^{n}}\left\{x \in \mathbb{R}^{n+1} \mid\langle x, \nu\rangle \leq \gamma(\nu)\right\}
$$

Every function $\gamma: S^{n} \rightarrow \mathbb{R}_{>0}$ determines its Wulff shape uniquely. In addition, if we denote

$$
\widehat{W}_{\gamma}:=\bigcap_{\nu \in S^{n}}\left\{x \in \mathbb{R}^{n+1} \mid\langle x, \nu\rangle \leq \gamma(\nu)\right\}
$$

the set $\widehat{W}_{\gamma}$ is always a convex body, i.e., $\widehat{W}_{\gamma}$ is compact, convex and has non-empty interior $0 \in \operatorname{Int} \widehat{W}_{\gamma}$.

Remark. The set $\widehat{W}_{\gamma}$ is also called the Wulff shape for $\gamma$. In the following we will identify $W_{\gamma}$ and $\widehat{W}_{\gamma}$, and use the same notation $W_{\gamma}$ if it is clear from the context.

Conversely, for any convex body $W$ with $0 \in \operatorname{Int} W$, we can construct a function $\gamma_{W}: S^{n} \rightarrow \mathbb{R}_{>0}$ as its support function:

$$
\gamma_{W}(\nu):=\max _{x \in W}\langle x, \nu\rangle, \quad \nu \in S^{n}
$$

We can check that $\gamma_{W}$ is continuous and the Wulff shape for $\gamma_{W}$ coincides with $W$. However, such a function is NOT unique in general.

Example 1.2.1. Let $W=[-1,1] \times[-1,1] \subset \mathbb{R}^{2}$ be a cube. Obviously $W$ is a convex body with $0 \in \operatorname{Int} W$ and we can show

$$
\gamma_{W}\left(\nu_{1}, \nu_{2}\right)=\left|\nu_{1}\right|+\left|\nu_{2}\right|
$$

by some calculation. However, the function $\gamma\left(\nu_{1}, \nu_{2}\right)=\left(\left|\nu_{1}\right|+\left|\nu_{2}\right|\right)^{2}$ also gives the same Wulff shape $W$.

This problem naturally leads to the notion of the convexity for functions. Let us recall the basic facts from the convex geometry.

For an energy density $\gamma: S^{n} \rightarrow \mathbb{R}_{\geq 0}$ we define the homogeneous extension of $\gamma$ as

$$
\bar{\gamma}: \mathbb{R}^{n+1} \rightarrow \mathbb{R}_{\geq 0}, \quad \bar{\gamma}(r \nu):=r \gamma(\nu), \quad r \geq 0, \nu \in S^{n} .
$$

Then we can easily check the following well-known fact:
Lemma 1.2.2. The following conditions are equivalent:
(1) $\bar{\gamma}$ is a convex function on $\mathbb{R}^{n+1}$, i.e., $\bar{\gamma}$ satisfies the triangle inequality : $\bar{\gamma}(x+y) \leq$ $\bar{\gamma}(x)+\bar{\gamma}(y)$ for all $x, y \in \mathbb{R}^{n+1}$.
(2) The set $F_{\gamma}:=\left\{x \in \mathbb{R}^{n+1} \mid \bar{\gamma}(x) \leq 1\right\}$ is a convex set in $\mathbb{R}^{n+1}$. $F_{\gamma}$ is sometimes called the Frank shape or the dual Wulff shape.
We say $\gamma: S^{n} \rightarrow \mathbb{R}_{\geq 0}$ to be convex if its homogeneous extension $\bar{\gamma}$ satisfies an either condition in the lemma. The second condition in the lemma gives a visual interpretation of the convexity (see the figure)

(a) Frank shape for $\gamma(\nu)=\left|\nu_{1}\right|+\left|\nu_{2}\right|$

(b) Frank shape for $\gamma(\nu)=\left(\left|\nu_{1}\right|+\left|\nu_{2}\right|\right)^{2}$

In particular, for a convex body $W$ with $0 \in \operatorname{Int} W$, the function $\gamma_{W}$ is convex. Moreover, this is the unique convex function among all continuous functions $\gamma: S^{n} \rightarrow$ $\mathbb{R}_{>0}$ which satisfies $W_{\gamma}=W$. For this reason, the function $\gamma_{W}$ is sometimes called the convex integrand.

### 1.2.3 Cahn-Hoffman map and anisotropic mean curvature

From now on we assume that $\gamma: S^{n} \rightarrow \mathbb{R}_{>0}$ is a positive-valued $C^{2}$ function. We define the map which parametrizes the Wulff shape. The map defined by

$$
\xi_{\gamma}: S^{n} \rightarrow \mathbb{R}^{n+1}, \quad \xi_{\gamma}(\nu):=D \gamma(\nu)+\gamma(\nu) \nu, \quad \nu \in S^{n}
$$

is called the Cahn-Hoffman map on $S^{n}$. Here $D \gamma$ denotes the gradient of $\gamma$ on $S^{n}$ and we identify it as a vector in $\mathbb{R}^{n+1}$. We put $\widetilde{W}_{\gamma}:=\xi_{\gamma}\left(S^{n}\right)$

Remark. Let $\bar{\gamma}: \mathbb{R}^{n+1} \rightarrow \mathbb{R}_{\geq 0}$ be the homogeneous extension of $\gamma$. Then the CahnHoffman map $\xi_{\gamma}$ can be expressed as

$$
\begin{equation*}
\xi_{\gamma}(\nu)=\left.\bar{\nabla} \bar{\gamma}\right|_{\nu}, \quad \nu \in S^{n} \tag{2.2}
\end{equation*}
$$

where $\bar{\nabla}$ is the gradient on $\mathbb{R}^{n+1}$. To see this, we take any $x \in \mathbb{R}^{n+1}$ and $v \in T_{x} \mathbb{R}^{n+1} \simeq$ $\mathbb{R}^{n+1}$, and we have

$$
\begin{aligned}
\langle\bar{\nabla} \bar{\gamma}, v\rangle_{\mathbb{R}^{n+1}} & =d \bar{\gamma}(v)=\gamma \cdot d(|x|)(v)+|x| \cdot d \gamma(v) \\
& =\gamma\langle\bar{\nabla}| x|, v\rangle+|x|\langle\bar{\nabla} \gamma, v\rangle=\langle D \gamma \cdot| x|+\gamma \cdot(x /|x|), v\rangle .
\end{aligned}
$$

By taking $x \in S^{n}$ we have the desired result.
The convexity condition is also naturally appeared in this situation since we have the following fact:

Proposition 1.2.3 ([24], Theorem 4.1, (ii)). The image of the Cahn-Hoffman map $\widetilde{W}_{\gamma}$ coincides with the Wulff shape $W_{\gamma}$ if and only if the funtion $\gamma$ is the convex integrand.

Now we define the anisotropic mean curvature for piecewise- $C^{2}$ hypersurfaces.
Definition 1.2.4 ([39], [25]). Let $X: M=\bigcup_{i=1}^{k} M_{i} \rightarrow \mathbb{R}^{n+1}$ be a piecewise- $C^{2}$ weak immersion and $\mathcal{S}(X)$ be its singularities. For the unit normal vector field $\nu=\nu_{X}$ along $X$, we define the anisotropic Gauss map $\widetilde{\xi}_{\gamma}: M \backslash \mathcal{S}(X) \rightarrow \mathbb{R}^{n+1}$ by $\widetilde{\xi}_{\gamma}:=\xi_{\gamma} \circ \nu_{X}$. If $\gamma \equiv 1$, this is the usual Gauss map of $X$.

We call the differential $-d \widetilde{\xi}_{\gamma}$ the anisotropic shape operator and the anisotropic mean curvature $\Lambda$ of $X$ is defined as

$$
\Lambda=\Lambda_{\gamma}:=\frac{1}{n} \operatorname{trace}\left(-d \widetilde{\xi}_{\gamma}\right)=\frac{1}{n}\left(n H \gamma-\operatorname{div}_{M} D \gamma\right)
$$

where $H$ is the usual mean curvature of $X$.
Remark. The eigenvalues of the anisotropic shape operator $-d \widetilde{\xi}_{\gamma}$ are called the anisotropic principal curvatures and we denote them as $\kappa_{1}^{\gamma}, \ldots, \kappa_{n}^{\gamma}$. The $j$-th mean curvature $H_{j}^{\gamma}$ of $X$ is defined as

$$
\binom{n}{j} H_{j}^{\gamma}:=\sum_{1 \leq i_{1}<\cdots<i_{j} \leq n} \kappa_{i_{1}}^{\gamma} \cdots \kappa_{i_{j}}^{\gamma} .
$$

By using the above notation, we have $\Lambda=H_{1}^{\gamma}$.
Example 1.2.5. Let us consider the planar $(n=1)$ case. Assume that the energy density $\gamma: S^{1} \rightarrow \mathbb{R}_{>0}$ is of class $C^{2}$, and denote $\gamma(\theta):=\gamma(\cos \theta, \sin \theta)$. Then, the Cahn-Hoffman map $\xi_{\gamma}(\theta):=\xi_{\gamma}(\cos \theta, \sin \theta)$ and its differential $\xi_{\gamma}^{\prime}(\theta)$ can be represented as follows:

$$
\xi_{\gamma}(\theta)=\gamma^{\prime}(\theta)(-\sin \theta, \cos \theta)+\gamma(\theta)(\cos \theta, \sin \theta), \quad \xi_{\gamma}^{\prime}(\theta)=\left(\gamma^{\prime \prime}(\theta)+\gamma(\theta)\right)(-\sin \theta, \cos \theta)
$$

For an arclength parametrized regular curve $X: I \subset \mathbb{R} \rightarrow \mathbb{R}^{2}$, we denote the unit normal of $X$ as $\nu(s)=(\cos \theta(s), \sin \theta(s))$. Then the curvature of $X$ is characterized by $\kappa=-d \theta / d s$ (note that we take the outward-pointing unit normal) and the differential $d\left(\xi_{\gamma} \circ \nu\right)$ can be computed as follows:

$$
d \xi_{\gamma}\binom{-\sin \theta(s)}{\cos \theta(s)}=\left(\gamma_{\theta \theta}+\gamma\right)(\theta(s)) \cdot \theta_{s}(s)\binom{-\sin \theta(s)}{\cos \theta(s)} .
$$

Therefore we have $\Lambda_{\gamma}(s)=\left(\gamma_{\theta \theta}+\gamma\right)(\theta(s)) \cdot \kappa(s)$. On the other hand, a basic calculation shows the curvature of the Cahn-Hoffman map $\xi_{\gamma}(\theta)$ is $\kappa_{\xi_{\gamma}}(\theta)=-1 /\left(\gamma^{\prime \prime}(\theta)+\gamma(\theta)\right)$ at the point satisfying $\gamma^{\prime \prime}(\theta)+\gamma(\theta) \neq 0$. Note that the image of the Cahn-Hoffman map is convex where $\gamma^{\prime \prime}+\gamma>0$, concave where $\gamma^{\prime \prime}+\gamma<0$ with respect to the outward-pointing unit normal, and the curvature diverges at the point satisfying $\gamma^{\prime \prime}+\gamma=0$. Therefore we have

$$
\Lambda_{\gamma}(s)=-\kappa(s) / \kappa_{\xi_{\gamma}}(\theta(s)), \quad \text { where } \quad \gamma_{\theta \theta}(\theta(s))+\gamma(\theta(s)) \neq 0
$$

In particular, the Cahn-Hoffman map $\xi_{\gamma}$ has constant anisotropic curvature -1 . Conversely, if we assume the curve $X$ has constant anisotropic curvature $\Lambda(s) \equiv-1$, then

$$
\kappa(s)=\kappa_{\xi_{\gamma}}(\theta(s)) \quad \text { where } \quad \gamma_{\theta \theta}(\theta(s))+\gamma(\theta(s)) \neq 0
$$

Therefore, by the fundamental theorem for plane curves, $X$ must be equal to a part of the Cahn-Hoffman map $\xi_{\gamma}$ (up to translation), where $\gamma_{\theta \theta}(\theta(s))+\gamma(\theta(s)) \neq 0$. In particular, if $\gamma^{\prime \prime}+\gamma>0$ on $S^{1}$, $X$ is equal to the Wulff shape with integer multiplicity. Note that this fact holds in more general settings [30].

### 1.3 The first variation formula of the anisotropic energy

Assume that $\gamma: S^{n} \rightarrow \mathbb{R}_{>0}$ is of class $C^{2}$. Let $M$ is a oriented connected compact $n$-dimensional $C^{\infty}$ manifold with smooth boundary $\partial M$ and $X: M \rightarrow \mathbb{R}^{n+1}$ be a $C^{2}$ mapping. We assume that following conditions:
(1) $\left.X\right|_{\text {Int } M}$ is a $C^{2}$-immersion.
(2) The unit normal vector field $\nu: \operatorname{Int} M \rightarrow S^{n}$ along Int $M$ can be extended to a $C^{1}$-mapping $\nu: M \rightarrow S^{n}$. Here, if $\left(u^{1}, \ldots, u^{n}\right)$ is a local coordinate system in $M$, $\left\{\nu, \partial / \partial u^{1}, \ldots, \partial / \partial u^{n}\right\}$ gives the canonical orientation in $\mathbb{R}^{n+1}$.

We consider a $C^{2}$ variation

$$
X_{t}: M \rightarrow \mathbb{R}^{n+1}, \quad X_{t}=X+t(\psi \nu+\eta)+O\left(t^{2}\right)
$$

where $\psi \nu, \eta$ are the normal and tangential components of the variational vector field respectively (hence these are of $C^{1}$ ):

$$
\delta X:={\frac{d X_{t}}{d t}}_{\mid t=0}=\psi \nu+\eta
$$

With these notations we have the following proposition:
Proposition 1.3.1 (First variation formula of $\mathscr{F}_{\gamma},[32]$ for $n=2,[24]$ for general $n$ ).

$$
\begin{aligned}
\delta \mathscr{F}_{\gamma} & =\int_{M} \psi\left(\operatorname{div}_{M} D \gamma-n H \gamma\right) d A_{X}+\int_{\partial M}\langle\delta X,-R(P(\xi))\rangle d \widetilde{s} \\
& =-n \int_{M} \Lambda\langle\delta X, \nu\rangle d A_{X}+\int_{\partial M}\langle\delta X,-R(P(\xi))\rangle d \widetilde{s}
\end{aligned}
$$

where $N$ is the outer unit conormal of $X$ along $\partial M, d \widetilde{s}$ is the ( $n-1$ )-dimensional volume form of $\partial M$ induced by $X$ and $R$ is the $\pi / 2$-rotation in ( $N, \nu$ )-plane.

Remark. We will discuss about the discrete curves and surfaces and these can be considered as a special case of the piecewise weak immersions. For example, in the triangulated surface case, each piece $M_{i}$ is a flat triangle and therefore we only have the boundary terms. However, if we consider these as a discrete objects, the "boundary terms" are not necessarily the boundary. Therefore the discrete differential geometry developed in the next chapter always lies in between the smooth and discrete objects, and its importance is the following viewpoint : how to extract the discrete curvature information from the "boundary terms" of the first variation formula.

Recall that $\mathcal{S}(X)$ is the set of all singular points of $X$. For any $p \in M_{i} \backslash \mathcal{S}(X)$, we may write $\widetilde{\xi}(p):=\widetilde{\xi}_{\gamma}(p):=\widetilde{\xi}_{i}(p)$.

Theorem 1.3.2 (Euler-Lagrange equations, Koiso [24]. For $n=2$, see B. Palmer [32]). A piecewise- $C^{2}$ weak immersion $X: M=\bigcup_{i=1}^{k} M_{i} \rightarrow \mathbb{R}^{n+1}$ is a critical point of the anisotropic energy $\mathscr{F}_{\gamma}(X)=\int_{M} \gamma(\nu) d A$ for volume-preserving variations if and only if
(1) The anisotropic mean curvature $\Lambda$ of $X$ is constant on $M \backslash S(X)$, and
(2) $\widetilde{\xi}_{i}(\zeta)-\widetilde{\xi}_{j}(\zeta) \in T_{\zeta} M_{i} \cap T_{\zeta} M_{j}=T_{\zeta}\left(\partial M_{i} \cap \partial M_{j}\right)$ holds at any $\zeta \in \partial M_{i} \cap \partial M_{j}$, where a tangent space of a submanifold of $\mathbb{R}^{n+1}$ is naturally identified with a linear subspace of $\mathbb{R}^{n+1}$.

Example 1.3.3. Let us continue Example 1.2.5. When $n=1$, the condition (2) in the above theorem is equivalent to $\widetilde{\xi}_{i}(\zeta)=\widetilde{\xi}_{j}(\zeta)$ at the meeting point of the curves $X\left(M_{i}\right)$ and $X\left(M_{j}\right)$. Therefore, we have the following statement:

Proposition 1.3.4. Let $\gamma: S^{1} \rightarrow \mathbb{R}_{>0}$ be of class $C^{2}$ and $\gamma^{\prime \prime}+\gamma=0$ only on the set $\mathcal{S}_{\gamma} \subset S^{1}$, where $\mathcal{S}_{\gamma}=\left\{\theta_{1}, \ldots, \theta_{m}\right\}$ is a set of finite points. Then a closed piecewise $C^{2}$-smooth curve $X: I \subset \mathbb{R} \rightarrow \mathbb{R}^{2}$ which is a critical point of the anisotropic energy $\mathscr{F}_{\gamma}$ with $\Lambda=-1$ coincides with a part of the Cahn-Hoffman map $\xi_{\gamma}$ up to translation with integer multiplicity. In particular, if $\gamma^{\prime \prime}+\gamma>0$ on $S^{1} \backslash \mathcal{S}_{\gamma}$, then $X$ coincides with the Wulff shape up to translation with integer multiplicity.

Remark . The last statement is a special case of the result of Morgan [30].

Proof. We already saw that the regular part of such a curve $X(s)$ coincides with the Cahn-Hoffman map $\xi_{\gamma}$ except for the points satisfying $\theta(s)=\theta_{j}$ for some $j \in\{1, \ldots, m\}$, where $\theta(s)$ is defined by the unit normal vector field $\nu(s)=(\cos \theta(s), \sin \theta(s))$. By the assumption, such points $s \in I$ are finite and must lie on the boundaries of each piece since such points correspond to the singularities of the Cahn-Hoffman map $\xi_{\gamma}$. By the condition of (2) in Theorem 1.3.2, each piece must glue together.

### 1.4 An example of dimension 1

We give an example which will be used to prove Theorems 1.1.1, 1.1.2, and 1.1.3 in the following sections. Throughout this section $\gamma: S^{1} \rightarrow \mathbb{R}_{>0}$ is the function defined by

$$
\begin{equation*}
\gamma\left(e^{i \theta}\right):=\cos ^{6} \theta+\sin ^{6} \theta \tag{4.1}
\end{equation*}
$$

where $\mathbb{R}^{2}$ is identified with $\mathbb{C}$.
Lemma 1.4.1. The Cahn-Hoffman map $\xi_{\gamma}: S^{1} \rightarrow \mathbb{R}^{2}$ for $\gamma$ is represented as follows.

$$
\begin{align*}
\xi_{\gamma}\left(e^{i \theta}\right)= & \left((\cos \theta)\left(\cos ^{6} \theta+6 \cos ^{4} \theta \sin ^{2} \theta-5 \sin ^{6} \theta\right)\right. \\
& \left.(\sin \theta)\left(-5 \cos ^{6} \theta+6 \cos ^{4} \theta \sin ^{2} \theta+\sin ^{6} \theta\right)\right) \tag{4.2}
\end{align*}
$$

In other words, we have

$$
\begin{equation*}
\xi_{\gamma}(\nu)=\left(\nu_{1}\left(\nu_{1}^{6}+6 \nu_{1}^{4} \nu_{2}^{2}-5 \nu_{2}^{6}\right), \nu_{2}\left(-5 \nu_{1}^{6}+6 \nu_{1}^{2} \nu_{2}^{4}+\nu_{2}^{6}\right)\right), \quad \nu=\left(\nu_{1}, \nu_{2}\right) \in S^{1} \tag{4.3}
\end{equation*}
$$

Set $\nu_{2}= \pm \sqrt{1-\nu_{1}^{2}}$. Then

$$
\begin{align*}
\xi_{\gamma}(\nu) & =\left(-\nu_{1}\left(9 \nu_{1}^{4}-15 \nu_{1}^{2}+5\right), \mp \sqrt{1-\nu_{1}^{2}}\left(9 \nu_{1}^{4}-3 \nu_{1}^{2}-1\right)\right)  \tag{4.4}\\
& =\left(-\left(9 \cos ^{4} \theta-15 \cos ^{2} \theta+5\right) \cos \theta,-\left(9 \cos ^{4} \theta-3 \cos ^{2} \theta-1\right) \sin \theta\right) \tag{4.5}
\end{align*}
$$

holds.
Proof. We use the formula (2.2), that is

$$
\xi_{\gamma}(\nu)=\left.\bar{\nabla} \bar{\gamma}\right|_{\nu}, \quad \nu \in S^{n}
$$

where $\bar{\gamma}$ is given by

$$
\begin{equation*}
\bar{\gamma}\left(\nu_{1}, \nu_{2}\right)=\frac{\nu_{1}^{6}+\nu_{2}^{6}}{\left(\nu_{1}^{2}+\nu_{2}^{2}\right)^{5 / 2}} . \tag{4.6}
\end{equation*}
$$

Hence, we have

$$
\begin{align*}
& \bar{\gamma}_{\nu_{1}}=\frac{\nu_{1}\left(\nu_{1}^{6}+6 \nu_{1}^{4} \nu_{2}^{2}-5 \nu_{2}^{6}\right)}{\left(\nu_{1}^{2}+\nu_{2}^{2}\right)^{7 / 2}}  \tag{4.7}\\
& \bar{\gamma}_{\nu_{2}}=\frac{\nu_{2}\left(-5 \nu_{1}^{6}+6 \nu_{1}^{2} \nu_{2}^{4}+\nu_{2}^{6}\right)}{\left(\nu_{1}^{2}+\nu_{2}^{2}\right)^{7 / 2}} \tag{4.8}
\end{align*}
$$

(2.2) with (4.2) and (4.8) gives

$$
\begin{equation*}
\xi_{\gamma}(\nu)=\left(\nu_{1}\left(\nu_{1}^{6}+6 \nu_{1}^{4} \nu_{2}^{2}-5 \nu_{2}^{6}\right), \nu_{2}\left(-5 \nu_{1}^{6}+6 \nu_{1}^{2} \nu_{2}^{4}+\nu_{2}^{6}\right)\right), \quad \nu=\left(\nu_{1}, \nu_{2}\right) \in S^{1} \tag{4.9}
\end{equation*}
$$

which gives the desired formulas (4.2), (4.3). Inserting $\nu_{2}= \pm \sqrt{1-\nu_{1}^{2}}$ to (4.3), we obtain (4.4), (4.5).

The image $\xi_{\gamma}\left(S^{1}\right)$ of the Cahn-Hoffman map $\xi_{\gamma}$ and the Wulff shape $W_{\gamma}$ are shown in Figure 1.4. Lemma 1.4.3 below gives rigorous explanations of these shapes. Now we


Figure 1.3: The image $\xi_{\gamma}\left(S^{1}\right)$ of the Cahn-Hoffman map (left) and the Wulff shape $W_{\gamma}$ (right) for $\gamma$ defined by (4.1). $W_{\gamma}$ is a subset of $\xi_{\gamma}\left(S^{1}\right)$.
study the singular points of $\xi_{\gamma}$.
Lemma 1.4.2. (i) The the set $S\left(\xi_{\gamma}\right)$ of all singular points of $\xi_{\gamma}$ is given by

$$
\left.\begin{array}{rl}
S\left(\xi_{\gamma}\right)=\left\{p_{1}:=\left(\frac{\sqrt{5}+1}{2 \sqrt{3}}, \frac{\sqrt{5}-1}{2 \sqrt{3}}\right), p_{2}:=\left(\frac{\sqrt{5}-1}{2 \sqrt{3}}, \frac{\sqrt{5}+1}{2 \sqrt{3}}\right)\right. \\
p_{3} & :=\left(-\frac{\sqrt{5}-1}{2 \sqrt{3}}, \frac{\sqrt{5}+1}{2 \sqrt{3}}\right), p_{4}:=\left(-\frac{\sqrt{5}+1}{2 \sqrt{3}}, \frac{\sqrt{5}-1}{2 \sqrt{3}}\right) \\
p_{5} & :=\left(-\frac{\sqrt{5}+1}{2 \sqrt{3}},-\frac{\sqrt{5}-1}{2 \sqrt{3}}\right), p_{6}:=\left(-\frac{\sqrt{5}-1}{2 \sqrt{3}},-\frac{\sqrt{5}+1}{2 \sqrt{3}}\right) \\
p_{7} & :=\left(\frac{\sqrt{5}-1}{2 \sqrt{3}},-\frac{\sqrt{5}+1}{2 \sqrt{3}}\right), p_{8} \tag{4.10}
\end{array}:=\left(\frac{\sqrt{5}+1}{2 \sqrt{3}},-\frac{\sqrt{5}-1}{2 \sqrt{3}}\right)\right\} .
$$

(ii) Set

$$
p_{j}=\left(\cos \theta_{j}, \sin \theta_{j}\right), \quad j=1, \ldots, 8,
$$

where $\theta_{1}$ is chosen so that $0<\theta_{1}<\pi / 2$ holds. Then, by choosing suitable principal values for $\theta_{1}, \cdots, \theta_{8}$, one can write

$$
\begin{align*}
& \theta_{2}=\pi / 2-\theta_{1}, \theta_{3}=\pi / 2+\theta_{1}, \theta_{4}=\pi-\theta_{1}, \theta_{5}=\pi+\theta_{1},  \tag{4.11}\\
& \theta_{6}=3 \pi / 2-\theta_{1}, \theta_{7}=3 \pi / 2+\theta_{1}, \theta_{8}=-\theta_{1} . \tag{4.12}
\end{align*}
$$

(iii)

$$
\begin{align*}
& \xi_{\gamma}\left(p_{8}\right)=\xi_{\gamma}\left(p_{3}\right)=\frac{2}{\sqrt{3}}(1,1), \xi_{\gamma}\left(p_{2}\right)=\xi_{\gamma}\left(p_{5}\right)=\frac{2}{\sqrt{3}}(-1,1),  \tag{4.13}\\
& \xi_{\gamma}\left(p_{4}\right)=\xi_{\gamma}\left(p_{7}\right)=\frac{2}{\sqrt{3}}(-1,-1), \xi_{\gamma}\left(p_{6}\right)=\xi_{\gamma}\left(p_{1}\right)=\frac{2}{\sqrt{3}}(1,-1) \tag{4.14}
\end{align*}
$$

holds.
Remark. Choose the principal value for $\theta_{1}$ so that $0<\theta_{1}<\pi / 2$ is satisfied. Then, by computation using Mathematica ver.11.2.0.0, we get

$$
\begin{equation*}
\theta_{1} \approx 0.1161397636 \pi \tag{4.15}
\end{equation*}
$$

Proof of Lemma 1.4.2. Using (4.4), we obtain

$$
\begin{equation*}
\xi_{\gamma}^{\prime}\left(\nu_{1}\right)=\left(-5\left(9 \nu_{1}^{4}-9 \nu_{1}^{2}+1\right), \pm \frac{5 \nu_{1}}{\sqrt{1-\nu_{1}^{2}}}\left(9 \nu_{1}^{4}-9 \nu_{1}^{2}+1\right)\right) . \tag{4.16}
\end{equation*}
$$

Hence, $\left(\nu_{1}, \nu_{2}\right) \in S^{1}$ is a singular point of $\xi_{\gamma}$ if and only if

$$
\begin{equation*}
9 \nu_{1}^{4}-9 \nu_{1}^{2}+1=0 \tag{4.17}
\end{equation*}
$$

holds. (4.17) is equivalent to

$$
\begin{equation*}
\nu_{1}=\frac{\sqrt{5} \pm 1}{2 \sqrt{3}},-\frac{\sqrt{5} \pm 1}{2 \sqrt{3}} \tag{4.18}
\end{equation*}
$$

which proves (i).
(ii) is a consequence of (i).
(iii) is obtained by inserting (4.10) to (4.4).

Set

$$
\xi_{\gamma}(\theta)=\left(\left(\xi_{\gamma}\right)_{x}(\theta),\left(\xi_{\gamma}\right)_{y}(\theta)\right):=\xi_{\gamma}\left(e^{i \theta}\right)
$$

Lemma 1.4.3. Increase and decrease of $\left(\xi_{\gamma}\right)_{x}(\theta),\left(\xi_{\gamma}\right)_{y}(\theta)$, and the sign of $A=D^{2} \gamma+\gamma \cdot 1$ are given by the following table.

| $\theta$ | $\theta_{8}$ |  | 0 |  | $\theta_{1}$ |  | $\theta_{2}$ |  | $\pi / 2$ |  | $\theta_{3}$ |  | $\theta_{4}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\left(\xi_{\gamma}\right)_{x}^{\prime}(\theta)$ | 0 | - | 0 | + | 0 | - | 0 | + | 5 | + | 0 | - | 0 |
| $\left(\xi_{\gamma}\right)_{y}^{\prime}(\theta)$ | 0 | - | -5 | - | 0 | + | 0 | - | 0 | + | 0 | - | 0 |
| $\left(\xi_{\gamma}\right)_{x}(\theta)$ | $2 / \sqrt{3}$ | $\searrow$ | 1 | $\nearrow$ | $2 / \sqrt{3}$ | $\searrow$ | $-2 / \sqrt{3}$ | $\nearrow$ | 0 | $\nearrow$ | $2 / \sqrt{3}$ | $\searrow$ | $-2 / \sqrt{3}$ |
| $\left(\xi_{\gamma}\right)_{y}(\theta)$ | $2 / \sqrt{3}$ | $\searrow$ | 0 | $\searrow$ | $-2 / \sqrt{3}$ | $\nearrow$ | $2 / \sqrt{3}$ | $\searrow$ | 1 | $\nearrow$ | $2 / \sqrt{3}$ | $\searrow$ | $-2 / \sqrt{3}$ |
| $A=D^{2} \gamma+\gamma \cdot 1$ | 0 | - | - | - | 0 | + | 0 | - | - | - | 0 | + | 0 |


| $\theta$ | $\theta_{4}$ |  | $\pi$ |  | $\theta_{5}$ |  | $\theta_{6}$ |  | $3 \pi / 2$ |  | $\theta_{7}$ |  | $\theta_{8}+2 \pi$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\left(\xi_{\gamma}\right)_{x}^{\prime}(\theta)$ | 0 | + | 0 | - | 0 | + | 0 | - | -5 | - | 0 | + | 0 |
| $\left(\xi_{\gamma}\right)_{y}^{\prime}(\theta)$ | 0 | + | 5 | + | 0 | - | 0 | + | 0 | - | 0 | + | 0 |
| $\left(\xi_{\gamma}\right)_{x}(\theta)$ | $-2 / \sqrt{3}$ | $\nearrow$ | -1 | $\searrow$ | $-2 / \sqrt{3}$ | $\nearrow$ | $2 / \sqrt{3}$ | $\searrow$ | 0 | $\searrow$ | $-2 / \sqrt{3}$ | $\nearrow$ | $2 / \sqrt{3}$ |
| $\left(\xi_{\gamma}\right)_{y}(\theta)$ | $-2 / \sqrt{3}$ | $\nearrow$ | 0 | $\nearrow$ | $2 / \sqrt{3}$ | $\searrow$ | $-2 / \sqrt{3}$ | $\nearrow$ | -1 | $\searrow$ | $-2 / \sqrt{3}$ | $\nearrow$ | $2 / \sqrt{3}$ |
| $A=D^{2} \gamma+\gamma \cdot 1$ | 0 | - | - | - | 0 | + | 0 | - | - | - | 0 | + | 0 |

Proof. Using (4.5), we obtain

$$
\begin{align*}
\xi_{\gamma}^{\prime}(\theta) & =\left(\left(\xi_{\gamma}\right)_{x}^{\prime}(\theta),\left(\xi_{\gamma}\right)_{y}^{\prime}(\theta)\right) \\
& =\left(5\left(9 \cos ^{4} \theta-9 \cos ^{2} \theta+1\right) \sin \theta,-5\left(9 \cos ^{4} \theta-9 \cos ^{2} \theta+1\right) \cos \theta\right) \tag{4.19}
\end{align*}
$$

For $\gamma(\theta):=\gamma\left(e^{i \theta}\right)=\cos ^{6} \theta+\sin ^{6} \theta$, we compute

$$
\begin{gathered}
\gamma^{\prime}(\theta)=-6 \cos ^{5} \theta \sin \theta+6 \sin ^{5} \theta \cos \theta \\
\gamma^{\prime \prime}(\theta)=6\left(5 \cos ^{2} \theta \sin ^{2} \theta-\cos ^{6} \theta-\sin ^{6} \theta\right)
\end{gathered}
$$

and hence

$$
\begin{equation*}
A=\gamma^{\prime \prime}(\theta)+\gamma(\theta)=-5\left(9 \cos ^{4} \theta-9 \cos ^{2} \theta+1\right)=-\left(\xi_{\gamma}\right)_{x}^{\prime}(\theta) \cdot(\sin \theta)^{-1} \tag{4.20}
\end{equation*}
$$

holds, here the last equality is valid only when $\sin \theta \neq 0$. These observations with Lemmas 1.4.1, 1.4.2 give the desired result.

Using Lemma 1.4.1, by simple calculations, we obtain the following result.
Lemma 1.4.4. Let $\rho_{1}, \rho_{2} \in(0, \pi / 2)$ be the solutions of

$$
\begin{equation*}
\cos \rho_{1}=\sqrt{\frac{1+\sqrt{5}}{6}}, \cos \rho_{2}=\sqrt{\frac{5-\sqrt{5}}{6}} \tag{4.21}
\end{equation*}
$$

respectively. Then
(i) The inequality

$$
\begin{equation*}
0<\theta_{1}<\rho_{1}<\pi / 4<\rho_{2}<\theta_{2}<\pi / 2 \tag{4.22}
\end{equation*}
$$

holds.
(ii) Set

$$
\begin{equation*}
\alpha:=\sqrt{\frac{2+2 \sqrt{5}}{3}}(\sqrt{5}-2) . \tag{4.23}
\end{equation*}
$$

The "inner self-intersection points" of $\xi_{\gamma}$ are the following four points.
$Q_{1}:=(\alpha, 0)=\xi_{\gamma}\left(\rho_{1}\right)=\xi_{\gamma}\left(\rho_{2}+3 \pi / 2\right), Q_{2}:=(0, \alpha)=\xi_{\gamma}\left(\rho_{2}\right)=\xi_{\gamma}\left(\rho_{1}+\pi / 2\right)$,
$Q_{3}:=(-\alpha, 0)=\xi_{\gamma}\left(\rho_{1}+\pi\right)=\xi_{\gamma}\left(\rho_{2}+\pi / 2\right), Q_{4}:=(0,-\alpha)=\xi_{\gamma}\left(\rho_{2}+\pi\right)=\xi_{\gamma}\left(\rho_{1}+3 \pi / 2\right)$.
$Q_{1}, Q_{2}, Q_{3}, Q_{4}$ are the vertices of $W_{\gamma}$ (Figure 1.4, right).

Remark . By computation using Mathematica ver.11.2.0.0, we get

$$
\begin{equation*}
\rho_{1} \approx 0.2374632441 \pi, \quad \rho_{2}=\pi / 2-\rho_{1} \approx 0.2625367559 \pi, \quad \alpha \approx 0.3467370642 \tag{4.24}
\end{equation*}
$$

By using the previous lemmas 1.4.1, 1.4.2, 1.4.3, and 1.4.4 with Fact ?? (i), we obtain the following:

Proposition 1.4.5. Let $\gamma: S^{1} \rightarrow \mathbb{R}_{>0}$ be the function defined by

$$
\begin{equation*}
\gamma\left(e^{i \theta}\right):=\cos ^{6} \theta+\sin ^{6} \theta \tag{4.25}
\end{equation*}
$$

Let $\theta_{1}, \theta_{2}$ be the constants defined in Lemma 1.4.2, and let $\rho_{1}, \rho_{2}$ be the constants defined in Lemma1.4.4.
(i) The Cahn-Hoffman map $\xi_{\gamma}: S^{1} \rightarrow \mathbb{R}^{2}$ is represented as

$$
\begin{align*}
\xi_{\gamma}\left(e^{i \theta}\right)= & \left((\cos \theta)\left(\cos ^{6} \theta+6 \cos ^{4} \theta \sin ^{2} \theta-5 \sin ^{6} \theta\right)\right. \\
& \left.(\sin \theta)\left(-5 \cos ^{6} \theta+6 \cos ^{4} \theta \sin ^{2} \theta+\sin ^{6} \theta\right)\right) \tag{4.26}
\end{align*}
$$

(Figure 1.4).
(ii) The Wulff shape $W_{\gamma}$ is given by

$$
\begin{align*}
W_{\gamma}= & \xi_{\gamma}\left(\left\{\rho_{1} \leq \theta \leq \rho_{2}\right\} \cup\left\{\rho_{1}+\pi / 2 \leq \theta \leq \rho_{2}+\pi / 2\right\}\right. \\
& \left.\cup\left\{\rho_{1}+\pi \leq \theta \leq \rho_{2}+\pi\right\} \cup\left\{\rho_{1}+3 \pi / 2 \leq \theta \leq \rho_{2}+3 \pi / 2\right\}\right) \tag{4.27}
\end{align*}
$$

(Figure 1.4a).
(iii) The following four closed curves which are subsets of $\xi_{\gamma}\left(S^{1}\right)$ are closed piecewise$C^{\infty}$ CAMC curves (Figures 1.4b, 1.4c, 1.4d, 1.4e).

$$
\begin{gather*}
\left(C_{\gamma}\right)_{1}:=\xi_{\gamma}\left(\left\{\theta_{2} \leq \theta \leq \theta_{3}\right\} \cup\left\{\theta_{4} \leq \theta \leq \theta_{5}\right\} \cup\left\{\theta_{6} \leq \theta \leq \theta_{7}\right\} \cup\left\{\theta_{8} \leq \theta \leq \theta_{1}\right\}\right)  \tag{4.28}\\
\left(C_{\gamma}\right)_{2}:=\xi_{\gamma}\left(\left\{\theta_{1} \leq \theta \leq \theta_{2}\right\} \cup\left\{\theta_{5} \leq \theta \leq \theta_{6}\right\}\right)  \tag{4.29}\\
\left(C_{\gamma}\right)_{3}:=\xi_{\gamma}\left(\left\{\theta_{8} \leq \theta \leq \theta_{1}\right\} \cup\left\{\theta_{3} \leq \theta \leq \rho_{2}+\pi / 2\right\} \cup\left\{\rho_{1}+\pi \leq \theta \leq \theta_{6}\right\}\right),  \tag{4.30}\\
\left(C_{\gamma}\right)_{4}:=\xi_{\gamma}\left(\left\{\theta_{1} \leq \theta \leq \rho_{1}\right\} \cup\left\{\rho_{2} \leq \theta \leq \theta_{2}\right\} \cup\left\{\theta_{3} \leq \theta \leq \rho_{1}+\pi / 2\right\}\right. \\
\cup\left\{\rho_{2}+\pi / 2 \leq \theta \leq \theta_{4}\right\} \cup\left\{\theta_{5} \leq \theta \leq \rho_{1}+\pi\right\} \cup\left\{\rho_{2}+\pi \leq \theta \leq \theta_{6}\right\} \\
\left.\cup\left\{\theta_{7} \leq \theta \leq \rho_{1}+3 \pi / 2\right\} \cup\left\{\rho_{2}+3 \pi / 2 \leq \theta \leq \theta_{8}+2 \pi\right\}\right) . \tag{4.31}
\end{gather*}
$$

The anisotropic (mean) curvature for the outward-pointing normal is -1 .
(iv) The following closed curve which is a subset of $\xi_{\gamma}\left(S^{1}\right)$ is a closed piecewise- $C^{\infty}$ curve (Figure 1.4f).

$$
\begin{equation*}
\left(C_{\gamma}\right)_{5}:=\xi_{\gamma}\left(\left\{-\rho_{1} \leq \theta \leq \rho_{1}\right\}\right) \tag{4.32}
\end{equation*}
$$

Its anisotropic (mean) curvature is not constant, because, for the outward-pointing normal, it is -1 at each point in the solid curve:

$$
\begin{equation*}
\xi_{\gamma}\left(\left\{\theta_{8} \leq \theta \leq \theta_{1}\right\}\right) \tag{4.33}
\end{equation*}
$$

while it is 1 at each point in the dashed curves:

$$
\begin{equation*}
\xi_{\gamma}\left(\left\{-\rho_{1} \leq \theta \leq \theta_{8}\right\} \cup\left\{\theta_{1} \leq \theta \leq \rho_{1}\right\}\right) \tag{4.34}
\end{equation*}
$$


(a) Wulff shape $W_{\gamma}$

(d) CAMC curve $\left(C_{\gamma}\right)_{3}$

Figure 1.4: Some of the closed curves which are subsets of $\xi_{\gamma}\left(S^{1}\right)$ for $\gamma$ defined by (4.1) (cf. Figure 1.4). (a), (b), (c), (d), (e) : The anisotropic (mean) curvature for the outward-pointing normal is -1 . (f) : For the outward-pointing normal, the anisotropic (mean) curvature is -1 at each point in the solid curve, while it is 1 at each point in the dashed curves. Hence, this curve is not CAMC.

Remark . Proposition 1.4.5 proves Theorem 1.1.1 for $n=1$.
Remark . In Proposition 1.4.5, we gave six closed piecewise- $C^{\infty}$ curves which are subsets of $\xi_{\gamma}\left(S^{1}\right)$. Five of them were CAMC and the other one was not CAMC. We should remark that there are more piecewise- $C^{\infty}$ CAMC closed curves and non-CAMC closed curves. Figure 1.5 gives all of the other closed CAMC curves included in $\xi_{\gamma}\left(S^{1}\right)$ (up to congruence in $\mathbb{R}^{2}$ ).

### 1.5 Higher dimensional examples

In this section we give two higher dimensional examples by rotating $\gamma$ which was defined by (4.1) and studied in the section 1.4 details.

Regarding (4.6), we consider the function

$$
\bar{\gamma}_{1}^{0}\left(\nu_{1}, \nu_{3}\right)=\frac{\nu_{1}^{6}+\nu_{3}^{6}}{\left(\nu_{1}^{2}+\nu_{3}^{2}\right)^{5 / 2}}
$$



Figure 1.5: Closed CAMC curves which are subsets of $\xi_{\gamma}\left(S^{1}\right)$ for $\gamma$ defined by (4.1) (cf. Figure 1.4). For all of these eight curves, the anisotropic (mean) curvature for the outward-pointing normal is -1 . They and the five curves in Figures 1.4a, 1.4b, 1.4c, $1.4 \mathrm{~d}, 1.4 \mathrm{e}$ give all closed CAMC curves included in $\xi_{\gamma}\left(S^{1}\right)$ (up to congruence in $\mathbb{R}^{2}$
defined on the $\left(\nu_{1}, \nu_{3}\right)$-plane. We denote the restriction of $\gamma_{1}^{0}$ to $S^{1}$ by $\gamma_{1}^{0}$. Then, by Lemma 1.4.1, the Cahn-Hoffman map $\xi_{\gamma_{1}^{0}}: S^{1} \rightarrow \mathbb{R}^{2}$ for $\gamma_{1}^{0}$ is represented as follows (Figure 1.6a).

$$
\begin{align*}
\xi_{\gamma_{1}^{0}}(\nu)= & \left(\nu_{1}\left(\nu_{1}^{6}+6 \nu_{1}^{4} \nu_{3}^{2}-5 \nu_{3}^{6}\right), \nu_{3}\left(-5 \nu_{1}^{6}+6 \nu_{1}^{2} \nu_{3}^{4}+\nu_{3}^{6}\right)\right)  \tag{5.1}\\
= & \left((\cos \theta)\left(\cos ^{6} \theta+6 \cos ^{4} \theta \sin ^{2} \theta-5 \sin ^{6} \theta\right),\right. \\
& \left.(\sin \theta)\left(-5 \cos ^{6} \theta+6 \cos ^{4} \theta \sin ^{2} \theta+\sin ^{6} \theta\right)\right) . \tag{5.2}
\end{align*}
$$

$\left(\nu=\left(\nu_{1}, \nu_{3}\right)=(\cos \theta, \sin \theta) \in S^{1}\right)$. The higher dimensional example obtained by rotating $\bar{\gamma}_{1}^{0}$ around the $\nu_{3}$-axis is given by

$$
\begin{equation*}
\bar{\gamma}_{1}\left(\nu_{1}, \nu_{2}, \nu_{3}\right)=\frac{\left(\nu_{1}^{2}+\nu_{2}^{2}\right)^{3}+\nu_{3}^{6}}{\left(\nu_{1}^{2}+\nu_{2}^{2}+\nu_{3}^{2}\right)^{5 / 2}}, \quad\left(\nu_{1}, \nu_{2}, \nu_{3}\right) \in \mathbb{R}^{3} \tag{5.3}
\end{equation*}
$$

The corresponding Cahn-Hoffman map $\xi_{\gamma_{1}}: S^{2} \rightarrow \mathbb{R}^{3}$ is given as follows (Figure 1.6b).

$$
\begin{align*}
\xi_{\gamma_{1}}(\nu)= & \left((\cos \theta)\left(\cos ^{6} \theta+6 \cos ^{4} \theta \sin ^{2} \theta-5 \sin ^{6} \theta\right)(\cos \rho),\right. \\
& (\cos \theta)\left(\cos ^{6} \theta+6 \cos ^{4} \theta \sin ^{2} \theta-5 \sin ^{6} \theta\right)(\sin \rho), \\
& \left.(\sin \theta)\left(-5 \cos ^{6} \theta+6 \cos ^{4} \theta \sin ^{2} \theta+\sin ^{6} \theta\right)\right), \tag{5.4}
\end{align*}
$$

$\left(\nu=(\cos \theta \cos \rho, \cos \theta \sin \rho, \sin \theta) \in S^{2}\right)$.
By the same way as in the section 1.4, we get closed piecewise- $C^{\infty}$ CAMC surfaces and closed piecewise- $C^{\infty}$ non-CAMC surfaces for $\gamma_{1}$ which are subsets of $\xi_{\gamma_{1}}\left(S^{2}\right)$ and which are not the Wulff shape $W_{\gamma_{1}}$ (up to homothety and translation). In fact, we have the following:

Proposition 1.5.1. Consider $\gamma_{1}: S^{2} \rightarrow \mathbb{R}_{>0}$ defined by

$$
\begin{equation*}
\gamma_{1}\left(\nu_{1}, \nu_{2}, \nu_{3}\right)=\left(\nu_{1}^{2}+\nu_{2}^{2}\right)^{3}+\nu_{3}^{6}, \quad\left(\nu_{1}, \nu_{2}, \nu_{3}\right) \in S^{2} \tag{5.5}
\end{equation*}
$$

The Wulff shape $W_{\gamma_{1}}$ is the surface of revolution (Figure 1.7a) given by rotating $W_{\gamma_{1}^{0}}$ (Figure 1.4a) around the vertical axis. The two piecewise- $C^{\infty}$ closed surfaces (Figures $1.7 \mathrm{~b}, 1.7 \mathrm{c}$ ) given by rotating the closed curves $\left(C_{\gamma}\right)_{1},\left(C_{\gamma}\right)_{4}$ (Figures $1.4 \mathrm{~b}, 1.4 \mathrm{e}$ ) around the vertical axis are CAMC. The piecewise- $C^{\infty}$ closed surface (Figures 1.7d) given by rotating the closed curve $\left(C_{\gamma}\right)_{5}$ (Figure 1.4f) is not CAMC.

(a) The image of $\xi_{\gamma_{1}^{0}}$

(b) The image of $\xi_{\gamma_{1}}$

Figure 1.6: (a): The image of the Cahn-Hoffman map $\xi_{\gamma_{1}^{0}}: S^{1} \rightarrow \mathbb{R}^{2}$ for $\gamma_{1}^{0}: S^{1} \rightarrow \mathbb{R}_{>0}$. (b): The image of the Cahn-Hoffman map $\xi_{\gamma_{1}}: S^{2} \rightarrow \mathbb{R}^{3}$ for $\gamma_{1}: S^{2} \rightarrow \mathbb{R}_{>0}$.


Figure 1.7: Some of the closed surfaces which are subsets of $\xi_{\gamma_{1}}\left(S^{2}\right)$ for $\gamma_{1}$ defined by (5.5) (Figure 1.6b). They are surfaces given by rotating the curves $W_{\gamma},\left(C_{\gamma}\right)_{1},\left(C_{\gamma}\right)_{4}$, $\left(C_{\gamma}\right)_{5}$, respectively. (a), (b), (c): The anisotropic mean curvature for the outwardpointing normal is -1 . (d) : The anisotropic mean curvature is -1 on the 'outer part', while it is 1 on the 'inner part'. Hence, this surface is not CAMC.

Let us give another example. We rotate $\bar{\gamma}_{1}^{0}$ around the origin by $\pi / 4$ and obtain the following example.

$$
\begin{equation*}
\bar{\gamma}_{2}^{0}\left(\nu_{1}, \nu_{3}\right)=\frac{\nu_{1}^{6}+15 \nu_{1}^{4} \nu_{3}^{2}+15 \nu_{1}^{2} \nu_{3}^{4}+\nu_{3}^{6}}{4\left(\nu_{1}^{2}+\nu_{3}^{2}\right)^{5 / 2}} \tag{5.6}
\end{equation*}
$$

defined on the $\left(\nu_{1}, \nu_{3}\right)$-plane. We denote the restriction of $\bar{\gamma}_{2}^{0}$ to $S^{1}$ by $\gamma_{2}^{0}$. The CahnHoffman map $\xi_{\gamma_{2}^{0}}: S^{1} \rightarrow \mathbb{R}^{2}$ for $\gamma_{2}^{0}$ is obtained by rotating $\xi_{\gamma_{1}^{0}}$ around the origin by
$\pi / 4$, and so it is represented as follows (Figure 1.8a).

$$
\begin{align*}
\xi_{\gamma_{2}^{0}}(\nu)= & \frac{1}{4}\left(\nu_{1}\left(\nu_{1}^{6}-9 \nu_{1}^{4} \nu_{3}^{2}+15 \nu_{1}^{2} \nu_{3}^{4}+25 \nu_{3}^{6}\right), \nu_{3}\left(25 \nu_{1}^{6}+15 \nu_{1}^{4} \nu_{3}^{2}-9 \nu_{1}^{2} \nu_{3}^{4}+\nu_{3}^{6}\right)\right)  \tag{5.7}\\
= & \frac{1}{4}\left((\cos \theta)\left(\cos ^{6} \theta-9 \cos ^{4} \theta \sin ^{2} \theta+15 \cos ^{2} \theta \sin ^{4} \theta+25 \sin ^{6} \theta\right),\right. \\
& \left.(\sin \theta)\left(25 \cos ^{6} \theta+15 \cos ^{4} \theta \sin ^{2} \theta-9 \cos ^{2} \theta \sin ^{4} \theta+\sin ^{6} \theta\right)\right) \tag{5.8}
\end{align*}
$$

$\left(\nu=\left(\nu_{1}, \nu_{3}\right)=(\cos \theta, \sin \theta) \in S^{1}\right)$. The higher dimensional example obtained by rotating $\bar{\gamma}_{2}^{0}$ around the $\nu_{3}$-axis is given by

$$
\begin{equation*}
\bar{\gamma}_{2}\left(\nu_{1}, \nu_{2}, \nu_{3}\right)=\frac{\left(\nu_{1}^{2}+\nu_{2}^{2}\right)^{3}+15\left(\nu_{1}^{2}+\nu_{2}^{2}\right)^{2} \nu_{3}^{2}+15\left(\nu_{1}^{2}+\nu_{2}^{2}\right) \nu_{3}^{4}+\nu_{3}^{6}}{4\left(\nu_{1}^{2}+\nu_{2}^{2}+\nu_{3}^{2}\right)^{5 / 2}}, \quad\left(\nu_{1}, \nu_{2}, \nu_{3}\right) \in \mathbb{R}^{3} . \tag{5.9}
\end{equation*}
$$

The restriction $\gamma_{2}$ of $\bar{\gamma}_{2}: \mathbb{R}^{3} \rightarrow \mathbb{R}_{\geq 0}$ to $S^{2}$ can be written as

$$
\begin{equation*}
\gamma_{2}\left(\nu_{1}, \nu_{2}, \nu_{3}\right)=\left(\nu_{1}^{2}+\nu_{2}^{2}\right)^{3}+15\left(\nu_{1}^{2}+\nu_{2}^{2}\right)^{2} \nu_{3}^{2}+15\left(\nu_{1}^{2}+\nu_{2}^{2}\right) \nu_{3}^{4}+\nu_{3}^{6}, \quad\left(\nu_{1}, \nu_{2}, \nu_{3}\right) \in S^{2} . \tag{5.10}
\end{equation*}
$$

The corresponding Cahn-Hoffman map $\xi_{\gamma_{2}}: S^{2} \rightarrow \mathbb{R}^{3}$ is given as follows (Figure 1.8b).

$$
\begin{align*}
\xi_{\gamma_{2}}(\nu)= & \frac{1}{4}\left((\cos \theta)\left(\cos ^{6} \theta-9 \cos ^{4} \theta \sin ^{2} \theta+15 \cos ^{2} \theta \sin ^{4} \theta+25 \sin ^{6} \theta\right)(\cos \rho)\right. \\
& (\cos \theta)\left(\cos ^{6} \theta-9 \cos ^{4} \theta \sin ^{2} \theta+15 \cos ^{2} \theta \sin ^{4} \theta+25 \sin ^{6} \theta\right)(\sin \rho) \\
& \left.(\sin \theta)\left(25 \cos ^{6} \theta+15 \cos ^{4} \theta \sin ^{2} \theta-9 \cos ^{2} \theta \sin ^{4} \theta+\sin ^{6} \theta\right)\right) \tag{5.11}
\end{align*}
$$

$\left(\nu=(\cos \theta \cos \rho, \cos \theta \sin \rho, \sin \theta) \in S^{2}\right)$.
By the same way as above, we get closed piecewise- $C^{\infty}$ CAMC surfaces (Figure 1.9) and closed piecewise- $C^{\infty}$ non-CAMC surfaces for $\gamma_{2}$ which are subsets of $\xi_{\gamma_{2}}\left(S^{2}\right)$ and which are not the Wulff shape $W_{\gamma_{2}}$ (up to homothety and translation).

(a) The image of $\xi_{\gamma_{2}^{0}}$

(b) The image of $\xi_{\gamma_{2}}$

Figure 1.8: (a): The image of the Cahn-Hoffman map $\xi_{\gamma_{2}^{0}}: S^{1} \rightarrow \mathbb{R}^{2}$ for $\gamma_{2}^{0}: S^{1} \rightarrow \mathbb{R}_{>0}$. (b): The image of the Cahn-Hoffman map $\xi_{\gamma_{2}}: S^{2} \rightarrow \mathbb{R}^{3}$ for $\gamma_{2}: S^{2} \rightarrow \mathbb{R}_{>0}$.

### 1.6 Proofs of Theorems 1.1.1, 1.1.2

The functions $\gamma_{1}^{0}: S^{1} \rightarrow \mathbb{R}_{>0}, \gamma_{2}^{0}: S^{1} \rightarrow \mathbb{R}_{>0}$ and their rotations around the vertical axis give examples which prove Theorems 1.1.1, 1.1.2. In fact, Propositions 1.4.5, 1.5.1


Figure 1.9: Some of the closed surfaces which are subsets of $\xi_{\gamma_{2}}\left(S^{2}\right)$ for $\gamma_{2}$ defined by (5.10) (Figure 1.8b). The anisotropic mean curvature for the outward-pointing normal is -1 .
give suitable examples for $n=1,2$, respectively. Also, higher dimensional examples are obtained by the method given in the section 1.5.

### 1.7 Applications to anisotropic mean curvature flow: proofs of Theorems 1.1.4, 1.1.3

Let $\gamma: S^{n} \rightarrow \mathbb{R}_{>0}$ be of $C^{2}$ with Cahn-Hoffman map $\xi_{\gamma}$. Let $X_{t}: M \rightarrow \mathbb{R}^{n+1}$ be one-parameter family of embedded piecewise- $C^{2}$ hypersurfaces with anisotropic mean curvature $\Lambda_{t}$. Assume that the Cahn-Hoffman field $\tilde{\xi}_{t}$ along $X_{t}$ is defined on $M$. If $X_{t}$ satisfies $\partial X_{t} / \partial t=\Lambda_{t} \tilde{\xi}_{t}$, it is called an anisotropic mean curvature flow, which diminishes the anisotropic energy if $\Lambda_{t} \not \equiv 0$. In fact,

$$
\begin{align*}
\frac{d \mathcal{F}_{\gamma}\left(X_{t}\right)}{d t} & =-\int_{M} n \Lambda_{t}\left\langle\frac{\partial X_{t}}{\partial t}, \nu_{t}\right\rangle d A_{t}=-\int_{M} n \Lambda_{t}^{2}\left\langle D \gamma+\gamma\left(\nu_{t}\right) \nu_{t}, \nu_{t}\right\rangle d A_{t} \\
& =-\int_{M} n \Lambda_{t}^{2} \gamma\left(\nu_{t}\right) d A_{t}<0 \tag{7.1}
\end{align*}
$$

holds.
Proof of Theorem 1.1.4. Since the anisotropic mean curvature of $\xi_{\gamma}$ is $-1, \Lambda_{t}=-1 / \sqrt{2(c-t)}$ holds. On the other hand, $\widetilde{\xi}_{t}=\xi_{\gamma}$ holds. These two facts imply that (i) and (ii) hold.

Proof of Theorem 1.1.3. Examples stated in Propositions 1.4.5, 1.5.1 give the desired result.

## Chapter 2

# A discretization of the anisotropic energy for curves and surfaces 


#### Abstract

In this chapter, we study discrete curves and surfaces from the variational problem for the anisotropic energy. For the planar curve case, by extracting a notion of discrete curvature vector, we show that this vector can be used effectively for an unified interpretation of various discrete curvature notions, Steiner-type formula and the stability problem of critical points of the length functional. Moreover, we formulate discrete CAMC curves and give an application. For the surface case, we derive a generalization of the cotangent formula obtained in [34] and formulate discrete CAMC surfaces. We also visualize some examples and derive a stability problem which is a generalization of the result in [37].


### 2.1 Introduction

Discrete differential geometry is an active research field because of the interaction among the theory, the algorithm and the visualization on the computer. In this chapter, we mainly focus on discrete objects themselves, not approximation problems of smooth objects or convergence problems. In this field, from the theoretical point of view, one approach is based on the variational methods and the other focus on the integrability of equations. We take the former approach in this chapter and try to develop a theory of discrete curves and simplicial surfaces based on the variational problem for the anisotropic energy as we discussed in Chapter 1. One of the benefit of this approach is that we can consider not only a surface itself but also vector fields on the surface, thus it has many applications (see e.g. [35], [21], [36], [38]). Moreover, this approach with the anisotropic energy gives an unified formulation of discrete CMC surfaces in the Euclidean space, the Lorentz-Minkowski space and the isotropic space (see Appendix). On the other hand, a discrete surface theory using the integrable systems is originated
from [6]. The study of the relation between the integrable system approach and the variational approach has just begun recently, see e.g. [27].

In $\S 2.2$, we recall some basic notions from the (combinatorial) simplicial complexes and its geometric realization [35]. By a discrete curve or surface we mean a geometric realization of a combinatorial simplicial curve or surface, respectively. We first derive the first variation formula for the anisotropic energy and extract the vector from the formula which should be called "discrete curvature vector" in §2.3.1. In §2.3.2, by using this curvature vector, we show an unifed interpretation which derive various kinds of discrete curvature notions defined in [19], [15]. The important viewpoint here is that there is no natural notion of the line element on the vertices. In §2.3.3, we derive the Steiner-type formula for parallel curves by using the "vertex normal" constructed from the curvature vector derived in §2.3.1. Since there is no universal definition of the discrete curvature and our main interest is CAMC curves and surfaces, we will define discrete CAMC curves and surfaces without defining the curvature itself. Our definitions of discrete CAMC curves and surfaces are inspired by discrete CMC surfaces defined in [37]. In §2.3.4, we characterize closed discrete constant curvature curves as regular polygons which are possibly non-convex. In $\S 2.3 .5$, we formulate discrete CAMC curves and derive the conservation law for the Euler-Lagrange equation defining the discrete CAMC curve. Then, as an application of the conservation law, we give another proof for the Chakerian-Lange theorem [8]. In $\S 2.3 .6$, we will consider the stability problem of the discrete CAMC curves. For the isotropic case, we derive the second variation formula similar to the smooth case by decomposing the variation vectors into the vertex normal directions introduced in $\S 2.3 .3$ and "tangential" directions. Moreover we show the instability for the non-convex regular polygons in $\S 2.3 .7$ by using the normal variation and discrete Wirtinger's inequality [12].

Then, as one interpretation of this conservation law, we give another proof for the Chakerian-Lange theorem [8]. Moreover, we will give a "non-uniqueness example" for the discrete CAMC curves as in the Chapter 1. In $\S 2.3 .6$, we will consider the stability problem of the discrete CAMC curves. We derive a simple criteria for the instability for the discrete CAMC curves and visualize its energy descent deformation for some cases. In the isotropic case, we can derive the second variation formula similar to the smooth case by decomposing the variation vector field to the "normal" and "tangential" directions derived in the section. Moreover, we show the instability for the non-convex regular polygons by using the second variation formula for "normal" variations and discrete Wirtinger's inequality [12] in §2.3.7.

From $\S 2.4$, we treat the surface case. We first derive the first variation formula of the anisotropic energy for simplicial surfaces (Theorem 2.4.1) which generalizes the cotangent formula obtained in [34]. Then we define the discrete CAMC surfaces as in the curve case. In §2.4.3, we give many examples of discrete CAMC surfaces from various point of view. In $\S 2.4 .4$, we derive the second variation formula for the anisotropic energy and prove the stability of the very small part of the discrete CAMC surfaces if the energy density is of class $C^{2}$ and convex (Corollary 2.4.16). This result is a generalization of the result obtained in [37].

### 2.2 Preliminaries

We will develop a discretization theory for the anisotropic energy by using simplicial complexes. To do so, we need to recall basic notions from simplicial complexes (cf. [35]).

### 2.2.1 Abstract simplicial complexes

Definition 2.2.1 (abstract simplex). Let $V=\left\{v_{0}, \ldots, v_{m}\right\}$ be a set, where $v_{j}$ are not necessarily in the Euclidean space. Then we call the ordered set $\left[v_{0}, \ldots, v_{m}\right]$ be an abstract $m$-simplex or simply abstract simplex. The number $m$ is called the dimension of the abstract simplex. For the completeness, we admit the empty set $\emptyset$.

Remark . An abstract simplex is defined just as a set. For example, we consider a triangle $T=(p, q, r)$ in $\mathbb{R}^{2}$. This is the convex hull of 3 vertices $\{p, q, r\}$, and the edges and faces can be reconstructed from the vertices, i.e., these can be considered as an element of the powered set of vertices. For example, the edge $\overline{p q}$ corresponds to a set $\{p, q\} \subset\{p, q, r\}$.

Definition 2.2.2 (face). Let $\sigma=\left[v_{0}, \ldots, v_{m}\right]$ be an abstract $m$-simplex. Then we call the subsimplex consists of $(k+1)$ points $\left[v_{i_{0}}, \ldots, v_{l_{k}}\right]$ the $k$-face of $\sigma$. An $(m-1)$-face of $\sigma$ is called proper. For completeness, we consider the empty set $\emptyset=[]$ and the simplex itself as a face of $\sigma$.

Remark. Every abstract $m$-simplex has $2^{m+1}$ faces up to the order.
Example 2.2.3 (triangle (2-simplex)). An 2-simplex $\left[v_{0}, v_{1}, v_{2}\right]$ has 8 faces, that is, three 0 -simplices (vertices) $\left[v_{0}\right],\left[v_{1}\right],\left[v_{2}\right]$, three 1 -simplices (edges) $\left[v_{0}, v_{1}\right],\left[v_{1}, v_{2}\right],\left[v_{2}, v_{0}\right]$, the empty set [] and the triangle itself $\left[v_{0}, v_{1}, v_{2}\right.$ ].

Definition 2.2.4 ((abstract) simplicial complex). Let $V=\left\{v_{1}, v_{2}, \ldots\right\}$ be an abstract point set. An (abstract) simplicial complex $K=K(V)=\{\sigma, \tau, \ldots\}$ is a set of simplices formed by finite subset of $V$ and satisfies the following condition : if $\sigma \in K$ is a simplex, then any face $\tau$ of $\sigma$ is also an element of $K$.

If two simplices in $K$ share a common non-empty face, then they are called neighbors. The boundary of $K$ is formed by any proper face that belongs to only one simplex, and its faes. The boundary is also an abstract simplicial complex.

Example 2.2.5. Let us consider the simplest case. Let $V=\left\{v_{0}, v_{1}\right\}$ be a set. Then $K_{1}:=\left\{\left[v_{0}\right],\left[v_{0}, v_{1}\right]\right\}$ is not a simplicial complex since the face $\left[v_{1}\right] \subset\left[v_{0}, v_{1}\right]$ does not belong to $K_{1}$. If we consider $K:=K_{1} \cup\left\{\left[v_{1}\right]\right\}$, we have a simplicial complex.

### 2.2.2 Subsimplex, star and link

Definition 2.2.6 (subsimplex). Let $K$ be a simplicial complex. Then a subset $L \subset K$ is called the subcomplex of $K$ if $L$ itself is a simplicial complex.

Definition 2.2.7 ( $p$-skelton). We denote the set of all $p$-dimensional simplices by $K^{(p)}$ and call the $p$-skelton of $K$.
Example 2.2.8. For a simplicial complex $K, K^{(0)}$ is the set of vertices and $K^{(2)}$ is the set of triangles.

Definition 2.2.9 (star, link). Let $K$ be an simplicial complex and $\sigma \in K$ be a simplex. Then the star of $\sigma$ and the link of $\sigma$ is defined as follows:

$$
\begin{aligned}
\operatorname{star}(\sigma) & :=\{\eta \in K \mid \sigma \text { includes } \eta \text { and every face of } \eta\}, \\
\operatorname{link}(\sigma) & :=\{\eta \in \operatorname{star}(\sigma) \mid \eta \cap \sigma=\emptyset\}
\end{aligned}
$$

### 2.2.3 Simplicial map

Definition 2.2.10 (simplicial map). Let $K, L$ be simplicial complexes. Then a map $\varphi: K \rightarrow L$ is said to be the simplicial map if $\varphi$ satisfies the following conditions:
(1) $\left.\varphi\right|_{K^{0}}: K^{(0)} \rightarrow L^{(0)}$, i.e, each vertex of $K$ is mapped to a vertex of $L$.
(2) For any simplex $\sigma=\left[v_{0}, \ldots, v_{k}\right] \in K$, the image $\varphi(\sigma)=\left[\varphi\left(v_{0}\right), \ldots, \varphi\left(v_{k}\right)\right]$, which can be degenerate, is a simplex in $L$.
A simplicial map $\varphi: K \rightarrow L$ is said to be simplicial isomorphism if the map $\left.\varphi\right|_{K^{(p)}}$ : $K^{(p)} \rightarrow L^{(p)}$ is bijective for every $p$. In this case $K$ and $L$ are called simplicially isomorphic.

### 2.2.4 Simplicial curves, simplicial surfaces

Definition 2.2.11 (path, closed curve). Let $n$ be a non-negative integer. A standard $n$-path is a simplicial complex formed by
(1) $(n+1)$ abstract points : $V=\left\{v_{0}, \ldots, v_{n}\right\}$.
(2) the set of $n$ edges $e_{k}=\left[v_{k}, v_{k+1}\right], k=0, \ldots, n-1$.

A standard $n$-circle is the union of a standard $n$-path and the "final" edge $e_{n}=\left[v_{n}, v_{0}\right]$. We sometimes call a standard $n$-path (or standad $n$-circle) a simplicial curve and denote it by $G$. Similarly we sometimes consider $\mathbb{N}$ or $\mathbb{Z}$ as a standard path (non-compact case).
Definition 2.2.12 (abstract $m$-path). Let $K$ be a simplicial complex and $\delta$ be a standard $n$-path. Then a simplicial map $\varphi: \delta \rightarrow \varphi(\delta) \subset K$ (or its image $\varphi(\delta)$ ) is called the abstract $n$-path. We define the abstract $n$-circle similarly. An abstract $n$-circle is called simple if the map $\varphi$ is injective.
Definition 2.2 .13 (closed simplicial surface). A closed simplicial surface $\Sigma$ is an abstract simplicial complex formed by finite number of triangles and satisfying the following conditions:
(1) Any point $p \in \Sigma$ belongs to at least one triangle $T \in \Sigma$.
(2) For any point $p \in \Sigma$, the link $\operatorname{link}(p)$ is an abstract circle.

### 2.2.5 Geometric simplicial complexes

Our main object is defined by a geometric realization of an abstract simplicial complex. First we give a "parameter domain" of abstract simplices.

Definition 2.2.14 (standard simplex). A standard simplex $\Delta^{m} \subset \mathbb{R}^{m+1}$ is the convex hull of the end points of the canonical basis $\left\{e_{0}, \ldots, e_{m}\right\} \subset \mathbb{R}^{m+1}$. Formally $\Delta^{m}$ is defined by

$$
\Delta^{m}=\left\{\sum_{i=0}^{m} \lambda_{i} e_{i} \mid 0 \leq \lambda_{i} \leq 1, \sum_{i=0}^{m} \lambda_{i}=1\right\}
$$

Here the coefficient $\left(\lambda_{0}, \ldots, \lambda_{m}\right)$ is called the barycentric coordinates.
Definition 2.2.15 (geometric simplicial complex). A geometric simplicial complex $(P, K(V))$ consists of the following three objects:
(1) (Combinatorial structure) A simplicial complex $K(V)$ over a set of abstract points $V=\left\{v_{0}, v_{1}, \ldots\right\}$.
(2) (Geometric realization) A set of geometric vertices $P=\left\{p_{0}, p_{1}, \ldots\right\} \subset \mathbb{R}^{n}$ and a bijection $\Phi: V \rightarrow P, \Phi\left(v_{i}\right)=p_{i}$.
(3) (Local parametrization) For each $k$-simplex $\sigma=\left[v_{i_{0}}, \ldots, v_{i_{k}}\right]$ there exists a simplicial map

$$
\varphi_{\sigma}: \Delta^{k} \rightarrow \sigma, \quad \varphi_{\sigma}\left(e_{j}\right)=p_{i_{j}} .
$$

The set of affine maps $\left\{\varphi_{\sigma}\right\}_{\sigma \in K}$ gives a piecewise affine structure on the geometric simplicial complex $K$. This induced from the bijection $\Phi$ we don't denote explicitly in the pair $(P, K(V))$.

In the following sections we use the word "a simplicial curve (a discrete curve)" and "a simplicial surface (a discrete surface, a triangulated surface)" as a geometric realization of a simplicial curve and a simplicial surface, respectively. Note that the geometric realization gives the induced metric on the abstract simplicial curves or surfaces.

### 2.3 Geometry of discrete anisotropic curve energy

### 2.3.1 Anisotropic curve energy and its first variation

For a discrete curve there is a natural question:
What is the unit normal, curvature and line element at the vertices?
In this section we derive will derive the first variation formula for the anisotropic energy and extract the "curvature vector" at a vertex in order to approach to this question. However non-uniqueness of the line element at vertices gives different notions
of the discrete curvature (so-called "no free lunch" story), and we will show that various kinds of discrete curvature notions, e.g. in [19], [15], can be derived from this viewpoint.

Let $\gamma: S^{1} \rightarrow \mathbb{R}_{>0}$ be a continuous function and $\Gamma_{h}$ be a simplicial curve in $\mathbb{R}^{2}$, i.e., $\Gamma_{h}$ is a geometric realization of a standard $(n-1)$-path or $(n-1)$-circle $G$. The set of vertices of $G$ is denoted by $V$ and its image is denoted by $\left\{p_{k}\right\}_{k=0}^{n-1}$. For each oriented edge $e_{k}:=\left[p_{k}, p_{k+1}\right]$ of $\Gamma_{h}$ we can assign an unit normal vector

$$
\nu_{k}:=R\left(\frac{p_{k+1}-p_{k}}{l_{k}}\right):=R\left(\frac{p_{k+1}-p_{k}}{\left|p_{k+1}-p_{k}\right|}\right),
$$

where $R$ is the $\pi / 2$-rotation or $-\pi / 2$-rotation in $\mathbb{R}^{2}$. It does not matter whichever we choose but we choose the same $R$ for all $k$. Then the anisotropic energy of $\Gamma_{h}$ is defined


Figure 2.1: A part of a discrete curve with $-\pi / 2$-rotation and its variation.
by

$$
\mathscr{F}_{\gamma}\left(\Gamma_{h}\right):=\sum_{k} \gamma\left(\nu_{k}\right)\left|p_{k+1}-p_{k}\right|=\sum_{k} \gamma\left(\nu_{k}\right) l_{k}
$$

By the definition of simplicial curves if we fix the orientation of the unit normal, the anisotropic energy $\mathscr{F}_{\gamma}$ for a discrete curve can be considered as a function $\mathscr{F}_{\gamma}: \mathbb{R}^{2 n} \rightarrow$ $\mathbb{R}_{>0}$.

First we derive the first variation formula, that is, we compute the directional derivative in $\mathbb{R}^{2 n}$. In the following we assume that $\gamma$ is of class $C^{1}$. We consider a variation

$$
p_{k}(t)=p_{k}+t v_{k}+O\left(t^{2}\right), \quad k=0, \ldots, n-1
$$

where ${ }^{t} \vec{v}=\left({ }^{t} v_{0}, \ldots,{ }^{t} v_{n-1}\right) \in \mathbb{R}^{2 n}$ is the "variation vector field". If $p_{k}$ is a boundary point of $\Gamma_{h}$, then we assume $v_{k}=0$.

We would like to find a vector $\nabla \mathscr{F}_{\gamma} \in \mathbb{R}^{2 n}$ which satisfies

$$
\frac{d}{d t}_{\mid t=0} \mathscr{F}_{\gamma}=\left\langle\vec{v}, \nabla \mathscr{F}_{\gamma}\right\rangle_{\mathbb{R}^{2 n}}=\sum_{k}\left\langle v_{k}, \nabla_{p_{k}} \mathscr{F}_{\gamma}\right\rangle_{\mathbb{R}^{2}}
$$

where we write ${ }^{t} \nabla \mathscr{F}_{\gamma}=\left({ }^{t} \nabla_{p_{0}} \mathscr{F}_{\gamma}, \ldots,{ }^{t} \nabla_{p_{n-1}} \mathscr{F}_{\gamma}\right) \in \mathbb{R}^{2 n}$. In this setting we have the following proposition:


Proposition 2.3.1. Let $\gamma: S^{1} \rightarrow \mathbb{R}_{>0}$ be a $C^{1}$ function and $\Gamma_{h}=\left\{p_{k}\right\}_{k}$ be a discrete curve. Then at each interior vertex $p_{k}$ the gradient of the anisotropic energy with respect to variations of the vertices can be expressed

$$
\begin{equation*}
\nabla_{p_{k}} \mathscr{F}_{\gamma}=R\left(\xi_{\gamma}\left(\nu_{k}\right)-\xi_{\gamma}\left(\nu_{k-1}\right)\right), \tag{3.1}
\end{equation*}
$$

where $\xi_{\gamma}: S^{1} \rightarrow \mathbb{R}^{2}$ is the Cahn-Hoffman map for $\gamma$.
Remarks . By using this formula we have

$$
\delta \mathscr{F}_{\gamma}=\sum_{k}\left\langle R\left(\xi_{\gamma}\left(\nu_{k}\right)-\xi_{\gamma}\left(\nu_{k-1}\right)\right), v_{k}\right\rangle=-\sum_{k}\left\langle\frac{1}{L_{k}} R\left(\xi_{\gamma}\left(\nu_{k-1}\right)-\xi_{\gamma}\left(\nu_{k}\right)\right), v_{k}\right\rangle L_{k},
$$

where we inserted some auxilary function $L: V_{n} \rightarrow \mathbb{R}_{>0}$. This kind of observation is already remarked in the paper [9]. We give some remarks for the isotropic case, i.e, $\gamma \equiv 1$ case.
(1) In this case, we have $\xi_{\gamma}(\nu)=\nu$ and the vector

$$
\widetilde{N}_{k}:=-\frac{1}{L_{k}} R\left(\nu_{k}-\nu_{k-1}\right)=\frac{1}{L_{k}}\left(\frac{p_{k+1}-p_{k}}{\left|p_{k+1}-p_{k}\right|}+\frac{p_{k-1}-p_{k}}{\left|p_{k-1}-p_{k}\right|}\right)
$$

can be considered as a curvature vector on the vertex $p_{k}$ because of the following reasons. First, the vector $\widetilde{N}_{k}$ is independent of the orientation of the unit normal. Moreover, if we assume that $X: I(\subset \mathbb{R}) \rightarrow \mathbb{R}^{2}$ is an arclength parametrized smooth regular curve and $\nu_{X}: I \rightarrow S^{1}$ is the (outward-pointing) unit normal, denoted by $\nu_{X}(s)=(\cos \theta(s), \sin \theta(s))$, the curvature $\kappa$ is characterized by

$$
\kappa=-\frac{d \theta}{d s} .
$$

Then we can regard the norm $\left|\tilde{N}_{k}\right|$ as a discretization of the above equation. From another point of view, we can regard the relation

$$
\frac{1}{L_{k}}\left(T_{k}-T_{k-1}\right)=\tilde{N}_{k}, \quad T_{k}:=\frac{p_{k+1}-p_{k}}{\left|p_{k+1}-p_{k}\right|}
$$

as a discrete version of the Frenet-Serret formula. As we will remark later, there are many possibilities of $L_{k}$, so we will define constant (anisotropic) curvature curves without defining the curvature notion. However, since the scaling factor $L_{k}$ does not change the direction of $-\nabla_{p_{k}}$ Length, we take a position that the vector $-\nabla_{p_{k}}$ Length gives the normal direction at the vertex $p_{k}$ throughout this chapter.
(2) In the length gradient case, we can give a geometric interpretation. If we fix the vertex $p_{k+1}$ and move the vertex $p_{k}$ and consider the length of the edge $e_{k}=\left[p_{k}, p_{k+1}\right]$. Then the direction which the length of the edge $l_{k}$ most decreases is $\left(p_{k+1}-p_{k}\right) /\left|p_{k+1}-p_{k}\right|$. The vector $-\nabla_{p_{k}}$ Length is a sum of two vectors $\left(p_{k+1}-p_{k}\right) /\left|p_{k+1}-p_{k}\right|$ and $\left(p_{k-1}-p_{k}\right) /\left|p_{k-1}-p_{k}\right|$ means the steepest direction of the length function at $p_{k}$.

Before giving the proof, we derive the first variation of the unit normal.
Lemma 2.3.2. $l_{k} \delta \nu_{k}=R\left(v_{k+1}-v_{k}\right)-\left\langle R\left(v_{k+1}-v_{k}\right), \nu_{k}\right\rangle \nu_{k}$.
Remark. If we define the notation $\nabla v_{k}=\left(v_{k+1}-v_{k}\right) / l_{k}$, then

$$
\delta \nu_{k}=\nabla\left(R v_{k}\right)-\left\langle\nabla\left(R v_{k}\right), \nu_{k}\right\rangle \nu_{k}
$$

and this can be considered as the "covariant derivative" of $R v$ at the edge $e_{k}=\left[p_{k}, p_{k+1}\right]$. Proof. By using the fact

$$
\delta l_{k}^{2}=\delta\left|p_{k+1}-p_{k}\right|^{2}=2\left\langle p_{k+1}-p_{k}, v_{k+1}-v_{k}\right\rangle, \quad \delta l_{k}^{2}=2 l_{k} \delta l_{k},
$$

we have $\delta l_{k}=\left\langle-R \nu_{k}, v_{k+1}-v_{k}\right\rangle=\left\langle\nu_{k}, R\left(v_{k+1}-v_{k}\right)\right\rangle$ and therefore

$$
\begin{aligned}
\delta \nu_{k} & =\delta R\left(\frac{p_{k+1}-p_{k}}{l_{k}}\right)=R\left(\frac{v_{k+1}-v_{k}}{l_{k}}-\frac{p_{k+1}-p_{k}}{l_{k}^{2}} \delta l_{k}\right) \\
& =\frac{1}{l_{k}}\left(R\left(v_{k+1}-v_{k}\right)-\nu_{k}\left\langle\nu_{k}, R\left(v_{k+1}-v_{k}\right)\right\rangle\right)
\end{aligned}
$$

Proof of Proposition 2.3.1. By the chain rule and the Leibniz rule we have

$$
\delta \mathscr{F}_{\gamma}:=\frac{d}{d t}_{\mid t=0} \mathscr{F}_{\gamma}=\sum_{k}\left\langle D \gamma\left(\nu_{k}\right), \delta \nu_{k}\right\rangle l_{k}+\gamma\left(\nu_{k}\right) \delta l_{k} .
$$

By using the previous lemma and the fact $D \gamma(\nu) \perp \nu$, we have

$$
\begin{aligned}
\delta \mathscr{F}_{\gamma} & =\sum_{k}\left\langle D \gamma\left(\nu_{k}\right), R\left(v_{k+1}-v_{k}\right)\right\rangle+\gamma\left(\nu_{k}\right)\left\langle\nu_{k}, R\left(v_{k+1}-v_{k}\right)\right\rangle \\
& =\sum_{k}\left\langle\xi_{\gamma}\left(\nu_{k}\right), R\left(v_{k+1}-v_{k}\right)\right\rangle=\sum_{k}\left\langle R\left(\xi_{\gamma}\left(\nu_{k}\right)-\xi_{\gamma}\left(\nu_{k-1}\right)\right), v_{k}\right\rangle .
\end{aligned}
$$

Remark. We can consider the length gradient for non-manifold case, i.e.,

$$
\nabla_{p} \text { Length }=-\sum_{k=1}^{m} \frac{p_{k}-p}{\left|p_{k}-p\right|} .
$$

For example, when $m=3$ (trivalent graph), $\nabla_{p}$ Length $=0$ if and only if the angles of adjacent edges are $2 \pi / 3$ ( 120 degree).

### 2.3.2 Relation with other notions of the discrete curvature

In the lecture note by Hoffman [19], three kinds of notions of the curvature for discrete curves are defined:
(1) The curvature at vertices by using the vertex osculating circle method,
(2) The curvature at edges by using the edge osculating circle method,
(3) The curvature at vertices by using edge osculating circle for "arclength parametrized" curve.

Moreover, Hatakeyama [15] also defined the curvature for discrete curves another way. We will show that these curvature notions can be derived from our curvature vector by taking the different line elements $L_{k}$ at vertices. In other words, these difference comes from the non-uniqueness of the "natural" line element at vertices.

As we have seen before, from the first variation formula

$$
\delta \text { Length }=\sum_{k}\left\langle R\left(\nu_{k}-\nu_{k-1}\right), v_{k}\right\rangle=-\sum_{k}\left\langle\frac{1}{L_{k}} R\left(\nu_{k-1}-\nu_{k}\right), v_{k}\right\rangle L_{k},
$$

we shall consider the vector

$$
\widetilde{N}_{k}=\frac{1}{L_{k}} R\left(\nu_{k-1}-\nu_{k}\right)=\frac{1}{L_{k}}\left(\frac{p_{k+1}-p_{k}}{\left|p_{k+1}-p_{k}\right|}+\frac{p_{k-1}-p_{k}}{\left|p_{k-1}-p_{k}\right|}\right)
$$

as the curvature vector on the vertex $p_{k}$. However, there are many possibilities of $L_{k}$. Therefore we call $\widetilde{N}_{k}$ the discrete curvature vector with respect to $L_{k}$. For example, the choice $2 L_{k}=l_{k+1}+l_{k}=\left|p_{k+1}-p_{k}\right|+\left|p_{k-1}-p_{k}\right|$, which satisfies $\sum_{k} L_{k}=\operatorname{Length}\left(\Gamma_{h}\right)$, is compatible with the curve-shortening flow.

To describe the curvature notions, we have to define the angles at vertices. If we define (the absolute value of) the angle between $\nu_{k-1}$ and $\nu_{k}$ as $\theta_{k}$, i.e.,

$$
\cos \theta_{k}:=\left\langle\nu_{k}, \nu_{k-1}\right\rangle=\frac{\left\langle p_{k+1}-p_{k}, p_{k}-p_{k-1}\right\rangle}{\left|p_{k+1}-p_{k}\right| \cdot\left|p_{k}-p_{k-1}\right|} .
$$

In the following we assume $\cos \theta_{k} \neq 1$ for every $k$. We have to care about the signature of $\theta_{k}$. Let $R_{\theta}$ be the $\theta$-rotation in $\mathbb{R}^{2}$. We assign the signature depends on the choice of the rotation $R$ :

$$
\sigma:= \begin{cases}+1 & \text { if } \quad R=R_{\pi / 2} \\ -1 & \text { if } \quad R=R_{-\pi / 2}\end{cases}
$$

In this situation the signature of $\theta_{k}$ is determined by the equation $R_{\sigma \theta_{k}}\left(\nu_{k-1}\right)=\nu_{k}$.
Remark. For any closed curve we have $\sigma \sum_{k} \theta_{k}=2 m \pi$ for some interger $m$. Because if we note that

$$
R_{\sigma\left(\theta_{0}+\cdots+\theta_{n-1}\right)} \nu_{0}=\nu_{n}=\nu_{0},
$$

then we have $\sigma\left(\theta_{0}+\cdots+\theta_{n-1}\right)=2 m \pi$.


Figure 2.2: The angle at the vertex $p_{k}$ ( $\theta_{k}<0$ in this figure).

By virtue of the fact

$$
\left|R\left(\nu_{k-1}-\nu_{k}\right)\right|=2 \sin \left(\left|\theta_{k}\right| / 2\right),
$$

we call the value

$$
\kappa\left(p_{k}\right):=\frac{2 \sin \left(\theta_{k} / 2\right)}{L_{k}}
$$

the discrete curvature at $p_{k}$ with respect to $L_{k}$.
Proposition 2.3.3 (The vertex osculating circle method [19]). If we choose

$$
L_{k}=\frac{\left|p_{k+1}-p_{k-1}\right|}{2 \cos \left(\theta_{k} / 2\right)}=\frac{\left|p_{k+1}-p_{k}+p_{k}-p_{k-1}\right|}{2 \cos \left(\theta_{k} / 2\right)}
$$

then the discrete curvature with respect to $L_{k}$ becomes

$$
\kappa\left(p_{k}\right)=\frac{2 \sin \theta_{k}}{\left|p_{k+1}-p_{k-1}\right|}=\frac{2 \sin \theta_{k}}{\left|p_{k+1}-p_{k}+p_{k}-p_{k-1}\right|}
$$

and this value coincides with the curvature based on the vertex osculating circle method.
Proposition 2.3.4 (For the arclength parametrized curves [19]). Assume $l_{k}=l_{k-1}=$ $l_{0}$. Then if we choose

$$
L_{k}=l_{0} \cos \frac{\theta_{k}}{2}=\frac{l_{k}+l_{k-1}}{2} \cdot \cos \frac{\theta_{k}}{2}
$$

then the discrete curvature with respect to $L_{k}$ becomes

$$
\kappa\left(p_{k}\right)=\frac{2}{l_{0}} \tan \frac{\theta_{k}}{2}
$$

and this value coincides with the curvature of arclength parametrized curve.
In the paper [15], the discrete curvature at the vertex is defined as

$$
\kappa\left(p_{k}\right):=\frac{1}{\left|p_{k}-p_{k-1}\right|}\left|\frac{p_{k+1}-p_{k}}{\left|p_{k+1}-p_{k}\right|}-\frac{p_{k}-p_{k-1}}{\left|p_{k}-p_{k-1}\right|}\right|=-\frac{\mid \nabla_{p_{k} \text { Length } \mid}^{\left|p_{k}-p_{k-1}\right|} . . . ~}{\text {. }}
$$

Then we immediately have the following result:

Proposition 2.3.5. If we choose $L_{k}=l_{k-1}=\left|p_{k}-p_{k-1}\right|$, then the discrete curvature with respect to $L_{k}$ coincides with the discrete curvature defined by Hatakeyama [15].

Before considering the edge osculating circle method, we shall modify the first variation formula from the vertex-based expression to the edge-based expression. If we put $v_{k}=\left(w_{k}+w_{k-1}\right) / 2$, then we have

$$
\begin{aligned}
\delta \text { Length } & =\frac{1}{2} \sum_{k}\left\langle\nabla_{p_{k}} \text { Length, } w_{k}+w_{k-1}\right\rangle \\
& =\frac{1}{2} \sum_{k}\left\langle\nabla_{p_{k}} \text { Length }+\nabla_{p_{k+1}} \text { Length, } w_{k}\right\rangle=-\sum_{k}\left\langle\frac{R\left(\nu_{k-1}-\nu_{k+1}\right)}{2 L_{k}^{\prime}}, w_{k}\right\rangle L_{k}^{\prime}
\end{aligned}
$$

where $L_{k}^{\prime}$ is some auxiliary function. As in the vertex case, we call the value

$$
\kappa\left(e_{k}\right):=\frac{1}{L_{k}^{\prime}} \cdot \sin \frac{\theta_{k}+\theta_{k+1}}{2}
$$

the discrete curvature at the edge $e_{k}=\left[p_{k}, p_{k+1}\right]$ with respect to $L_{k}^{\prime}$.
Proposition 2.3.6 (The edge osculating circle method [19]). If we choose

$$
L_{k}^{\prime}=l_{k} \cos \frac{\theta_{k}}{2} \cos \frac{\theta_{k+1}}{2}=\left|p_{k+1}-p_{k}\right| \cos \frac{\theta_{k}}{2} \cos \frac{\theta_{k+1}}{2},
$$

then the discrete curvature with respect to $L_{k}^{\prime}$ becomes

$$
\kappa\left(e_{k}\right)=\frac{\tan \left(\theta_{k} / 2\right)+\tan \left(\theta_{k+1} / 2\right)}{\left|p_{k+1}-p_{k}\right|}
$$

and this value coincides with the curvature based on the edge osculating circle method.
Remark. To define the discrete curvature, we have to choose $L_{k}$ (respectively $L_{k}^{\prime}$ ) properly. That means if $l_{k}, l_{k-1} \rightarrow d s$ and $\theta_{k} \rightarrow 0$, then $L_{k}$ (respectively $L_{k}^{\prime}$ ) must converge to $d s$, i.e., $L_{k}$ must be a "good" candidate for a discrete line element. We can check that $L_{k}$ and $L_{k}^{\prime}$ satisfy this condition in the above examples.

Remark ("No free lunch" for the discrete Laplacian, cf. [41]). We consider these kinds of "no free lunch" story for the discrete Laplacian which will be used in the second variation formula. Let $V$ be the vertices of a standard $n$-path and $X: V \rightarrow \mathbb{R}^{2}$ be its geometric realization (hence we have the induced metric). Then for a function $\psi: V \rightarrow \mathbb{R}^{d}$, the gradient and the Laplacian of $\psi$ can be defined as

$$
\nabla \psi_{k}:=\frac{\psi_{k+1}-\psi_{k}}{l_{k}}, \quad \Delta \psi_{k}:=\frac{1}{L_{k}}\left(\nabla \psi_{k}-\nabla \psi_{k-1}\right)=\frac{1}{L_{k}}\left(\frac{\psi_{k+1}-\psi_{k}}{l_{k}}-\frac{\psi_{k}-\psi_{k-1}}{l_{k-1}}\right),
$$

where we denote $\psi_{k}:=\psi(k)$. Note that the gradient is the "edge-based operator" but the Laplacian is the "vertex-based" operator, in addition, the discrete curvature vector $\widetilde{N}_{k}$ with respect to $L_{k}$ can be written as $\Delta p_{k}$.

From another point of view, if we define the Dirichlet energy of $\psi$ as

$$
E_{h}(\psi):=\frac{1}{2} \sum_{k}\left|\nabla \psi_{k}\right|^{2} l_{k}=\frac{1}{2} \sum_{k} \frac{\left|\psi_{k+1}-\psi_{k}\right|^{2}}{l_{k}}
$$

then the first variation of the energy becomes

$$
\delta E_{h}(\psi)=\sum_{k} \frac{\left\langle\psi_{k+1}-\psi_{k}, \varphi_{k+1}-\varphi_{k}\right\rangle}{l_{k}}=\sum_{k}\left\langle\nabla \psi_{k-1}-\nabla \psi_{k}, \varphi_{k}\right\rangle=-\sum_{k}\left\langle\Delta \psi_{k}, \varphi_{k}\right\rangle L_{k},
$$

where we take the variation of $\psi$ as $\psi_{k}(t)=\psi_{k}+t \varphi_{k}+O\left(t^{2}\right)$. Therefore $\delta E_{h}(\psi)=0$ if and only if $\Delta \psi_{k}=0$. Note that the condition $\Delta \psi_{k}=0$ is independent of the choice of $L_{k}$.

As in the curvature case, the Laplacian can be changed since there is no natural "line element divisor $L_{k}$ ". However, with another function $\varphi: V \rightarrow \mathbb{R}$, we still have the following properties since the quanties $\Delta \psi_{k} L_{k}$ are independent of $L_{k}$ :
(1) If $\psi$ is constant, then $\Delta \psi=0$.
(2) The condition $\Delta \psi=0$ is independent of the choice of $L_{k}$ and in this case we have the mean value property:

$$
\psi_{k}=\frac{l_{k-1}}{l_{k}+l_{k-1}} \psi_{k+1}+\frac{l_{k}}{l_{k}+l_{k-1}} \psi_{k-1} .
$$

(3) $L^{2}$ symmetric property:

$$
\sum_{k} \psi_{k} \cdot \Delta \varphi_{k} \cdot L_{k}=\sum_{k} \Delta \psi_{k} \cdot \varphi_{k} \cdot L_{k} .
$$

Note that the summation is vertex-based.
(4) Integration by parts:

$$
-\sum_{k} \psi_{k} \cdot \Delta \varphi_{k} \cdot L_{k}=\sum_{k} \nabla \psi_{k} \cdot \nabla \varphi_{k} \cdot l_{k}
$$

Note that the right hand side is the edge-based summation but the left hand side is the vertex-based summation. As a corollary, the operator $-\Delta$ is positive semi-definite.

### 2.3.3 Parallel curves and Steiner-type formula

In this section we will derive the discrete version of the Steiner-type formula. The following type of Steiner formula is essentially appeared in some papers, for example [10]. Although they try to extract the curvature notion from the Steiner-type formula, we
will derive the Steiner-type formula by using our vertex normal and connect with the well-known curvature notion in [19].

Let $\Gamma_{h}=\left\{p_{k}\right\}_{k}$ be a discrete curve and take an interior vertex $p_{k}$. If we recall the length gradient

$$
-\nabla_{p_{k}} \text { Length }=R\left(\nu_{k-1}-\nu_{k}\right)=\frac{p_{k+1}-p_{k}}{\left|p_{k+1}-p_{k}\right|}+\frac{p_{k-1}-p_{k}}{\left|p_{k-1}-p_{k}\right|}
$$

gives a candidate of the normal direction at the vertex $p_{k}$ and the fact $\left|R\left(\nu_{k-1}-\nu_{k}\right)\right|=$ $2 \sin \left(\left|\theta_{k}\right| / 2\right)$, we should divide $-\nabla_{p_{k}}$ Length by $2 \sin \left(\theta_{k} / 2\right)$ in order to get the unit normal. However this choice does not work well. In the discrete case, we should consider another factor $2 \sin \left(\theta_{k} / 2\right) \cos \left(\theta_{k} / 2\right)=\sin \theta_{k}$ and put

$$
N_{k}:=-\frac{\nabla_{p_{k}} \text { Length }}{\sin \theta_{k}}=\frac{R\left(\nu_{k-1}-\nu_{k}\right)}{\sin \theta_{k}}=\frac{1}{1+\cos \theta_{k}}\left(\nu_{k}+\nu_{k-1}\right),
$$

if $\theta_{k} \neq 0$. Then we shall call the vector $N_{k}$ as the vertex normal at the vertex $p_{k}$. Then we consider the following deformation of a discrete curve

$$
p_{k}(t)=p_{k}+t N_{k}=p_{k}+t \cdot \frac{R\left(\nu_{k-1}-\nu_{k}\right)}{\sin \theta_{k}},
$$

where the vertex $p_{k}$ with $\theta_{k} \neq 0$.
Lemma 2.3.7. Assume $\theta_{k}, \theta_{k+1} \neq 0$. Then we have $\left\langle p_{k+1}(t)-p_{k}(t), \nu_{k}\right\rangle=0$, therefore we call this deformation parallel curves.

Proof. In order to clarify the situation, we consider a more general deformation

$$
N_{k}:=-\frac{\nabla_{p_{k}} \text { Length }}{M_{k}}=\frac{R\left(\nu_{k-1}-\nu_{k}\right)}{M_{k}},
$$

where $M_{k}$ is some auxilary function. Then we have

$$
\begin{aligned}
\left\langle p_{k+1}(t)-p_{k}(t), \nu_{k}\right\rangle & =\left\langle\left(p_{k+1}-p_{k}\right)+t\left(N_{k+1}-N_{k}\right), \nu_{k}\right\rangle \\
& =t\left(\frac{\left\langle R\left(\nu_{k}-\nu_{k+1}\right), \nu_{k}\right\rangle}{M_{k+1}}-\frac{\left\langle R\left(\nu_{k-1}-\nu_{k}\right), \nu_{k}\right\rangle}{M_{k}}\right) \\
& =t\left(-\frac{\left\langle R \nu_{k+1}, \nu_{k}\right\rangle}{M_{k+1}}-\frac{\left\langle R \nu_{k-1}, \nu_{k}\right\rangle}{M_{k}}\right)=t\left(\frac{\sin \theta_{k+1}}{M_{k+1}}-\frac{\sin \theta_{k}}{M_{k}}\right) .
\end{aligned}
$$

Therefore if we choose $M_{k}=\sin \theta_{k}$, then $\left\langle p_{k+1}(t)-p_{k}(t), \nu_{k}\right\rangle=0$.
Theorem 2.3.8 (Discrete Steriner-type formula). For parallel curves $\left\{p_{k}(t)\right\}_{k}$, we have

$$
\left|p_{k+1}(t)-p_{k}(t)\right|=\left|p_{k+1}-p_{k}\right|\left(1+t \cdot \kappa\left(e_{k}\right)\right),
$$

where $\kappa\left(e_{k}\right)$ is the discrete curvature based on the edge osculating circle method [19]:

$$
\kappa\left(e_{k}\right)=\frac{\tan \left(\theta_{k} / 2\right)+\tan \left(\theta_{k+1} / 2\right)}{\left|p_{k+1}-p_{k}\right|} .
$$



Figure 2.3: Parallel curves

Before giving the proof, we give an intuitive explanation for a special case. In the Figure 2.3, the similarity ratio of the triangles gives

$$
\begin{aligned}
& \left|p_{k+1}(t)-p_{k}(t)\right|:\left|p_{k+1}-p_{k}\right|=\left(r_{k}+t\right): r_{k} \\
\Longleftrightarrow & \left|p_{k+1}(t)-p_{k}(t)\right|=\left|p_{k+1}-p_{k}\right|\left(1+t \cdot \frac{1}{r_{k}}\right) \\
\Longleftrightarrow & \left|p_{k+1}(t)-p_{k}(t)\right|=\left|p_{k+1}-p_{k}\right|\left(1+t \cdot \kappa\left(e_{k}\right)\right)
\end{aligned}
$$

Proof. As in the previous lemma we consider $N_{k}:=R\left(\nu_{k-1}-\nu_{k}\right) / M_{k}$ and $p_{k}(t)=$ $p_{k}+t N_{k}$. Then

$$
\begin{aligned}
\left\langle\frac{p_{k+1}-p_{k}}{\left|p_{k+1}-p_{k}\right|}, N_{k+1}-N_{k}\right\rangle & =\left\langle-R \nu_{k}, \frac{R\left(\nu_{k}-\nu_{k+1}\right)}{M_{k+1}}-\frac{R\left(\nu_{k-1}-\nu_{k}\right)}{M_{k}}\right\rangle \\
& =-\left(\frac{1-\cos \theta_{k+1}}{M_{k+1}}+\frac{1-\cos \theta_{k}}{M_{k}}\right)
\end{aligned}
$$

$$
\begin{aligned}
\left|N_{k+1}-N_{k}\right|^{2} & =\frac{1}{\cos ^{2}\left(\theta_{k+1} / 2\right)}+\frac{1}{\cos ^{2}\left(\theta_{k} / 2\right)}-2\left\langle N_{k}, N_{k+1}\right\rangle \\
& =2+\tan ^{2} \frac{\theta_{k+1}}{2}+\tan ^{2} \frac{\theta_{k}}{2}-2 \cdot \frac{\sin \theta_{k} \sin \theta_{k+1}-\left(1-\cos \theta_{k}\right)\left(1-\cos \theta_{k+1}\right)}{M_{k} M_{k+1}} .
\end{aligned}
$$

If we take $M_{k}=\sin \theta_{k}$, then we have

$$
\begin{aligned}
\left\langle\frac{p_{k+1}-p_{k}}{\left|p_{k+1}-p_{k}\right|}, N_{k+1}-N_{k}\right\rangle & =-\left(\tan \frac{\theta_{k}}{2}+\tan \frac{\theta_{k+1}}{2}\right), \\
\left|N_{k+1}-N_{k}\right|^{2} & =\left(\tan \frac{\theta_{k}}{2}+\tan \frac{\theta_{k+1}}{2}\right)^{2}
\end{aligned}
$$

Therefore we conclude

$$
\begin{aligned}
& \left|p_{k+1}(t)-p_{k}(t)\right|^{2} \\
= & \left|p_{k+1}-p_{k}\right|^{2}-2 t\left\langle p_{k+1}-p_{k}, N_{k+1}-N_{k}\right\rangle+t^{2}\left|N_{k+1}-N_{k}\right|^{2} \\
= & \left|p_{k+1}-p_{k}\right|^{2}+2 t\left|p_{k+1}-p_{k}\right|\left(\tan \frac{\theta_{k}}{2}+\tan \frac{\theta_{k+1}}{2}\right)+t^{2}\left(\tan \frac{\theta_{k}}{2}+\tan \frac{\theta_{k+1}}{2}\right)^{2} \\
= & \left|p_{k+1}-p_{k}\right|^{2}\left(1+t \cdot \frac{\tan \left(\theta_{k} / 2\right)+\tan \left(\theta_{k+1} / 2\right)}{\left|p_{k+1}-p_{k}\right|}\right)^{2}
\end{aligned}
$$

Remark. This formula itself is already appeared some papers to extract the discrete curvature at the vertices by using the offsets, see [5], [10]. If we take the offset of the edges, then we have the following (at least) three possiblities. By computing the length

(a) $\Gamma_{h, t}^{(1)}$, connect by a segment

(b) $\Gamma_{h, t}^{(2)}$, connect by an arc

(c) $\Gamma_{h, t}^{(3)}$, connect by a wedge

Figure 2.4: Three possibilities to construct a new curve
of dotted curves in the figure, we can write the total length of each offsets as follows:

$$
\begin{aligned}
\operatorname{Length}\left(\Gamma_{h, t}^{(1)}\right) & =\operatorname{Length}\left(\Gamma_{h}\right)+t \sum_{k} 2 \sin \frac{\theta_{k}}{2}, \quad \operatorname{Length}\left(\Gamma_{h, t}^{(2)}\right)=\operatorname{Length}\left(\Gamma_{h}\right)+t \sum_{k} \theta_{k} \\
\operatorname{Length}\left(\Gamma_{h, t}^{(3)}\right) & =\operatorname{Length}\left(\Gamma_{h}\right)+t \sum_{k} 2 \tan \frac{\theta_{k}}{2}
\end{aligned}
$$

Note that $\theta \approx 2 \sin (\theta / 2) \approx 2 \tan (\theta / 2)$ if $\theta \approx 0$. The second curve $\Gamma_{h}^{(2)}$ is nothing but the normal cone method (or the boundary of the Minkowski sum with the disk) known in the convex geometry. However, the only possible way which does not change the number of the vertices during the offset procedure is the third one, and by modifying the third formula gives our Steiner-type formula:

$$
\begin{aligned}
\operatorname{Length}\left(\Gamma_{h, t}^{(3)}\right) & =\sum_{k}\left(l_{k}+2 t \tan \left(\theta_{k} / 2\right)\right) \\
& =\sum_{k}\left(l_{k}+t\left(\tan \left(\theta_{k} / 2\right)+\tan \left(\theta_{k+1} / 2\right)\right)\right)=\sum_{k} l_{k}\left(1+t \kappa\left(e_{k}\right)\right),
\end{aligned}
$$

where we put $\kappa\left(e_{k}\right)=\left(\tan \left(\theta_{k} / 2\right)+\tan \left(\theta_{k+1} / 2\right)\right) / l_{k}$ (the curvature defined by the edge osculating circle).

### 2.3.4 Critical points of the length functional

In $\S 2.3 .2$, we saw that there is no universal notion of discrete curvature. However, since our main interest is constant (mean) curvature objects, we formulate such objects by using the variational method without defining the curvature itself. In this section we focus on the length functional case and consider closed discrete curves in the following. The 2-dimensional oriented volume (i.e., the area) bounded by a closed discrete curve $\Gamma_{h}$ is defined by

$$
\operatorname{Vol}\left(\Gamma_{h}\right):=\frac{1}{2} \sum_{k}\left\langle p_{k}, \nu_{k}\right\rangle\left|p_{k+1}-p_{k}\right|=\frac{1}{2} \sum_{k}\left\langle p_{k}, R p_{k+1}\right\rangle,
$$

Note that the signature of the volume can be negative. Then we have the following result:

Lemma 2.3.9. For any vertex $p_{k}$ of $\Gamma_{h}$ the gradient of the volume $\nabla_{p_{k}}$ Vol is given by

$$
\nabla_{p_{k}} \mathrm{Vol}=\frac{1}{2} R\left(p_{k+1}-p_{k-1}\right)
$$

Proof. We consider any variation ${ }^{t} \vec{v}=\left({ }^{t} v_{1}, \ldots,{ }^{t} v_{n}\right)$ and it follows from the second representation of the volume that

$$
\begin{aligned}
\delta \mathrm{Vol} & =\frac{1}{2} \sum_{k} \delta\left\langle p_{k}, R p_{k+1}\right\rangle=\frac{1}{2} \sum_{k}\left(\left\langle v_{k}, R p_{k+1}\right\rangle+\left\langle p_{k}, R v_{k+1}\right\rangle\right) \\
& =\frac{1}{2} \sum_{k}\left(\left\langle v_{k}, R p_{k+1}\right\rangle+\left\langle p_{k-1}, R v_{k}\right\rangle=\frac{1}{2} \sum_{k}\left\langle R\left(p_{k+1}-p_{k-1}\right), v_{k}\right\rangle .\right.
\end{aligned}
$$

Remark (Another "no free lunch" story). We can modify the first variation formula of the volume as follows:

$$
\delta \mathrm{Vol}=\frac{1}{2} \sum_{k}\left\langle R\left(p_{k+1}-p_{k}+p_{k}-p_{k-1}\right), v_{k}\right\rangle=\sum_{k}\left\langle\frac{l_{k} \nu_{k}+l_{k-1} \nu_{k-1}}{2 L_{k}}, v_{k}\right\rangle L_{k} .
$$

At a glance, it seems like a natural to choose $2 L_{k}=l_{k}+l_{k-1}=\operatorname{Length}(\operatorname{star}(p))$ and this is also frequently used as a "vertex normal" (weighted sum of the edge normals):

$$
N_{k}^{V}:=\frac{l_{k} \nu_{k}+l_{k-1} \nu_{k-1}}{l_{k}+l_{k-1}} .
$$

In addition, we have

$$
-\nabla_{p_{k}} \text { Length }=R\left(\nu_{k-1}-\nu_{k}\right)=\frac{\sin \theta_{k}}{1+\cos \theta_{k}}\left(\nu_{k}+\nu_{k-1}\right)=2 \tan \frac{\theta_{k}}{2} \cdot \frac{\nu_{k}+\nu_{k-1}}{2} .
$$

by a simple calculation. Therefore, unless the curve is arclength parameterized, there are (at least) two choices of the "vertex normal" from the variational viewpoint : using the length gradient (length descent direction) or using the volume gradient (volume descent direction). This suggests that, in contrast to the smooth case, we have to choose the "prefered" vertex normal according to the energy in question.

First, we will start from the $\gamma \equiv 1$ case.
Example 2.3.10 (Regular polygons). Let us take a discrete curve $\Gamma_{h}^{m, n}=\left\{p_{k}\right\}_{k}$ as in the following way (including non-convex regular $n$-gon with radius $a$ ):

$$
p_{k}=a \exp (2 \pi \sqrt{-1} m k / n), \quad k=0, \ldots, n-1, \quad 1 \leq m \leq n-1,
$$

where we assume that $m$ and $n$ are coprime. In particular, we sometimes call the curve $\Gamma_{h}^{1, n}$ as a convex regular $n$-gon. Note that $\Gamma_{h}^{n-1, n}$ is also convex but it has an opposite unit normal with $\Gamma_{h}^{1, n}$ (usually we assume that $\Gamma_{h}^{1, n}$ has the outward-pointing unit normal). Then the curve $\Gamma_{h}^{m, n}$ is a critical point of the functional Length $+\kappa \mathrm{Vol}$

(a) $\kappa=-1 / \cos (\pi / 5)$
(b) $\kappa=-1 / \cos (2 \pi / 5)$
(c) $\kappa=-1 / \cos (3 \pi / 5)$
(d) $\kappa=-1 / \cos (4 \pi / 5)$

Figure 2.5: Convex and non-convex regular 5 -gons with radii $a=1$.
with $\kappa=-1 /(a \cos (m \pi / n))$. This value is the reciprocal of the radius of the inscribed circle of the polygon (up to the signature). We sometimes say that a convex regular $n$-gon with radius $a$ (and outward-pointing unit normal) has constant curvature $\kappa_{n}=$ $-1 /(a \cos (\pi / n))$. Note that $\cos (\pi m / n)=-\cos (\pi(n-m) / n)$ and $\kappa_{n} \rightarrow-1 / a$ when $n \rightarrow \infty$. We also remark that this value naturally appears in the discrete cylinder (Example 2.4.9).

We will show that these regular polygons are the only critical points for the functional Length $+\kappa$ Vol.

Theorem 2.3.11. Let $\Gamma_{h}=\left\{p_{k}\right\}_{k=0}^{n-1}$ be a closed discrete curve and take $\kappa \in \mathbb{R} \backslash\{0\}$. Then the following two conditions are equivalent:
(1) $\Gamma_{h}$ is a critical point of the functional Length $+\kappa$ Vol.
(2) The edge length $l_{k}$ and the angles $\theta_{k}$ are constant $l_{0}$ and $\theta_{0}$ which satisfies $\kappa l_{0}=$ $2 \tan \left(\theta_{0} / 2\right)$.

Proof. We put

$$
A_{k}:=\left(\nu_{k}-\nu_{k-1}\right)+\frac{\kappa}{2}\left(p_{k+1}-p_{k-1}\right), \quad k=0, \ldots, n-1 .
$$

Then the discrete curve $\Gamma_{h}$ is a critical point of the functional Length $+\kappa \mathrm{Vol}$ if and only if $A_{k}=0$ for all $k$. By a simple calculation we have

$$
\begin{align*}
\left\langle A_{k}, \nu_{k-1}\right\rangle & =\sin \theta_{k}\left(\frac{\kappa l_{k}}{2}-\tan \frac{\theta_{k}}{2}\right)  \tag{3.2}\\
\left\langle A_{k}, \nu_{k}\right\rangle & =\sin \theta_{k}\left(\tan \frac{\theta_{k}}{2}-\frac{\kappa l_{k-1}}{2}\right),  \tag{3.3}\\
\left\langle A_{k+1}, \nu_{k+1}\right\rangle & =\sin \theta_{k+1}\left(\tan \frac{\theta_{k+1}}{2}-\frac{\kappa l_{k}}{2}\right) . \tag{3.4}
\end{align*}
$$

For the necessity, that is, if we assume $A_{k}=0$ for all $k$, then it follows from (3.2) and (3.3) that $\kappa l_{k}=2 \tan \left(\theta_{k} / 2\right)=\kappa l_{k-1}$. And it also follows from (3.3), (3.4) and using $l_{k}=l_{k-1}$ that

$$
\tan \frac{\theta_{k}}{2}=\frac{\kappa l_{k}}{2}=\frac{\kappa l_{k-1}}{2}=\tan \frac{\theta_{k-1}}{2} .
$$

To prove the sufficiency, since $\nu_{k}$ and $p_{k}-p_{k-1}$ forms a basis of $\mathbb{R}^{2}$ and $\left\langle A_{k}, \nu_{k}\right\rangle=0$ by using (3.2) and the assumption, all we have to prove is $\left\langle A_{k}, p_{k}-p_{k-1}\right\rangle=0$ for all $k$. By using the assumption $l_{k}=l_{k-1}=l_{0}$ and $\theta_{k}=\theta_{0}$, we have

$$
\begin{aligned}
\left\langle A_{k}, \frac{p_{k}-p_{k-1}}{l_{k}}\right\rangle & =-\sin \theta_{k}+\frac{\kappa}{2}\left(l_{k} \cos \theta_{k}+l_{k-1}\right) \\
& =\left(1+\cos \theta_{0}\right)\left(\frac{\kappa l_{0}}{2}-\frac{\sin \theta_{0}}{1+\cos \theta_{0}}\right)=\left(1+\cos \theta_{0}\right)\left(\frac{\kappa l_{0}}{2}-\tan \frac{\theta_{0}}{2}\right)=0
\end{aligned}
$$

This shows $A_{k}=0$ and proves the statement.
We found that critical points of the functional Length $+\kappa$ Vol must satisfy $l_{k} \equiv l_{0}$, i.e., they must have "good coordinates (arclength parameter)". If we note that for an arclength parametrized curve we can define the curvature at vertices, the previous result can be restated as follows:

Corollary 2.3.12. Let $\Gamma_{h}$ be an arclength parametrized discrete closed curve, i.e., $l_{k} \equiv l_{0}$, and take $\kappa \in \mathbb{R} \backslash\{0\}$. Then the following two conditions are equivalent:
(1) $\Gamma_{h}$ is a critical point of the functional Length $+\kappa$ Vol.
(2) The discrete curvature $\left(2 / l_{0}\right) \tan \left(\theta_{k} / 2\right)$ is constant $\kappa$.

Corollary 2.3.13. Let $\Gamma_{h}=\left\{p_{k}\right\}$ be a critical point of the functional Length $+\kappa$ Vol. Then $\Gamma_{h}$ must be a regular polygon (including a non-convex polygon).

Proof. Let $r_{k}$ be the radius of the triangle $\left(p_{k-1}, p_{k}, p_{k+1}\right)$. By using the sine law, we have $r_{k}=1 /\left(|\kappa| \cos \left(\theta_{k} / 2\right)\right)$ and this is independent of $k$ by the previous result. Then all vertices of $\Gamma_{h}$ must lie on the same circle with radius $r_{0}=r_{k}$. By a rotation
and translation, we can put $p_{k}=r_{0} \exp \left(i \varphi_{k}\right)$ with $\varphi_{1}=0$. By using the condition $\kappa l_{0}=2 \tan \left(\theta_{0} / 2\right)$, we have

$$
\begin{aligned}
\left|p_{k+1}-p_{k}\right|=l_{k}=l_{0} & \Longleftrightarrow r_{0} \cdot 2 \sin \frac{\varphi_{k+1}-\varphi_{k}}{2}=\frac{2}{|\kappa|} \tan \frac{\left|\theta_{0}\right|}{2} \\
& \Longleftrightarrow \sin \frac{\varphi_{k+1}-\varphi_{k}}{2}=\sin \frac{\left|\theta_{0}\right|}{2} \Longleftrightarrow \varphi_{k+1}-\varphi_{k}=\left|\theta_{0}\right|
\end{aligned}
$$

Therefore we have $\varphi_{k}=(k-1)\left|\theta_{0}\right|$ and this proves the statement.
Remark. We will prove the instability of the non-convex regular polygons (Theorem ).

### 2.3.5 Discrete constant anisotropic curvature curves

In this section we will define the discrete constant anisotropic (mean) curvature (CAMC) curves by using the variational principle inspired by [37]. We will derive the conservation law for the Euler-Lagrange equation and Minkowski-type formula. As an application, we will show another proof of the Chakerian-Lange theorem by using our method (Corollary 2.3.18). Finally we will give some examples of discrete CAMC curves. In connection with Chapter 1, we can give "non-uniqueness" of discrete CAMC curves.

Definition 2.3.14 (Discrete CAMC curve). Let $\gamma: S^{1} \rightarrow \mathbb{R}$ be a $C^{1}$ function. Then a simplicial curve $\Gamma_{h}$ has constant anisotropic (mean) curvature (CAMC) if there exists a constant $\Lambda$ such that

$$
\nabla_{p} \mathscr{F}_{\gamma}+\Lambda \nabla_{p} \mathrm{Vol}=0
$$

holds for every interior vertex $p$. This definition is equivalent to the following equations:

$$
\begin{equation*}
\xi_{\gamma}\left(\nu_{k}\right)-\xi_{\gamma}\left(\nu_{k-1}\right)+\frac{\Lambda}{2}\left(p_{k+1}-p_{k-1}\right)=0 \tag{3.5}
\end{equation*}
$$

for any interior vertex $p_{k}$.
Remark. This definition is inspired by the definition of discrete CMC surfaces in [37]. Note that we defined the constant anisotropic curvature curve without defining the discrete anisotropic curvature. We sometimes call such a curve as CAMC- $\Lambda$ curve.

We first remark the following fact.
Proposition 2.3.15. Let $\gamma: S^{1} \rightarrow \mathbb{R}_{>0}$ be of $C^{2}$ and assume $\gamma^{\prime \prime}(\theta)+\gamma(\theta)>0$ where we define $\gamma(\theta):=\gamma(\cos \theta, \sin \theta)$. Then any bounded curve with zero anisotropic curvature must be a segment.

Proof. When $\gamma \equiv 1$, the statement is trivial. If we assume the function $\gamma$ satisfies $\gamma^{\prime \prime}+\gamma>0$, then the Cahn-Hoffman map $\xi_{\gamma}: S^{1} \rightarrow W_{\gamma}$ becomes diffeomorphism. Therefore, the zero curvature condition $\xi_{\gamma}\left(\nu_{k-1}\right)=\xi_{\gamma}\left(\nu_{k}\right)$ is equivalent to $\nu_{k-1}=\nu_{k}$.

Lemma 2.3.16 (Conservation law). A closed discrete curve $\Gamma_{h}=\left\{p_{k}\right\}_{k}$ is CAMC- $\Lambda$ if and only if there exists a vector $c \in \mathbb{R}^{2}$ such that $\xi_{\gamma}\left(\nu_{k}\right)+(\Lambda / 2)\left(p_{k+1}+p_{k}\right) \equiv c$ for all $k$.

Proof. In the critical point condition (3.5), by inserting $p_{k}-p_{k}$ we have

$$
\xi_{\gamma}\left(\nu_{k}\right)+\frac{\Lambda}{2}\left(p_{k+1}+p_{k}\right)=\xi_{\gamma}\left(\nu_{k-1}\right)+\frac{\Lambda}{2}\left(p_{k}+p_{k-1}\right) \quad \text { for all } k
$$

This implies the vector $\xi_{\gamma}\left(\nu_{k}\right)+(\Lambda / 2)\left(p_{k+1}+p_{k}\right)$ is independent of $k$.
Since the constant vector $c$ is just a translation of the curve, it is enough to consider the case $c=0$, then we have

$$
\begin{equation*}
\frac{p_{k+1}+p_{k}}{2}=-\frac{1}{\Lambda} \xi_{\gamma}\left(\nu_{k}\right) \quad \text { for all } k \tag{3.6}
\end{equation*}
$$

Moreover, since the Cahn-Hoffman map $\xi_{\gamma}: S^{1} \rightarrow \mathbb{R}^{2}$ is a wave front (cf. [24], Proposition 5.1), every midpoint of the edge of the curve is tangent to the rescaled image of the Cahn-Hoffman map $\xi_{\gamma}$. Therefore discrete CAMC- $\Lambda$ curves for a energy density $\gamma$ must be constructed from the lines

$$
\left\{x \in \mathbb{R}^{2} \left\lvert\,\langle x, \nu\rangle=-\frac{1}{\Lambda} \gamma(\nu)\right.\right\}, \quad \nu \in S^{1}
$$

which are the tangent lines of the rescaled Cahn-Hoffman map $(-1 / \Lambda) \xi_{\gamma}: S^{1} \rightarrow \mathbb{R}^{2}$.
Remark. The condition in the above lemma can be considered as a conservation law for the Euler-Lagrange equation (3.5). Note that the Euler-Lagrange equation (3.5) is the vertex-based condition but the conservation law is the edge-based condition. Moreover, some calculation shows the following relation:

$$
\delta\left(\mathscr{F}_{\gamma}+\Lambda \mathrm{Vol}\right)=\sum_{k}\left\langle\xi_{\gamma}\left(\nu_{k}\right)+\frac{\Lambda}{2}\left(p_{k}+p_{k+1}\right), R\left(v_{k}-v_{k+1}\right)\right\rangle .
$$

From another point of view, if the curve $\left\{p_{k}\right\}_{k}$ approximates a smooth curve $p(s)$, then $p_{k} \approx p_{k+1} \approx p(s)$ and therefore $p(s)=-(1 / \Lambda) \xi_{\gamma}(\nu)$. In particular $p(s)=\xi_{\gamma}(\nu)$ when $\Lambda=-1$. This fact is closely related to the result of F. Morgan [30], but since we did not assume the convexity for $\gamma$, the Wulff shape is replaced by the image of the Cahn-Hoffman map.

Taking the inner product (3.6) with $\nu_{k}$ gives the following relation:

$$
\gamma\left(\nu_{k}\right)=\left\langle\xi_{\gamma}\left(\nu_{k}\right), \nu_{k}\right\rangle=-\frac{\Lambda}{2}\left\langle p_{k+1}+p_{k}, \nu_{k}\right\rangle=-\Lambda\left\langle p_{k}, \nu_{k}\right\rangle .
$$

In particular, we have the following Minkowski-type formula:

Proposition 2.3.17 (Discrete Minkowski-type formula). Let $\gamma: S^{1} \rightarrow \mathbb{R}$ be a $C^{1}$ function and $\Gamma_{h}=\left\{p_{k}\right\}_{k}$ be a closed discrete CAMC- $\Lambda$ curve. Then we have the following Minkowski-type formula:

$$
\sum_{k}\left(\gamma\left(\nu_{k}\right)+\Lambda\left\langle p_{k}, \nu_{k}\right\rangle\right)\left|p_{k+1}-p_{k}\right|=0
$$

In connection with the convex geometry, we mention about the discrete isoperimetric problem, i.e., we consider the following problem:

For a given convex region $W$ and number $n$, find the convex n-gon which have the least area containing $W$.

This kind of problem is interested not only from the pure mathematics, but also from the robotics and computer aided design field, see e.g. [1]. Here we will show that our discretization method gives another proof of the following theorem:

Corollary 2.3.18 (Chakerian-Lange [8]). Let $W$ be a strictly convex domain, $n \geq 3$ and $\Gamma_{h}=\left\{p_{k}\right\}_{k}$ be a convex $n$-gon of minimum area containing $W$. Then, the midpoints of the edges of $\Gamma_{h}$ must lie on $\partial W$.

Remark. The original theorem assumed that $W$ is just a convex region. As we will see, we have to remove the regularity assumption for the energy density function in order to overcome this restriction.

For any strictly convex domain $W$, it is known that there exists a convex and $C^{1}$ energy density $\gamma: S^{1} \rightarrow \mathbb{R}_{>0}$ with $\partial W=\xi_{\gamma}\left(S^{1}\right)([14])$. We say a closed discrete curve $\Gamma_{h}=\left\{p_{k}\right\}_{k}$ is circumscribed about $W_{\gamma}$ if for every $k$ there exist $\lambda_{k} \in \mathbb{R}$ such that $\left(1-\lambda_{k}\right) p_{k}+\lambda_{k} p_{k+1}=\xi_{\gamma}\left(\nu_{k}\right)$. Note that we have

$$
\mathscr{F}_{\gamma}\left(\Gamma_{h}\right)=2 \operatorname{Vol}\left(\Gamma_{h}\right) \Longleftrightarrow \mathscr{F}_{\gamma}\left(\Gamma_{h}\right)-\operatorname{Vol}\left(\Gamma_{h}\right)=\operatorname{Vol}\left(\Gamma_{h}\right)
$$

for any discrete curve $\Gamma_{h}$ circumscribed about $W_{\gamma}$. Therefore we have the following lemma:

Lemma 2.3.19. For any discrete curve circumscribed about $W_{\gamma}$, the following problems are equivalent:

- Find the minimizer of the energy $\mathscr{F}_{\gamma}-\mathrm{Vol}$ (CAMC-(-1) curve).
- Find the minimizer of the 2 -dimensional volume Vol.

In particular, the existence of the minimizer of the second problem guarantees the existence of the minimizer of the first problem.

Proof of Corollary 2.3.18. By the assumption, there exists a convex and $C^{1}$ energy density $\gamma: S^{1} \rightarrow \mathbb{R}_{>0}$ with $\partial W=\xi_{\gamma}\left(S^{1}\right)$. Consider the variational problem of discrete curves for the functional $\mathscr{F}_{\gamma}-$ Vol. Then by the previous lemma, among discrete curves circumscribed about $\partial W$, the minimizer of the bounded area coincides with the minimizer of the energy $\mathscr{F}_{\gamma}-$ Vol. From the conservation law (Lemma 2.3.16), we have

$$
\frac{p_{k}+p_{k+1}}{2}=\xi_{\gamma}\left(\nu_{k}\right) \quad \text { for all } k
$$

for the minimizer $\Gamma_{h}$. This proves the statement.
We will give some examples of discrete CAMC curves. As we will see, non-convexity of the energy density also affects the "non-uniqueness" of the discrete CAMC curves.

Example 2.3.20. We take the energy density $\gamma: S^{1} \rightarrow \mathbb{R}_{>0}$ as $\gamma\left(\nu_{1}, \nu_{2}\right)=\sqrt{a^{2} \nu_{1}^{2}+b^{2} \nu_{2}^{2}}$, where $a, b \in \mathbb{R}_{>0}$ are constant. Then the corresponding Wulff shape becomes the ellipse $(x / a)^{2}+(y / b)^{2}=1$. For positive numbers $A, B>0$, we define a map $T_{A, B}: \mathbb{R}^{2} \rightarrow \mathbb{R}^{2}$ as $T_{A, B}\left(x_{1}, x_{2}\right):=\left(A x_{1}, B x_{2}\right)$. By a typical calculation, we can check the following fact:

Proposition 2.3.21 (for smooth case, see [25], Example 4.4). Let $\Gamma_{h}=\left\{p_{k}\right\}_{k}$ be a closed discrete curve and $\Lambda \neq 0$. Then $\Gamma_{h}$ is a critical point of the functional $\mathscr{F}_{\gamma}+\Lambda \mathrm{Vol}$ if and only if $\widetilde{\Gamma}_{h}:=\left\{T_{1 / a, 1 / b}\left(p_{k}\right)\right\}$ is a critical point of Length $+\Lambda$ Vol.

Therefore, by using the previous result, every discrete CAMC curve for the density $\gamma$ is "rescaled" regular polygons, i.e., if $\Gamma_{h}$ is a regular polygon, then $T_{a, b}\left(\Gamma_{h}\right)$ is CAMC for $\gamma$.

The following Figure 2.6 give examples of the discrete CAMC curve $(\Lambda=-1)$ for the energy density $\gamma\left(\nu_{1}, \nu_{2}\right)=\sqrt{4 \nu_{1}^{2}+\nu_{2}^{2}}$. The dotted curve represents the image of the Cahn-Hoffman map, i.e., the ellipse $x^{2} / 4+y^{2}=1$. In the first figure, the vertices are

$$
(2,1),(-2,1),(-2,-1),(2,-1)
$$

In the second figure, the vertices are

$$
\begin{aligned}
& (2, \sqrt{2}-1),(2(\sqrt{2}-1), 1),(2(1-\sqrt{2}), 1),(-2, \sqrt{2}-1) \\
& (-2,1-\sqrt{2}),(2(1-\sqrt{2}),-1),(2(\sqrt{2}-1),-1),(2,1-\sqrt{2})
\end{aligned}
$$

Note that the first and third figure have the same volume, therefore this means the minimizer for the functional $\mathscr{F}_{\gamma}-\mathrm{Vol}$ is not unique.

Example 2.3.22. Let us take a non-convex energy density which appeared in the Chapter 1, i.e., we put $\gamma\left(\nu_{1}, \nu_{2}\right)=\nu_{1}^{6}+\nu_{2}^{6}$. The dotted curve represents the image of the Cahn-Hoffman map for $\gamma$. As in the smooth case, we have discrete CAMC- $(-1)$ curves which are quite different from the Wulff shape like in the second and third figures. As we will see later, these curves are unstable (Proposition 2.3.27).


Figure 2.6: Examples of discrete CAMC- $(-1)$ curves for $\gamma\left(\nu_{1}, \nu_{2}\right)=\sqrt{4 \nu_{1}^{2}+\nu_{2}^{2}}$.


Figure 2.7: Examples of discrete CAMC- $(-1)$ curves for $\gamma\left(\nu_{1}, \nu_{2}\right)=\nu_{1}^{6}+\nu_{2}^{6}$.

### 2.3.6 Second variation of the anisotropic energy

In this section we consider the second variation of the anisotropic energy and stability of discrete CAMC curves. We will follow the argument developed in [37]. We assume that the energy density $\gamma$ is of class $C^{2}$ in the following.

Let $G=(V, E)$ be a standard $(n-1)$-circle and $p: V \rightarrow \mathbb{R}^{2}$ be its realization and denote its image by $\Gamma_{h}=\left\{p_{k}\right\}_{k}$. We consider a smooth variation $p_{k}(t)=p_{k}+t v_{k}+$ $O\left(t^{2}\right)$. Recall the first variation formula of the anisotropic energy and the 2-dimensional volume:

$$
\begin{aligned}
\frac{d}{d t} \mathscr{F}_{\gamma}=\sum_{k}\left\langle\nabla_{p_{k}} \mathscr{F}_{\gamma}, p_{k}^{\prime}\right\rangle, & \nabla_{p_{k}} \mathscr{F}_{\gamma}=R\left(\xi_{\gamma}\left(\nu_{k}\right)-\xi_{\gamma}\left(\nu_{k}\right)\right), \\
\frac{d}{d t} \mathrm{Vol}=\sum_{k}\left\langle\nabla_{p_{k}} \mathrm{Vol}, p_{k}^{\prime}\right\rangle, & \nabla_{p_{k}} \mathrm{Vol}=\frac{1}{2} R\left(p_{k+1}-p_{k-1}\right) .
\end{aligned}
$$

Note that if the variation is volume-preserving, then

$$
0=\delta^{2} \mathrm{Vol}=\sum_{k}\left\langle\delta\left(\nabla_{p_{k}} \mathrm{Vol}\right), \delta p_{k}\right\rangle+\left\langle\nabla_{p_{k}} \mathrm{Vol}, \delta^{2} p_{k}\right\rangle .
$$

We say a variation is admissible (or permissible) if the variation is volume-preserving and fixes the boundary.
Lemma 2.3.23. Let $\Gamma_{h}=\left\{p_{k}\right\}_{k}$ be a discrete CAMC- $\Lambda$ curve and $p_{k}(t)=p_{k}+t v_{k}+$ $O\left(t^{2}\right)$ be an admissible variation. Then we have

$$
\delta^{2} \mathscr{F}_{\gamma}:=\frac{d^{2}}{d t^{2}}{ }_{\mid t=0} \mathscr{F}_{\gamma}=\sum_{k}\left\langle\delta\left(\nabla_{p_{k}} \mathscr{F}_{\gamma}+\Lambda \nabla_{p_{k}} \mathrm{Vol}\right), v_{k}\right\rangle .
$$

Proof. From the CAMC condition (3.5) we have

$$
\begin{aligned}
\delta^{2} \mathscr{F}_{\gamma} & =\sum_{k} \delta\left\langle\nabla_{p_{k}} \mathscr{F}_{\gamma}, p_{k}^{\prime}\right\rangle=\sum_{k}\left\langle\delta \nabla_{p_{k}} \mathscr{F}_{\gamma}, \delta p_{k}\right\rangle+\left\langle\nabla_{p_{k}} \mathscr{F}_{\gamma}, \delta^{2} p_{k}\right\rangle \\
& =\sum_{k}\left\langle\delta \nabla_{p_{k}} \mathscr{F}_{\gamma}, v_{k}\right\rangle+\left\langle-\Lambda \nabla_{p_{k}} \operatorname{Vol}, \delta^{2} p_{k}\right\rangle \\
& =\sum_{k}\left\langle\delta \nabla_{p_{k}} \mathscr{F}_{\gamma}, v_{k}\right\rangle+\left\langle\Lambda \delta \nabla_{p_{k}} \operatorname{Vol}, \delta p_{k}\right\rangle=\sum_{k}\left\langle\delta\left(\nabla_{p_{k}} \mathscr{F}_{\gamma}+\Lambda \nabla_{p_{k}} \operatorname{Vol}\right), v_{k}\right\rangle
\end{aligned}
$$

Definition 2.3.24 (Stability of discrete CAMC curve). Let $\gamma: S^{1} \rightarrow \mathbb{R}$ be a $C^{2}$ function and $\Gamma_{h}=\left\{p_{k}\right\}_{k}$ be a closed discrete CAMC curve. Then $\Gamma_{h}$ is said to be stable if $\delta^{2} \mathscr{F}_{\gamma} \geq 0$ for any admissible variation.

We introduce the matrix $Q^{\gamma}$ and $Q^{V}$ as follows:

$$
{ }^{t} \vec{v} Q^{\gamma} \vec{v}=\sum_{k}\left\langle\delta \nabla_{p_{k}} \mathscr{F}_{\gamma}, v_{k}\right\rangle, \quad{ }^{t} \vec{v} Q^{V} \vec{v}=\sum_{k}\left\langle\delta \nabla_{p_{k}} \operatorname{Vol}, v_{k}\right\rangle .
$$

Then we can write $\delta^{2} \mathscr{F}_{\gamma}={ }^{t} \vec{v}\left(Q^{\gamma}+\Lambda Q^{V}\right) \vec{v}$.
Lemma 2.3.25. ${ }^{t} \vec{v} Q^{V} \vec{v}=\sum_{k}\left\langle v_{k}, R v_{k+1}\right\rangle$.
Proof. The proof follows from the direct computation:

$$
\begin{aligned}
{ }^{t} \vec{v} Q^{V} \vec{v} & =\sum_{k}\left\langle\delta \nabla_{p_{k}} \operatorname{Vol}, v_{k}\right\rangle=\frac{1}{2} \sum_{k}\left\langle\delta R\left(p_{k+1}-p_{k-1}\right), v_{k}\right\rangle \\
& =\frac{1}{2} \sum_{k}\left\langle R\left(v_{k+1}-v_{k-1}\right), v_{k}\right\rangle=\sum_{k}\left\langle v_{k}, R v_{k+1}\right\rangle .
\end{aligned}
$$

Lemma 2.3.26. ${ }^{t} \vec{v} Q^{\gamma} \vec{v}=\sum_{k}\left\langle d \xi_{\gamma}\left(\delta \nu_{k}\right), \delta \nu_{k}\right\rangle l_{k}$.
Proof. If we note that $d \xi_{\gamma}$ is an endomorphism of the tangent space and $\delta \nu_{k}$ is tangent to the edge $e_{k}=\left[p_{k}, p_{k+1}\right]$, then we have

$$
\begin{aligned}
{ }^{t} \vec{v} Q^{\gamma} \vec{v} & =\sum_{k}\left\langle\delta \nabla_{p_{k}} \mathscr{F}_{\gamma}, v_{k}\right\rangle=\sum_{k}\left\langle\delta R\left(\xi_{\gamma}\left(\nu_{k}\right)-\xi_{\gamma}\left(\nu_{k-1}\right)\right), v_{k}\right\rangle \\
& =\sum_{k}\left\langle R\left(d \xi_{\gamma}\left(\delta \nu_{k}\right)-d \xi_{\gamma}\left(\delta \nu_{k-1}\right)\right), v_{k}\right\rangle=\sum_{k}\left\langle R\left(d \xi_{\gamma}\left(\delta \nu_{k}\right)\right), v_{k}-v_{k+1}\right\rangle \\
& =\sum_{k}\left\langle d \xi_{\gamma}\left(\delta \nu_{k}\right), R\left(v_{k+1}-v_{k}\right)\right\rangle=\sum_{k}\left\langle d \xi_{\gamma}\left(\delta \nu_{k}\right), \delta \nu_{k}\right\rangle l_{k}
\end{aligned}
$$

Therefore we have an expression of second variation for the anisotropic energy:

$$
\delta^{2} \mathscr{F}_{\gamma}=\sum_{k}\left(\left\langle d \xi_{\gamma}\left(\delta \nu_{k}\right), \delta \nu_{k}\right\rangle l_{k}+\left\langle v_{k}, R v_{k+1}\right\rangle\right) .
$$

As a special case, we consider an one point variation, i.e., for a fixed vertex $p_{k}$, consider a variation vector field $v_{k} \in \mathbb{R}^{2}$ with $v_{j}=0$ for $j \neq k$. As we saw in Chaper 1 , the matrix $\left.d \xi_{\gamma}\right|_{\nu}$ can be written as $\gamma^{\prime \prime}(\theta)+\gamma(\theta)$ when we write $\nu=(\cos \theta, \sin \theta)$ and $\gamma(\theta)=\gamma(\cos \theta, \sin \theta)$. Then the second variation formula reduces to the following form:

$$
\delta^{2} \mathscr{F}_{\gamma}=\left\langle d \xi_{\gamma}\left(\delta \nu_{k}\right), \delta \nu_{k}\right\rangle l_{k}=\left.\left(\gamma^{\prime \prime}+\gamma\right)\right|_{\nu_{k}} \cdot\left|R v_{k}-\left\langle R v_{k}, \nu_{k}\right\rangle \nu_{k}\right|^{2} / l_{k} .
$$

Therefore, we have the following necessary condition for the stability:
Proposition 2.3.27. Let $\gamma: S^{1} \rightarrow \mathbb{R}_{>0}$ be of $C^{2}$ and $\Gamma_{h}$ be a discrete CAMC curve. Then $\Gamma_{h}$ is stable only if every edge of the curve is tangent to a portion of the CahnHoffman map for $\gamma$ which satisfies $\gamma^{\prime \prime}+\gamma \geq 0$, i.e., the convex part of the Cahn-Hoffman map.

Example 2.3.28. The second and third curves in the Figure 2.7 are unstable since the vertical and horizontal edges are tangent to the non-convex part of the CahnHoffman map (see Figure 2.8, 2.9 for the energy descent deformations under the volumeconstraint condition).

(a) energy $=8$.
(b) energy $\approx 5.63944664472859$.

Figure 2.8: Energy descent deformation (rotation) for unstable discrete CAMC curve.


Figure 2.9: Energy descent deformation for unstable discrete CAMC curve.

In the following we will focus on the $\gamma \equiv 1$ case and denote $Q^{\gamma}=Q$ in this case.

Proposition 2.3.29 (Second variation formula for the length functional).

$$
Q=\sum_{k} \frac{1}{l_{k}}\left(\left|v_{k+1}-v_{k}\right|^{2}-\left\langle v_{k+1}-v_{k}, R \nu_{k}\right\rangle^{2}\right)=\sum_{k}\left(\left|\nabla v_{k}\right|^{2}-\left\langle\nabla v_{k}, R \nu_{k}\right\rangle^{2}\right) l_{k},
$$

therefore we have the following second variation formula for the length functional:

$$
\begin{equation*}
\delta^{2} \text { Length }=\sum_{k}\left(\left|\nabla v_{k}\right|^{2}-\left\langle\nabla v_{k}, R \nu_{k}\right\rangle^{2}\right) l_{k}+\Lambda\left\langle v_{k}, R v_{k+1}\right\rangle . \tag{3.7}
\end{equation*}
$$

Proof. By using Lemma 2.3.2, we have

$$
\begin{aligned}
Q & =\sum_{k}\left\langle\delta \nu_{k}, R\left(v_{k+1}-v_{k}\right)\right\rangle=\sum_{k}\left\langle R\left(\nabla v_{k}\right)-\left\langle R\left(\nabla v_{k}\right), \nu_{k}\right\rangle \nu_{k}, R\left(\nabla v_{k}\right)\right\rangle l_{k} \\
& =\sum_{k}\left(\left|\nabla v_{k}\right|^{2}-\left\langle R\left(\nabla v_{k}\right), \nu_{k}\right\rangle^{2}\right) l_{k}=\sum_{k}\left(\left|\nabla v_{k}\right|^{2}-\left\langle\nabla v_{k}, R \nu_{k}\right\rangle^{2}\right) l_{k} .
\end{aligned}
$$

In the section of the Steiner-type formula, we used the vector

$$
N_{k}:=\frac{R\left(\nu_{k-1}-\nu_{k}\right)}{\sin \theta_{k}}=\frac{1}{\sin \theta_{k}}\left(\frac{p_{k+1}-p_{k}}{l_{k}}-\frac{p_{k}-p_{k-1}}{l_{k-1}}\right)
$$

as the "normal vector" at the vertex $p_{k}$. If we define the "tangent vector" $T_{k}$ as $T_{k}=-R N_{k}$, then we can decompose the variation vector $v_{k}$ as

$$
v_{k}=\psi_{k} N_{k}+\eta_{k} T_{k}=\psi_{k} N_{k}-\eta_{k} R N_{k},
$$

where $\psi, \eta: V \rightarrow \mathbb{R}$ is some functions on the vertices. If $\eta_{k}=0$ for all $k$, we sometimes call it the normal variation. In the following we will use this notation.

Lemma 2.3.30. The first variation formula of the volume can be written as

$$
\delta \mathrm{Vol}=\frac{1}{2} \sum_{k}\left(\psi_{k}\left(l_{k}+l_{k-1}\right)+\eta_{k}\left(l_{k}-l_{k-1}\right) \tan \frac{\theta_{k}}{2}\right) .
$$

In particular, if a curve $\left\{p_{k}\right\}$ satisfies $l_{k} \equiv l_{0}$, then $\sum_{k} \psi_{k}=0$ for any volume-preserving variation.

Remark. Conversely, for such a curve, if a function $\psi: V \rightarrow \mathbb{R}$ satiefies $\sum_{k} \psi_{k}=0$, then we can construct a volume-preserving normal variation whose variation vector is $\psi_{k} N_{k}$ (cf. [4]).

Proof. Using the first variation formula, we have

$$
\begin{aligned}
\delta \mathrm{Vol} & =\frac{1}{2} \sum_{k}\left\langle R\left(p_{k+1}-p_{k-1}\right), v_{k}\right\rangle=\frac{1}{2} \sum_{k}\left\langle R\left(p_{k+1}-p_{k-1}\right), \psi_{k} N_{k}+\eta_{k} T_{k}\right\rangle \\
& =\frac{1}{2} \sum_{k}\left[\psi_{k} l_{k} \frac{\left\langle\nu_{k}, R \nu_{k-1}\right\rangle}{\sin \theta_{k}}+\psi_{k} l_{k-1} \frac{\left\langle\nu_{k-1},-R \nu_{k}\right\rangle}{\sin \theta_{k}}+\frac{\eta_{k}}{\sin \theta_{k}}\left(l_{k}-l_{k-1}\right)\left(1-\cos \theta_{k}\right)\right] \\
& =\frac{1}{2} \sum_{k}\left(\psi_{k}\left(l_{k}+l_{k-1}\right)+\eta_{k}\left(l_{k}-l_{k-1}\right) \tan \frac{\theta_{k}}{2}\right) .
\end{aligned}
$$

Recall that if $\left\{p_{k}\right\}_{k}$ is a critical point of the functional Length $+\Lambda$ Vol, then we have $l_{k} \equiv l_{0}, \theta_{k} \equiv \theta_{0}$ and $\Lambda l_{0}=2 \tan \left(\theta_{0} / 2\right)$.

## Lemma 2.3.31.

$$
\left|v_{k+1}-v_{k}\right|^{2}-\left\langle v_{k+1}-v_{k}, R \nu_{k}\right\rangle^{2}=\left[\left(\psi_{k+1}-\psi_{k}\right)+\tan \left(\theta_{0} / 2\right)\left(\eta_{k+1}+\eta_{k}\right)\right]^{2}
$$

Therefore we have

$$
\left|\nabla v_{k}\right|^{2}-\left\langle\nabla v_{k}, R \nu_{k}\right\rangle^{2}=\left(\nabla \psi_{k}+\frac{\Lambda}{2}\left(\eta_{k}+\eta_{k+1}\right)\right)^{2}
$$

Proof. Recall that

$$
\begin{aligned}
\left\langle N_{k},-R N_{k+1}\right\rangle & =\tan \frac{\theta_{k}}{2}+\tan \frac{\theta_{k+1}}{2}=2 \tan \frac{\theta_{0}}{2}, \\
\left\langle N_{k}, N_{k+1}\right\rangle & =1-\tan \frac{\theta_{k}}{2} \tan \frac{\theta_{k+1}}{2}=1-\tan ^{2} \frac{\theta_{0}}{2} .
\end{aligned}
$$

If we note $\left|N_{k}\right|=1 / \cos \left(\theta_{0} / 2\right)$, then

$$
\begin{aligned}
& \left|v_{k+1}-v_{k}\right|^{2} \\
= & \left|\psi_{k+1} N_{k+1}-\psi_{k} N_{k}\right|^{2}+2\left\langle\psi_{k+1} N_{k+1}-\psi_{k} N_{k}, \eta_{k+1} T_{k+1}-\eta_{k} T_{k}\right\rangle+\left|\eta_{k+1} T_{k+1}-\eta_{k} T_{k}\right|^{2} \\
= & \frac{\psi_{k+1}^{2}}{\cos ^{2}\left(\theta_{0} / 2\right)}-2 \psi_{k} \psi_{k+1}\left(1-\tan ^{2}\left(\theta_{0} / 2\right)\right)+\frac{\psi_{k}^{2}}{\cos ^{2}\left(\theta_{0} / 2\right)}-2\left\langle N_{k}, R N_{k+1}\right\rangle\left(\psi_{k+1} \eta_{k}-\psi_{k} \eta_{k+1}\right) \\
& \quad+\frac{\eta_{k+1}^{2}}{\cos ^{2}\left(\theta_{0} / 2\right)}-2 \eta_{k} \eta_{k+1}\left(1-\tan ^{2}\left(\theta_{0} / 2\right)\right)+\frac{\eta_{k}^{2}}{\cos ^{2}\left(\theta_{0} / 2\right)} \\
= & \left(1+\tan ^{2}\left(\theta_{0} / 2\right)\right)\left(\psi_{k+1}^{2}+\psi_{k}^{2}+\eta_{k+1}^{2}+\eta_{k}^{2}\right) \\
& \quad-2\left(\psi_{k} \psi_{k+1}+\eta_{k} \eta_{k+1}\right)\left(1-\tan ^{2}\left(\theta_{0} / 2\right)\right)+4 \tan \left(\theta_{0} / 2\right)\left(\eta_{k} \psi_{k+1}-\eta_{k+1} \psi_{k}\right) .
\end{aligned}
$$

Similarly since

$$
\begin{aligned}
& \left\langle v_{k+1}-v_{k}, R \nu_{k}\right\rangle \\
= & \left\langle\psi_{k+1} N_{k+1}-\psi_{k} N_{k}+\eta_{k+1} T_{k+1}-\eta_{k} T_{k}, R \nu_{k}\right\rangle \\
= & \psi_{k+1} \frac{\left\langle\nu_{k}-\nu_{k+1}, \nu_{k}\right\rangle}{\sin \theta_{k+1}}-\psi_{k} \frac{\left\langle\nu_{k-1}-\nu_{k}, \nu_{k}\right\rangle}{\sin \theta_{k}}-\eta_{k+1} \frac{\left\langle\nu_{k+1}-\nu_{k}, R \nu_{k}\right\rangle}{\sin \theta_{k+1}}+\eta_{k} \frac{\left\langle\nu_{k}-\nu_{k-1}, R \nu_{k}\right\rangle}{\sin \theta_{k}} \\
= & \left(\psi_{k+1}+\psi_{k}\right) \tan \frac{\theta_{0}}{2}-\left(\eta_{k+1}-\eta_{k}\right),
\end{aligned}
$$

we have

$$
\left\langle v_{k+1}-v_{k}, R \nu_{k}\right\rangle^{2}=\left(\psi_{k+1}+\psi_{k}\right)^{2} \tan ^{2} \frac{\theta_{0}}{2}-2 \tan \frac{\theta_{0}}{2}\left(\psi_{k+1}+\psi_{k}\right)\left(\eta_{k+1}-\eta_{k}\right)+\left(\eta_{k+1}-\eta_{k}\right)^{2} .
$$

Substracting these factors we have

$$
\begin{aligned}
& \left|v_{k+1}-v_{k}\right|^{2}-\left\langle v_{k+1}-v_{k}, R \nu_{k}\right\rangle^{2} \\
= & \psi_{k+1}^{2}+\psi_{k}^{2}+\tan ^{2} \frac{\theta_{0}}{2}\left(\eta_{k+1}^{2}+\eta_{k}^{2}\right)-2 \psi_{k} \psi_{k+1} \tan ^{2} \frac{\theta_{0}}{2}+2 \eta_{k} \eta_{k+1} \\
& \quad+2 \tan \frac{\theta_{0}}{2}\left(2 \eta_{k} \psi_{k+1}-2 \eta_{k+1} \psi_{k}+\psi_{k+1} \eta_{k+1}-\psi_{k+1} \eta_{k}+\psi_{k} \eta_{k+1}-\psi_{k} \eta_{k}\right) \\
\quad & \quad-2\left(\psi_{k} \psi_{k+1}+\eta_{k} \eta_{k+1}\right)+2 \tan ^{2} \frac{\theta_{0}}{2}\left(\psi_{k} \psi_{k+1}+\eta_{k} \eta_{k+1}\right) \\
= & \left(\psi_{k+1}-\psi_{k}\right)^{2}+\tan ^{2} \frac{\theta_{0}}{2}\left(\eta_{k+1}+\eta_{k}\right)^{2}+2 \tan \frac{\theta_{0}}{2}\left(\psi_{k+1}-\psi_{k}\right)\left(\eta_{k+1}+\eta_{k}\right) \\
= & {\left[\left(\psi_{k+1}-\psi_{k}\right)+\tan \frac{\theta_{0}}{2}\left(\eta_{k+1}+\eta_{k}\right)\right]^{2} }
\end{aligned}
$$

## Lemma 2.3.32.

$$
\left\langle v_{k}, R v_{k+1}\right\rangle=-2 \tan \frac{\theta_{0}}{2}\left(\psi_{k} \psi_{k+1}+\eta_{k} \eta_{k+1}\right)-\left(1-\tan ^{2} \frac{\theta_{0}}{2}\right)\left(\eta_{k} \psi_{k+1}-\eta_{k+1} \psi_{k}\right)
$$

Proof. This is also a simple calculation:

$$
\begin{aligned}
\left\langle v_{k}, R v_{k+1}\right\rangle & =\left\langle\psi_{k} N_{k}-\eta_{k} R N_{k}, \psi_{k+1} R N_{k+1}+\eta_{k+1} N_{k+1}\right\rangle \\
& =\left(\psi_{k} \psi_{k+1}+\eta_{k} \eta_{k+1}\right)\left\langle N_{k}, R N_{k+1}\right\rangle-\left(\eta_{k} \psi_{k+1}-\eta_{k+1} \psi_{k}\right)\left\langle N_{k}, N_{k+1}\right\rangle \\
& =-2 \tan \frac{\theta_{0}}{2}\left(\psi_{k} \psi_{k+1}+\eta_{k} \eta_{k+1}\right)-\left(1-\tan ^{2} \frac{\theta_{0}}{2}\right)\left(\eta_{k} \psi_{k+1}-\eta_{k+1} \psi_{k}\right) .
\end{aligned}
$$

Theorem 2.3.33 (Second variation formula for the length functional).

$$
\begin{aligned}
\delta^{2} \text { Length } & =\sum_{k}\left[\left|\nabla \psi_{k}\right|^{2}-\Lambda^{2} \psi_{k} \psi_{k+1}+\Lambda \tan ^{2} \frac{\theta_{0}}{2} \nabla \psi_{k}\left(\eta_{k+1}+\eta_{k}\right)+\frac{\Lambda^{2}}{4}\left(\eta_{k+1}-\eta_{k}\right)^{2}\right] l_{0} \\
& =\sum_{k}\left[\left|\nabla \psi_{k}\right|^{2}-\Lambda^{2} \psi_{k} \psi_{k+1}+\tan ^{2} \frac{\theta_{0}}{2}\left(\Lambda \nabla \psi_{k}\left(\eta_{k+1}+\eta_{k}\right)+\left|\nabla \eta_{k}\right|^{2}\right)\right] l_{0}
\end{aligned}
$$

In particular, for a normal variation we have

$$
\delta^{2} \text { Length }=\sum_{k}\left(\left|\nabla \psi_{k}\right|^{2}-\Lambda^{2} \psi_{k} \psi_{k+1}\right) l_{0}=-\sum_{k} \psi_{k}\left(\Delta \psi_{k}+\Lambda^{2} \psi_{k+1}\right) l_{0},
$$

where we use the integration by parts and take the line element at the vertex as $L_{k}=$ $\left(l_{k}+l_{k-1}\right) / 2=l_{0}$.

Proof. By using the previous lemmas, we have

$$
\begin{aligned}
& \left(\left|\nabla v_{k}\right|^{2}-\left\langle\nabla v_{k}, R \nu_{k}\right\rangle^{2}\right) l_{k}+\Lambda\left\langle v_{k}, R v_{k+1}\right\rangle \\
= & \left(\left|\nabla \psi_{k}\right|^{2}+\Lambda \nabla \psi_{k}\left(\eta_{k+1}+\eta_{k}\right)+\frac{\Lambda^{2}}{4}\left(\eta_{k+1}+\eta_{k}\right)^{2}\right) l_{0} \\
& -\Lambda^{2} l_{0}\left(\psi_{k} \psi_{k+1}+\eta_{k} \eta_{k+1}\right)-\Lambda\left(1-\tan ^{2} \frac{\theta_{0}}{2}\right)\left(\eta_{k} \psi_{k+1}-\eta_{k+1} \psi_{k}\right) \\
= & \left|\nabla \psi_{k}\right|^{2} l_{0}+\Lambda \nabla \psi_{k}\left(\eta_{k+1}+\eta_{k}\right) l_{0}+\frac{\Lambda^{2}}{4}\left(\eta_{k+1}-\eta_{k}\right)^{2} l_{0}-\Lambda^{2} \psi_{k} \psi_{k+1} l_{0} \\
& \quad-\Lambda\left(1-\tan ^{2} \frac{\theta_{0}}{2}\right)\left(\left(\psi_{k+1}-\psi_{k}\right)\left(\eta_{k+1}+\eta_{k}\right)-\left(\psi_{k+1} \eta_{k+1}-\psi_{k} \eta_{k}\right)\right) \\
= & \left|\nabla \psi_{k}\right|^{2} l_{0}+\Lambda \tan ^{2} \frac{\theta_{0}}{2} \nabla \psi_{k}\left(\eta_{k+1}+\eta_{k}\right) l_{0}+\frac{\Lambda^{2}}{4}\left(\eta_{k+1}-\eta_{k}\right)^{2} l_{0}-\Lambda^{2} \psi_{k} \psi_{k+1} l_{0} \\
& \quad+\Lambda\left(1-\tan ^{2} \frac{\theta_{0}}{2}\right)\left(\psi_{k+1} \eta_{k+1}-\psi_{k} \eta_{k}\right) .
\end{aligned}
$$

Taking the summation, we have the desired result.

### 2.3.7 Instability of non-convex regular polygons

It is known that embedded regular polygons are global minimizer for the discrete isoperimetric problem (see e.g. [13]). In this section we will prove the instability of the nonconvex regular polygons. To prove this, we find a special variation with a help of the following discrete version of Wirtinger's inequality:

Theorem 2.3.34 (Discrete Wirtinger's inequality, [12]). Let $\psi_{0}, \ldots, \psi_{n}$ be $(n+1)$ real numbers such that

$$
\psi_{0}=\psi_{n}, \quad \sum_{k=0}^{n-1} \psi_{k}=0
$$

Then we have

$$
\sum_{k=0}^{n-1}\left(\psi_{k+1}-\psi_{k}\right)^{2} \geq 4 \sin ^{2} \frac{\pi}{n} \sum_{k=0}^{n-1} \psi_{k}^{2}
$$

and the equality holds if and only if there exist $A, B \in \mathbb{R}$ such that

$$
\psi_{k}=A \cos \frac{2 \pi k}{n}+B \sin \frac{2 \pi k}{n} .
$$

We inculde the proof of this theorem for the completeness. The essential part of the proof is the analysis of the discrete Laplacian by using the circulant matrix method.

Proof of Theorem 2.3.34. By some calculation we have

$$
\sum_{k}\left(\psi_{k+1}-\psi_{k}\right)^{2}=\sum_{k}\left(-\psi_{k-1}+2 \psi_{k}-\psi_{k+1}\right) \psi_{k}=\langle L \Psi, \Psi\rangle,
$$

where we put

$$
L=\left(\begin{array}{cccccc}
2 & -1 & 0 & \cdots & 0 & -1 \\
-1 & 2 & -1 & \cdots & 0 & 0 \\
0 & -1 & 2 & \cdots & 0 & 0 \\
\vdots & \vdots & \vdots & \ddots & \vdots & \vdots \\
-1 & 0 & 0 & \cdots & -1 & 2
\end{array}\right), \quad \Psi=\left(\begin{array}{c}
\psi_{0} \\
\psi_{2} \\
\psi_{3} \\
\vdots \\
\psi_{n-1}
\end{array}\right) .
$$

The $(n \times n)$-matrix $L$ is the circulant matrix, and the eigenvalues $\lambda_{k}$ and eigenvectors $e_{k}$ can be computed by the well known fact (cf. e.g. [11], p. 72) :

$$
\lambda_{k}=4 \sin ^{2} \frac{k \pi}{n}, \quad e_{k}=\left(1, \omega^{k}, \ldots, \omega^{k(n-1)}\right) \in \mathbb{C}^{n}, \quad \omega=\exp (2 \pi \sqrt{-1} / n)
$$

In particular, $L$ is positive semi-definite. The assumption $\sum_{k} \psi_{k}=0$ shows that the vector $\Psi$ is perpendicular to the vector ${ }^{t}(1, \ldots, 1)$ which is the eigenvector of the eigenvalue $\lambda_{n}=0$. Therefore for any $\Psi$ satisfying $\sum_{k} \psi_{k}=0$ we have

$$
\sum_{k}\left(\psi_{k+1}-\psi_{k}\right)^{2}=\langle L \Psi, \Psi\rangle \geq 4 \sin ^{2} \frac{\pi}{n}\langle\Psi, \Psi\rangle
$$

and the equality holds if and only if $\Psi=e_{1}=e_{n-1}={ }^{t}\left(1, \omega, \ldots, \omega^{n-1}\right)$.
In the following, we will consider a special type of normal variation, i.e., a variation which has the form:

$$
p_{k}(t)=p_{k}+t \psi_{k} N_{k}, \quad N_{k}=R\left(\nu_{k}-\nu_{k-1}\right) / \sin \theta_{k}
$$

By the second variation formula (Theorem 2.3.33) we have

$$
\delta^{2} \text { Length }=\sum_{k}\left(\left|\nabla \psi_{k}\right|^{2}-\Lambda^{2} \psi_{k} \psi_{k+1}\right) l_{0}=\sum_{k} \frac{1}{l_{0}}\left[\left(\psi_{k+1}-\psi_{k}\right)^{2}-4 \psi_{k} \psi_{k+1} \tan ^{2} \frac{m \pi}{n}\right]
$$

for any admissible variations, where we use the relation $\Lambda l_{0}=2 \tan \left(\theta_{0} / 2\right)$ and put $\theta_{0}=2 m \pi / n$ for some $m \in \mathbb{Z}$ and assume that $m$ and $n$ are coprime. Note that the area-preserving condition derives $\sum_{k} \psi_{k}=0$ (Lemma 2.3.30).

Theorem 2.3.35 (Instability of non-convex regular polygons). Let $n \geq 5$. By taking $\psi_{k}=A \cos (2 \pi k / n)+B \sin (2 \pi k / n),(A, B) \neq(0,0)$, we have

$$
\delta^{2} \text { Length }=\frac{4}{l_{0}}\left[\sin ^{2} \frac{\pi}{n}-\cos \frac{2 \pi}{n} \tan ^{2} \frac{m \pi}{n}\right] \sum_{k} \psi_{k}^{2}
$$

In particular, $\delta^{2}$ Length $<0$ for $2 \leq m \leq n-2$, i.e., non-convex regular polygons are unstable.

Proof. By the discrete Wirtinger's inequality we have

$$
\delta^{2} \text { Length } \geq \sum_{k} \frac{4}{l_{0}}\left(\psi_{k}^{2} \sin ^{2} \frac{\pi}{n}-\psi_{k} \psi_{k+1} \tan ^{2} \frac{m \pi}{n}\right)
$$

In the following we use the equality condition $\psi_{k}=A \cos (2 \pi k / n)+B \sin (2 \pi k / n)$ and put $\varphi_{k}=-A \sin (2 \pi k / n)+B \cos (2 \pi k / n)$. Then we have

$$
\psi_{k+1}=\psi_{k} \cos (2 \pi / n)+\varphi_{k} \sin (2 \pi / n), \quad \psi_{k} \varphi_{k}=\frac{1}{2}\left(B^{2}-A^{2}\right) \sin \frac{4 k \pi}{n}+A B \cos \frac{4 k \pi}{n} .
$$

If we note the fact $\sum_{k} \psi_{k} \varphi_{k}=0$, then

$$
\begin{aligned}
\delta^{2} \text { Length } & =\sum_{k} \frac{4}{l_{0}}\left[\psi_{k}^{2} \sin ^{2} \frac{\pi}{n}-\psi_{k}\left(\psi_{k} \cos \frac{2 \pi}{n}+\varphi_{k} \sin \frac{2 \pi}{n}\right) \tan ^{2} \frac{m \pi}{n}\right] \\
& =\frac{4}{l_{0}}\left[\sin ^{2} \frac{\pi}{n}-\cos \frac{2 \pi}{n} \tan ^{2} \frac{m \pi}{n}\right] \sum_{k} \psi_{k}^{2} .
\end{aligned}
$$

If $m=1$ or $m=n-1$, then

$$
\delta^{2} \text { Length } \geq \frac{4}{l_{0}} \sin ^{2} \frac{\pi}{n} \tan ^{2} \frac{\pi}{n} \sum_{k} \psi_{k}^{2} \geq 0 .
$$

On the other hand, for $2 \leq m \leq n-2$ and $(A, B) \neq(0,0)$ we have

$$
\delta^{2} \text { Length } \leq-\frac{4 \sin ^{2}(\pi / n)\left(1+2 \cos ^{2}(\pi / n)\right)}{l_{0} \cos (2 \pi / n)} \sum_{k} \psi_{k}^{2}<0,
$$

where we use the fact $\tan ^{2}(m \pi / n) \geq \tan ^{2}(2 \pi / n)$ if $2 \leq m \leq n-2$, and $\cos (2 \pi / n)>0$ if $n \geq 5$. This proves the statement.

### 2.4 Geometry of discrete anisotropic surface energy

### 2.4.1 Anisotropic surface energy and its first variation

Next we consider the anisotropic energy for simplicial surfaces. Let $M_{h}$ be a simplicial surface and $\left\{\nu_{T}\right\}_{T \in M_{h}^{(2)}}$ be the unit normal vector fields on $M_{h}$, where $M_{h}^{(2)}$ is the set of 2-skeltons, i.e., the set of triangles. For a continuous function $\gamma: S^{2} \rightarrow \mathbb{R}_{>0}$, we define the anisotropic surface energy of $M_{h}$ as

$$
\mathscr{F}_{\gamma}\left(M_{h}\right):=\sum_{T \in M_{h}^{(2)}} \gamma\left(\nu_{T}\right) \cdot \operatorname{Area}(T) .
$$

This functional can be considered as a function $\mathscr{F}_{\gamma}: \mathbb{R}^{3 n} \rightarrow \mathbb{R}$ as in the simplicial curve case, where $n$ is the number of vertices of $M_{h}$.

Consider a vector-valued function $v_{p_{j}} \in \mathbb{R}^{3}$ defined on the $n_{i}$ interior vertices $\mathscr{V}_{\text {int }}=$ $\left\{v_{p_{1}}, \ldots, v_{p_{n_{i}}}\right\}$ of $M_{h}$. We may extend this function to the boundary vertices of $M_{h}$ as well, by assuming $v_{p}=0 \in \mathbb{R}^{3}$ for each boundary vertex $p$. The vectors $v_{p_{j}}$ are the variation vector field of any boundary-fixing variation of the form

$$
p_{j}(t)=p_{j}+t \cdot v_{p_{j}}+O\left(t^{2}\right)
$$

that is, $p_{j}^{\prime}(0)=v_{p_{j}}$. We define the vector $\vec{v} \in \mathbb{R}^{3 n}$ by


$$
{ }^{t} \vec{v}=\left({ }^{t} v_{p_{1}}, \ldots,{ }^{t} v_{p_{n}}\right)
$$

If we modify this vector as a $(3 \times n)$-matrix $\left(v_{p_{1}}, \ldots, v_{p_{n}}\right)$, each row vector can be considered as an element in $\mathscr{L}_{0}$, where $\mathscr{L}_{0}$ is the set of conforming functions on $M_{h}$ which vanish on the boundary:

$$
\mathscr{L}_{0}:=\left\{f: M_{h} \rightarrow \mathbb{R}: f \in C^{0}\left(M_{h}\right), f \text { is linear on each } T \in M_{h} \text { and }\left.f\right|_{\partial M_{h}}=0\right\} .
$$

Now we derive the first variation formula for the anisotropic surface energy. We assume that $\gamma$ is of class $C^{1}$ in the following. That means we would like to find the vectors $\nabla_{p_{j}} \mathscr{F}_{\gamma} \in \mathbb{R}^{3}$ which satisfies

$$
\frac{d}{d t} \mathscr{F}_{\gamma=0}=\left\langle\vec{v}, \nabla \mathscr{F}_{\gamma}\right\rangle_{\mathbb{R}^{3 n}}=\sum_{j=1}^{n}\left\langle v_{p_{j}}, \nabla_{p_{j}} \mathscr{F}_{\gamma}\right\rangle_{\mathbb{R}^{3}},
$$

where we denote ${ }^{t} \nabla \mathscr{F}_{\gamma}=\left({ }^{t} \nabla_{p_{1}} \mathscr{F}_{\gamma}, \ldots,{ }^{t} \nabla_{p_{n}} \mathscr{F}_{\gamma}\right) \in \mathbb{R}^{3 n}$. In other words, we will calculate the directional derivative of the function $\mathscr{F}_{\gamma}$ in $\mathbb{R}^{3 n}$.

Theorem 2.4.1. Let $\gamma: S^{2} \rightarrow \mathbb{R}_{>0}$ be a $C^{1}$ function and $p$ be an interior vertex of a simplicial surface $M_{h}$. Then at each interior vertex $p$ the anisotropic energy gradient can be expressed in the following formula:

$$
\begin{equation*}
\left.\nabla_{p} \mathscr{F}_{\gamma}=\frac{1}{2} \sum_{j=1}^{k} \xi_{\gamma}\left(\nu_{j}\right) \times\left(q_{j+1}-q_{j}\right)=\frac{1}{2} \sum_{j=1}^{k}\left(\xi_{\gamma}\left(\nu_{j-1}\right)-\xi_{\gamma}\left(\nu_{j}\right)\right)\right) \times\left(q_{j}-p\right), \tag{4.1}
\end{equation*}
$$

where we write $\operatorname{star}(p)=\left\{T_{1}, \ldots, T_{k}\right\}, T_{j}=\left(p, q_{j}, q_{j+1}\right), \nu_{j}:=\nu_{T_{j}}$, and $\xi_{\gamma}: S^{2} \rightarrow \mathbb{R}^{3}$ is the Cahn-Hoffman map for $\gamma$.

Remarks . (1) If we take $\gamma \equiv 1$, this formula reduces to the well-known formula (so-called cotangent formula) in [34].
(2) We can define this quantity on any polyhedral surface, i.e., the quantity (4.1) is independent of the choice of the triangulation of polygons. Because the triangles lie on the same plane, they have the same unit normal. In general, we can treat the same way if the polyhedral surface contains a simply-connected subsurface lies in a plane (cf. [37], Remark 2.2).
(3) We can define the energy gradient at the edge as in the isotropic case [35], i.e., if we take an interior edge $e=q-p$ and denote the unit normal of adaject triangles $T_{1}=\left(p, r_{1}, q\right)$ and $T_{2}=\left(p, q, r_{2}\right)$ as $\nu_{1}, \nu_{2}$, then

$$
\nabla_{e} \mathscr{F}_{\gamma}:=\frac{1}{2}\left(\xi_{\gamma}\left(\nu_{1}\right)-\xi_{\gamma}\left(\nu_{2}\right)\right) \times e .
$$

The proof of well-definedness is completely the same as [35], Lemma 83. From this viewpoint we can consider that the bending of a surface happens at edges instead of vertices.

Lemma 2.4.2. For a triangle $T=(p, q, r)$ we have

$$
\begin{aligned}
\delta \operatorname{Area}(T) & =\frac{1}{2}\left[\left\langle\nu \times(r-q), v_{p}\right\rangle+\left\langle\nu \times(p-r), v_{q}\right\rangle+\left\langle\nu \times(q-p), v_{r}\right\rangle\right] \\
\delta \nu & =\frac{1}{2 \operatorname{Area}(T)}\left[(r-q) \times v_{p}+(p-r) \times v_{q}+(q-p) \times v_{r}-2(\delta \operatorname{Area}(T)) \nu\right]
\end{aligned}
$$

where we write $\nu=\nu_{T}$.
Remark. Note that the first variation of the unit normal $\delta \nu=\delta \nu_{T}$ is tangent to the triangle $T$.

Proof. We put $\widetilde{\nu}=(q-p) \times(r-p)$. Then by differentiating $4 \operatorname{Area}(T)^{2}=\langle\widetilde{\nu}, \widetilde{\nu}\rangle$, we have

$$
\begin{aligned}
& 4 \cdot 2 \operatorname{Area}(T) \delta \operatorname{Area}(T)=2\langle\delta[(q-p) \times(r-p)],(q-p) \times(r-p)\rangle \\
& \quad \Longleftrightarrow \delta \operatorname{Area}(T)=\frac{1}{2}\langle\delta[(q-p) \times(r-p)], \nu\rangle
\end{aligned}
$$

Then we have

$$
\begin{aligned}
\delta \operatorname{Area}(T) & =\frac{1}{2}\left\langle\left(v_{q}-v_{p}\right) \times(r-p)+(q-p) \times\left(v_{r}-v_{p}\right), \nu\right\rangle \\
& =\frac{1}{2}\left\langle\nu,(r-q) \times v_{p}+(p-r) \times v_{q}+(q-p) \times v_{r}\right\rangle \\
& =\frac{1}{2}\left[\left\langle\nu \times(r-q), v_{p}\right\rangle+\left\langle\nu \times(p-r), v_{q}\right\rangle+\left\langle\nu \times(q-p), v_{r}\right\rangle\right]
\end{aligned}
$$

The second statement follows from the fact

$$
\delta \nu=\frac{1}{\|\widetilde{\nu}\|}(\delta \widetilde{\nu}-2(\delta \operatorname{Area}(T)) \nu)=\frac{1}{2 \operatorname{Area}(T)}(\delta \widetilde{\nu}-2(\delta \operatorname{Area}(T)) \nu)
$$

Proof of Theorem 2.4.1. By using the lemma and the fact $D \gamma(\nu) \perp \nu$, on the triangle $T=(p, q, r)$ we have

$$
\begin{aligned}
\delta\left(\gamma\left(\nu_{T}\right) \operatorname{Area}(T)\right)= & \langle D \gamma(\nu), \delta \nu\rangle \operatorname{Area}(T)+\gamma\left(\nu_{T}\right) \delta \operatorname{Area}(T) \\
= & \frac{1}{2}\left[\left\langle D \gamma(\nu),(r-q) \times v_{p}+(p-r) \times v_{q}+(q-p) \times v_{r}\right\rangle\right. \\
& \left.\quad+\gamma(\nu)\left(\left\langle\nu \times(r-q), v_{p}\right\rangle+\left\langle\nu \times(p-r), v_{q}\right\rangle+\left\langle\nu \times(q-p), v_{r}\right\rangle\right)\right] \\
= & \frac{1}{2}\left(\left\langle\xi_{\gamma}(\nu) \times(r-q), v_{p}\right\rangle+\left\langle\xi_{\gamma}(\nu) \times(p-r), v_{q}\right\rangle+\left\langle\xi_{\gamma}(\nu) \times(q-p), v_{r}\right\rangle\right) .
\end{aligned}
$$

By summing up these terms we have

$$
\delta \mathscr{F}_{\gamma}=\sum_{p}\left\langle\frac{1}{2} \sum_{j=1}^{k} \xi_{\gamma}\left(\nu_{j}\right) \times\left(q_{j+1}-q_{j}\right), v_{p}\right\rangle,
$$

where we write the vertices of the $\operatorname{star}(p)$ as $\left\{q_{1}, \ldots, q_{k}\right\}, q_{k+1}=q_{1}, T_{j}:=\left(p, q_{j}, q_{j+1}\right)$ and $\nu_{j}:=\nu_{T_{j}}$.

Remark. We can consider this formula as a special case of the smooth version of the first variation formula. On the triangle $T_{j}$, the $\pi / 2$-rotation $R_{T_{j}}$ in $\left(N_{T_{j}}, \nu_{T_{j}}\right)$-plane can be written as

$$
R_{T_{j}}(x)=x \times \frac{\vec{c}_{j}}{\left|\overrightarrow{c_{j}}\right|}
$$

By using this fact, we have

$$
\begin{aligned}
R_{T_{j}}\left(P_{T_{j}}\left(\xi\left(\nu_{T_{j}}\right)\right)\right) & =R_{T_{j}}\left(\left\langle\xi\left(\nu_{T_{j}}\right), \nu_{T_{j}}\right\rangle \nu_{T_{j}}+\left\langle\xi\left(\nu_{T_{j}}\right), N_{T_{j}}\right\rangle N_{T_{j}}\right) \\
& =R_{T_{j}}\left(\gamma\left(\nu_{T_{j}}\right) \nu_{T_{j}}+\left\langle D \gamma\left(\nu_{T_{j}}\right), N_{T_{j}}\right\rangle N_{T_{j}}\right) \\
& =\gamma\left(\nu_{T_{j}}\right) \nu_{T_{j}} \times \frac{\vec{c}_{j}}{\left|\vec{c}_{j}\right|}+\left\langle D \gamma\left(\nu_{T_{j}}\right), N_{T_{j}}\right\rangle \nu_{T_{j}} \\
& =\frac{1}{\left|\vec{c}_{j}\right|}\left(\gamma\left(\nu_{T_{j}}\right) \nu_{T_{j}} \times \vec{c}_{j}+\left\langle D \gamma\left(\nu_{T_{j}}\right), \vec{c}_{j} \times \nu_{T_{j}}\right\rangle \nu_{T_{j}}\right) \\
& =\frac{1}{\left|\overrightarrow{c_{j}}\right|}\left(\gamma\left(\nu_{T_{j}}\right) \nu_{T_{j}} \times \vec{c}_{j}+\left\langle D \gamma\left(\nu_{T_{j}}\right) \times \vec{c}_{j}, \nu_{T_{j}}\right\rangle \nu_{T_{j}}\right) \\
& =\frac{1}{\left|\overrightarrow{c_{j}}\right|}\left(\gamma\left(\nu_{T_{j}}\right) \nu_{T_{j}} \times \vec{c}_{j}+D \gamma\left(\nu_{T_{j}}\right) \times \vec{c}_{j}\right) \\
& =\frac{1}{\left|\overrightarrow{c_{j}}\right|}\left(D \gamma\left(\nu_{T_{j}}\right)+\gamma\left(\nu_{T_{j}}\right) \nu_{T_{j}}\right) \times \vec{c}_{j}=\frac{1}{\left|\vec{c}_{j}\right|} \xi\left(\nu_{T_{j}}\right) \times \vec{c}_{j}
\end{aligned}
$$

Here we use the fact $D \gamma\left(\nu_{T_{j}}\right), \vec{c}_{j} \perp \nu_{T_{j}}$ (hence $D \gamma\left(\nu_{T_{j}}\right) \times \vec{c}_{j} \| \nu_{T_{j}}$ ). Therefore we have

$$
\sum_{j=1}^{k}\left\langle\vec{c}_{j} \times \xi\left(\nu_{T_{j}}\right), v\right\rangle=\sum_{j=1}^{k}\left\langle-R_{T_{j}}\left(P_{T_{j}}\left(\xi\left(\nu_{T_{j}}\right)\right)\right), v\right\rangle\left|\vec{c}_{j}\right|=\int_{\partial \operatorname{star}(p)}\langle-R(P(\xi)), v\rangle d \widetilde{s}
$$

### 2.4.2 Discrete constant anisotropic mean curvature surfaces

The volume of oriented surface $M_{h}$ is the oriented volume enclosed by the cone of the surface over the origin in $\mathbb{R}^{3}$

$$
\operatorname{Vol}\left(M_{h}\right):=\frac{1}{3} \sum_{T=(p, q, r) \in M_{h}}\left\langle p, \nu_{T}\right\rangle \cdot \operatorname{Area}(T)=\frac{1}{6} \sum_{T=(p, q, r) \in M_{h}}\langle p, q \times r\rangle,
$$

where we denote

$$
\nu_{T}=\frac{(q-p) \times(r-p)}{\|(q-p) \times(r-p)\|} .
$$

Then we have the first variation of the volume with respect to the variation vector field $\vec{v} \in \mathbb{R}^{3 n}$ by

$$
\left.\left.\frac{d}{d t} \right\rvert\, t=0\right) \operatorname{Vol}\left(M_{h}\right)=\sum_{p}\left\langle v_{p}, \nabla_{p} \mathrm{Vol}\right\rangle,
$$

where

$$
\nabla_{p} \mathrm{Vol}=\frac{1}{6} \sum_{T=(p, q, r) \in \operatorname{star}(p)} q \times r
$$

Remark. Since $\partial \operatorname{star}(p)$ is a closed curve for an interior vertex $p$, we have

$$
\sum_{T=(p, q, r) \in \operatorname{star}(p)} q \times r=\sum_{T=(p, q, r) \in \operatorname{star}(p)}(q-p) \times(r-p)
$$

Therefore we can modify the first variation of the volume as follows:

$$
\delta \mathrm{Vol}=\sum_{p}\left(\left\langle\frac{1}{3 A_{p}} \sum_{T \in \operatorname{star}(p)} \operatorname{Area}(T) \nu_{T}, v_{p}\right\rangle\right) A_{p}
$$

Then we have weighted summations of the face normals which are volume descent direction:

$$
N_{p}^{V}:=\frac{1}{3 A_{p}} \sum_{T \in \operatorname{star}(p)} \operatorname{Area}(T) \nu_{T},
$$

and it seems to be natural to take $3 A_{p}=\operatorname{Area}(\operatorname{star}(p))$. However, since there is no natural area element at the vertices, another choice can be allowed.

By virtue of the smooth case, we define the discrete constant anisotropic mean curvature surface as follows.

Definition 2.4.3 (Discrete CAMC surface (cf. [37], Definition 2.3)). Let $\gamma: S^{2} \rightarrow \mathbb{R}$ be a $C^{1}$ function. Then a simplicial surface $M_{h}$ has constant anisotropic mean curvature (CAMC) if there exists a constant $\Lambda$ such that

$$
\nabla_{p} \mathscr{F}_{\gamma}+2 \Lambda \nabla_{p} \mathrm{Vol}=0
$$

holds for all interior vertex $p$. If $\Lambda=0$, we call the surface discrete anisotropic minimal surface. When $\gamma \equiv 1$, this reduces to the definition of the discrete CMC surfaces ([37]).

Remark. As we remarked before, if $M_{h}$ is a discrete surface that contains a simplyconnected discrete subsurface $M_{h}^{\prime}$ that lies in a plane, then it follows easily from the formula (4.1) that the discrete anisotropic minimality is independent of the choice of triangulation of the trace of $M_{h}^{\prime}$ as in the isotropic case (see [37], Remark 2.2).

We have the following Minkowski-type formula.
Proposition 2.4.4 (Discrete Minkowski-type formula). Let $\gamma: S^{2} \rightarrow \mathbb{R}$ be a $C^{1}$ function and $M_{h}$ be a closed discrete CAMC- $\Lambda$ surface. Then we have the following Minkowski-type formula:

$$
\sum_{T=(p, q, r)}\left(\gamma\left(\nu_{T}\right)+\Lambda\left\langle p, \nu_{T}\right\rangle\right) \operatorname{Area}(T)=0
$$

Proof. We consider the homothety $p(t)=p+t p=(1+t) p$. Then we have $\delta \mathscr{F}_{\gamma}\left(M_{h, t}\right)=$ $2 \mathscr{F}_{\gamma}\left(M_{h}\right)$. On the other hand, by using the CAMC condition $\nabla_{p} \mathscr{F}_{\gamma}=-2 \Lambda \nabla_{p} \mathrm{Vol}$ and the first variation of the volume, we have

$$
\begin{aligned}
\delta \mathscr{F}_{\gamma}\left(M_{h, t}\right) & =\sum_{p}\left\langle\nabla_{p} \mathscr{F}_{\gamma}, p\right\rangle=-2 \Lambda \sum_{p}\left\langle\nabla_{p} \operatorname{Vol}, p\right\rangle \\
& =-\frac{\Lambda}{3} \sum_{p} \sum_{\operatorname{star}}\langle p) \ni T=(p, q, r) \\
& =-\Lambda \sum_{T=(p, q, r)}\langle p, q \times r\rangle \\
& \langle p\rangle=-2 \Lambda \sum_{T=(p, q, r)}\left\langle p, \nu_{T}\right\rangle \operatorname{Area}(T) .
\end{aligned}
$$

Therefore we conclude

$$
2 \mathscr{F}_{\gamma}\left(M_{h}\right)=-2 \Lambda \sum_{T=(p, q, r)}\left\langle p, \nu_{T}\right\rangle \operatorname{Area}(T) \Longleftrightarrow \sum_{T=(p, q, r)}\left(\gamma\left(\nu_{T}\right)+\Lambda\left\langle p, \nu_{T}\right\rangle\right) \operatorname{Area}(T)=0
$$

Remark. The following question is naturally arised from the observation in the planar curve case: can we prove the relation $(p+q+r) / 3=-(1 / \Lambda) \xi_{\gamma}\left(\nu_{T}\right)$ on every triangle $T=(p, q, r)$ ? The answer is trivially false since we have discrete cylinders (Example 2.4.9). However, the following theorem is known in convex geometry [13].

Theorem 2.4.5 ([28], [29]). Of all convex polyhedra with a fixed surface area and the same number of faces, the best polyhedron with the largest volume is circumscribed about a sphere which touches each face in its barycenter.

Some calculation shows the following relation:

$$
\nabla_{p} \mathscr{F}_{\gamma}+2 \Lambda \nabla_{p} \mathrm{Vol}=\frac{1}{2} \sum_{T=(p, q, r) \in \operatorname{star}(p)}\left(\xi_{\gamma}\left(\nu_{T}\right)+\frac{\Lambda}{3}(p+q+r)\right) \times(r-q),
$$

Therefore, if every barycenter of the triangles touches to the rescaled Cahn-Hoffman map $-(1 / \Lambda) \xi_{\gamma}$, then the surface is discrete CAMC- $\Lambda$. In particular, the regular tetrahedron, regular octahedron and regular icosahedron are discrete CMC surfaces as remarked in [35].

(a) Regular tetrahedron

(b) Regular octahedron

(c) Regular icosahedron

Figure 2.10: Discrete CMC surfaces
We also remark the following fact:

$$
\gamma\left(\nu_{T}\right)+\Lambda\left\langle p, \nu_{T}\right\rangle=0 \Longleftrightarrow-\frac{1}{\Lambda} \xi_{\gamma}\left(\nu_{T}\right) \in \operatorname{span}\{p, q, r\}
$$

for any triangle $T=(p, q, r)$. Moreover, $\gamma\left(\nu_{T}\right)+\Lambda\left\langle p, \nu_{T}\right\rangle=0$ for every triangle $T=(p, q, r)$ if and only if the vector field $X_{T}:=\xi_{\gamma}\left(\nu_{T}\right)+(\Lambda / 3)(p+q+r)$ defines a piecewise constant vector field on the simplicial surface. For example, a discrete cylinder (Example 2.4.9) satisfies this condition with $\gamma \equiv 1$ even though the triangles of the surface does not necessarily touch to the sphere.

Another viewpoint is to consider the energy gradient on edges. As we remarked before, we can define the energy gradient on the edge $e=q-p$ as follows:

$$
\nabla_{e} \mathscr{F}_{\gamma}:=\frac{1}{2}\left(\xi_{\gamma}\left(\nu_{1}\right)-\xi_{\gamma}\left(\nu_{2}\right)\right) \times e, \quad \nabla_{e} \mathrm{Vol}=\frac{1}{6}\left(r_{1}-r_{2}\right) \times e,
$$

where we write star $(e)=\left\{T_{1}, T_{2}\right\}, T_{1}=\left(p, r_{1}, q\right), T_{2}=\left(p, q, r_{2}\right)$. Then we have
$\nabla_{e} \mathscr{F}_{\gamma}-\Lambda \nabla_{e} \mathrm{Vol}=\frac{1}{2}\left[\left(\xi_{\gamma}\left(\nu_{1}\right)+\frac{\Lambda}{3}\left(p+q+r_{1}\right)\right)-\left(\xi_{\gamma}\left(\nu_{2}\right)+\frac{\Lambda}{3}\left(p+q+r_{2}\right)\right)\right] \times e$.
Therefore, if we assume $\nabla_{e} \mathscr{F}_{\gamma}=\Lambda \nabla_{e}$ Vol for every interior edge $e$, then

$$
\omega(e):=\frac{1}{2}\left(\xi_{\gamma}(\nu)+\frac{\Lambda}{3}(p+q+r)\right) \times e
$$

is a well-defined discrete 1 -form, where $T=(p, q, r)$ is either triangle of star $(e)$.

### 2.4.3 Examples of discrete CAMC surfaces

Examples of discrete anisotropic minimal surfaces:
Example 2.4.6 (Discrete anisotropic Plateau-type problem). For fixed points

$$
q_{1}=(-1,0,-1), \quad q_{2}=(0,-1,0), \quad q_{3}=(1,0,-1), \quad q_{4}=(0,1,0)
$$

we want to find a point $p=(0,0, h)$ which satisfies $\nabla_{p}$ Area $=0$. This is one of the simplest case of the discrete Plateau-type problem. A direct calculation shows

$$
\nabla_{p} \text { Area }=\frac{2(1+2 h)}{\sqrt{2 h^{2}+2 h+2}}(0,0,1)
$$

Therefore we have $\nabla_{p}$ Area $=0 \Longleftrightarrow h=-1 / 2$.
Next we consider the anisotropic version of this surface. If we choose

$$
\gamma: S^{2} \rightarrow \mathbb{R}_{>0}, \quad \gamma\left(\nu_{1}, \nu_{2}, \nu_{3}\right):=\sqrt{a^{2} \nu_{1}^{2}+b^{2} \nu_{2}^{2}+c^{2} \nu_{3}^{2}}
$$

then the Cahn-Hoffman map for $\gamma$ is represented as

$$
\xi_{\gamma}\left(\nu_{1}, \nu_{2}, \nu_{3}\right)=\frac{1}{\sqrt{a^{2} \nu_{1}^{2}+b^{2} \nu_{2}^{2}+c^{2} \nu_{3}^{2}}}\left(a^{2} \nu_{1}, b^{2} \nu_{2}, c^{2} \nu_{3}\right)
$$

therefore the Wulff shape for $\gamma$ is the ellipsoid $(x / a)^{2}+(y / b)^{2}+(z / c)^{2}=1$ and we have

$$
\nabla_{p} \mathscr{F}_{\gamma}=\frac{2\left(\left(a^{2}+b^{2}\right) h+a^{2}\right)}{\sqrt{(1+h)^{2} a^{2}+h^{2} b^{2}+c^{2}}}(0,0,1) .
$$

Therefore $\nabla_{p} \mathscr{F}_{\gamma}=0 \Longleftrightarrow h=-a^{2} /\left(a^{2}+b^{2}\right)$. We remark that if we choose $a=$ $b=c=1$, then $\gamma \equiv 1$ and it reduces to the previous result. These surfaces are stable since we will show that a discrete CAMC surface with only one interior vertex must be a stable (Corollary 2.4.16). Although it is difficult to find the discrete anisotropic


Figure 2.11: Discrete (stable) anisotropic minimal surfaces
minimal surfaces having more refined mesh by hand, we can visualize such surfaces numerically (gradient flow method) as in the following figures.


Figure 2.12: Discrete anisotropic minimal surfaces (numerically)
Example 2.4.7 (Square anisotropic catenoid). We take $\gamma\left(\nu_{1}, \nu_{2}, \nu_{3}\right)=\left(\nu_{1}^{8}+\nu_{2}^{8}+\right.$ $\left.\nu_{3}^{8}\right)^{1 / 8}$. For constructing Schwarz P surface (triply periodic minimal surface in $\mathbb{R}^{3}$ ), we sometimes use so-called square catenoid as in the Figure 2.13 (left). From the square catenoid, we can construct a square anisotopic catenoid numerically (Figure 2.13, right).


Figure 2.13: Discrete square catenoid and square anisotropic catenoid
Example 2.4.8 (Discrete timelike catenoid). If we take $\gamma\left(\nu_{1}, \nu_{2}, \nu_{3}\right)=\sqrt{1-2 \nu_{3}^{2}},\left|\nu_{3}\right|<$ $1 / \sqrt{2}$, the energy functional $\mathscr{F}_{\gamma}$ coincides with the area functional in the LorentzMinkowski space $\left(\mathbb{R}^{3}, d x_{1}^{2}+d x_{2}^{2}-d x_{3}^{2}\right)[17]$.

With this energy, we can consider a discrete version of the timelike catenoid which rotation axis is the timelike direction as in [37]. Consider the meridian curve $(x(t), 0, t)$ as

$$
x(t)=r \cos \left(\frac{1}{r} a t\right), \quad a=\frac{r}{\delta} \arccos \left(1-\frac{1}{r^{2}} \frac{\delta^{2}}{1+\cos (2 \pi / k)}\right)
$$

Note that we have
$\lim _{\delta \rightarrow+0} \frac{r}{\delta} \arccos \left(1-\frac{1}{r^{2}} \frac{\delta^{2}}{1+\cos (2 \pi / k)}\right)=\sqrt{\frac{2}{1+\cos (2 \pi / k)}}, \quad \lim _{k \rightarrow \infty} \sqrt{\frac{2}{1+\cos (2 \pi / k)}}=1$.
Then the discrete surface whose dihedral angle is $2 \pi / k, k \in \mathbb{N}, k \geq 3$ and the profile curve $\left(x_{j}, 0, z_{j}\right)$ defined by

$$
z_{j}:=j \delta, \quad x_{j}:=x\left(z_{j}\right)
$$

is a discrete anisotropic minimal surface until the trapezoids are degenerate (Figure 2.14). The planar trapezoids of the surface may be triangulated arbitrarily by the previous remark.


Figure 2.14: Discrete timelike catenoid (avoiding singularities)

Examples of discrete (non-zero) CAMC surfaces:
Example 2.4.9 (Discrete cylinder ([35], p. 116)). Let $X: \mathbb{Z}^{2} \rightarrow \mathbb{R}^{3}$ be a map

$$
X(j, l):=(a \cos (2 \pi j m / k), a \sin (2 \pi j m / k), e l), \quad k \geq 3,1 \leq m \leq k-1, a, e>0
$$

where we assume that $m$ and $k$ are coprime. The image of the map $X$, denoted by $M_{h}$, is a "discrete cylinder". For any vertex $p$ on the meridian line we have

$$
\nabla_{p} \text { Area }=e \sqrt{2(1-\cos (2 \pi m / k))}\left(\begin{array}{l}
1 \\
0 \\
0
\end{array}\right), \quad \nabla_{p} \mathrm{Vol}=\left(\begin{array}{c}
a e \sin (2 \pi m / k) \\
0 \\
0
\end{array}\right) .
$$

Therefore if we set

$$
H:=H(k, e):=-\frac{1}{2 a} \sqrt{\frac{2}{1+\cos (2 \pi m / k)}}=-\frac{1}{2 a \cos (\pi m / k)},
$$

then we have $\nabla_{p}$ Area $+2 H \nabla_{p} \mathrm{Vol}=0$ and $M_{h}$ is discrete CMC surface by the definition. It is interesting that $H$ does not depend on the value $e$ and we saw that the value

(a) $m=1, k=17$

(b) $m=5, k=17$

(c) $m=9, k=17$

Figure 2.15: Discrete cylinders
$-1 / a \cos (\pi m / k)$ appeared as the curvature for regular $k$-gons. For the case $m=1$, when $k \rightarrow \infty$ and $e \rightarrow 0$, we have a smooth cylinder with radius $a$ and $H(k, e) \rightarrow-1 /(2 a)$.

Example 2.4.10 (Discrete unduloid-type surface ([35], p. 117)). By the similar calculation we also have a discrete unduloid-type surface. We consider the following discrete surface:

$$
X(j, l):= \begin{cases}(a \cos (2 \pi j / k), a \sin (2 \pi j / k), e l) & \text { if } l \text { is odd } \\ (b \cos (2 \pi j / k), b \sin (2 \pi j / k), e l) & \text { if } l \text { is even } .\end{cases}
$$

For any vertex $p$ on the meridian curve we have

$$
\nabla_{p} \text { Area }=\frac{2(1-c)\left(e^{2}+(1+c) a(a-b)\right)}{\sqrt{2 e^{2}(1-c)+s^{2}(b-a)^{2}}}\left(\begin{array}{l}
1 \\
0 \\
0
\end{array}\right), \quad \nabla_{p} \mathrm{Vol}=\frac{s e(2 a+b)}{3}\left(\begin{array}{l}
1 \\
0 \\
0
\end{array}\right)
$$

where we put $c=\cos (2 \pi / k)$ and $s=\sin (2 \pi / k)$. Note that if we set $a=b$ this reduces to the discrete cylinder.

Example 2.4.11 (Visualization of generating crystals). As we remarked in the preface, the anisotropic energy gives a mathematical model of small crystals. By using the energy gradient descent method, we can visualize the generation of shapes (Figure 2.16). On the other hand, although we cannot handle a non-smooth energy density by


Figure 2.16: Anisotropic energy descent $\left(\gamma\left(\nu_{1}, \nu_{2}, \nu_{3}\right)=\left(\nu_{1}^{20}+\nu_{2}^{20}+\nu_{3}^{20}\right)^{1 / 20}\right)$.
out method, we can visualize such kind of objects by specifying the Cahn-Hoffman map except a measure zero set on the unit sphere (Figure 2.17).


Figure 2.17: Anisotropic energy descent $\left(\gamma\left(\nu_{1}, \nu_{2}, \nu_{3}\right)=\left|\nu_{1}\right|+\left|\nu_{2}\right|+\left|\nu_{3}\right|\right)$.

### 2.4.4 Second variation of the anisotropic energy

Next we consider the second variation of the anisotropic energy $\mathscr{F}_{\gamma}$ of a simplicial surface $M_{h}$. We assume that $\gamma$ is of class $C^{2}$ in the following. Consider a variation of the form

$$
p_{j}(t)=p_{j}+t v_{p_{j}}+O\left(t^{2}\right), \quad j=1, \ldots, n
$$

with the variation vector field ${ }^{t} v=\left({ }^{t} v_{1}, \ldots,{ }^{t} v_{n}\right)$. We say a variation is admissible (or permissible) if the variation is volume-preserving and fixes the boundary.

Definition 2.4.12 (Stability of discrete CAMC surface). Let $\gamma: S^{2} \rightarrow \mathbb{R}$ be a $C^{2}$ function and $M_{h}$ be a discrete CAMC surface. Then $M_{h}$ is said to be stable if $\delta^{2} \mathscr{F}_{\gamma} \geq 0$ for any admissible variation.

Lemma 2.4.13 (cf. [37], Lemma 5.1). For a discrete CAMC- $\Lambda$ surface $M_{h}$ with vertex set $\mathscr{V}$, we have

$$
\delta^{2} \mathscr{F}_{\gamma}:=\frac{d^{2}}{d^{2} t}{ }_{\mid t=0} \mathscr{F}_{\gamma}=\sum_{p \in \mathscr{V}}\left\langle v_{p}, \delta\left(\nabla_{p} \mathscr{F}_{\gamma}+2 \Lambda \nabla_{p} \mathrm{Vol}\right)\right\rangle .
$$

for any admissible variation.
The proof is completely the same as Lemma 2.3.23.
We now decompose $\delta^{2} \mathscr{F}_{\gamma}$ into the sum of two terms

$$
{ }^{t} \vec{v} Q^{\gamma} \vec{v}:=\sum_{p \in \mathscr{V}}\left\langle v_{p}, \delta\left(\nabla_{p} \mathscr{F}_{\gamma}\right)\right\rangle, \quad{ }^{t} \vec{v} Q^{V} \vec{v}:=\sum_{p \in \mathscr{V}}\left\langle v_{p}, \delta\left(\nabla_{p} \mathrm{Vol}\right)\right\rangle .
$$

Then $Q^{\gamma}$ and $Q^{V}$ are $3 n \times 3 n$-matrix and

$$
\delta^{2} \mathscr{F}_{\gamma}=\left.\frac{d^{2}}{d^{2} t}\right|_{t=0} \mathscr{F}_{\gamma}={ }^{t} \vec{v}\left(Q^{\gamma}+2 \Lambda Q^{V}\right) \vec{v} .
$$

Since

$$
\begin{aligned}
\left\langle v_{p}, \delta \nabla_{p} \mathrm{Vol}\right\rangle & =\frac{1}{6} \sum_{T=\left(p, q_{j}, q_{j+1}\right) \in \operatorname{star}(p)}\left\langle v_{p}, v_{q_{j}} \times q_{j+1}+q_{j} \times v_{q_{j+1}}\right\rangle \\
& =\frac{1}{6} \sum_{j}\left\langle v_{p},\left(q_{j-1}-q_{j+1}\right) \times v_{q_{j}}\right\rangle \\
& =\frac{1}{6} \sum_{j}{ }^{t} v_{p}\left[\left(q_{j-1}-q_{j+1}\right) \times \boldsymbol{e}_{1},\left(q_{j-1}-q_{j+1}\right) \times \boldsymbol{e}_{2},\left(q_{j-1}-q_{j+1}\right) \times \boldsymbol{e}_{3}\right] v_{q_{j}},
\end{aligned}
$$

we have the following result:
Proposition 2.4.14 ([37], Proposition 5.2). The matrix $Q^{V}$ can be considered as an $n \times n$ grid with a $3 \times 3$ entry $Q_{p, q}^{V}$ for each pair of vertices $p, q \in \mathscr{V}_{\text {Int }}$ of $M_{h}$, so that

$$
{ }^{t} \vec{v} Q^{V} \vec{v}=\sum_{p \in \mathscr{Y}}\left\langle v_{p}, \delta\left(\nabla_{p} \mathrm{Vol}\right)\right\rangle
$$

for the variation vector field $\vec{v}$ of any permissible variation. We have $Q_{p, p}^{V}=0$ and $Q_{p, q}^{V}=0$ when the vertices $p$ and $q$ are not adjacent, and

$$
\begin{aligned}
Q_{p, q}^{V} & =\frac{1}{6}\left(\begin{array}{ccc}
0 & r_{2,3}-r_{1,3} & r_{1,2}-r_{2,2} \\
r_{1,3}-r_{2,3} & 0 & r_{2,1}-r_{1,1} \\
r_{2,2}-r_{1,2} & r_{1,1}-r_{2,1} & 0
\end{array}\right) \\
& =\frac{1}{6}\left(\left(r_{1}-r_{2}\right) \times \boldsymbol{e}_{1},\left(r_{1}-r_{2}\right) \times \boldsymbol{e}_{2},\left(r_{1}-r_{2}\right) \times \boldsymbol{e}_{3}\right)
\end{aligned}
$$

for adjacent unequal $p$ and $q$, where $\left(p, q, r_{k}\right)$ are the two triangles in $\operatorname{star}(\overline{p q})$ and $r_{k}=\left(r_{k, 1}, r_{k, 2}, r_{k, 3}\right)$ for $k=1,2,\left\{\boldsymbol{e}_{1}, \boldsymbol{e}_{2}, \boldsymbol{e}_{3}\right\}$ is the canonical basis of $\mathbb{R}^{3}$, and ( $p, q, r_{2}$ ) is properly oriented and ( $p, q, r_{1}$ ) is not.

Next we consider the matrix $Q^{\gamma}$. In the proof of the first variation formula (Theorem 2.4.1), we showed

$$
\frac{d}{d t} \mathscr{F}_{\gamma}=\frac{1}{2} \sum_{T=(p, q, r)}\left[\left\langle\xi_{\gamma}\left(\nu_{T}\right) \times(r-q), v_{p}\right\rangle+\left\langle\xi_{\gamma}\left(\nu_{T}\right) \times(p-r), v_{q}\right\rangle+\left\langle\xi_{\gamma}\left(\nu_{T}\right) \times(q-p), v_{r}\right\rangle\right] .
$$

Therefore we put

$$
\begin{aligned}
= & \frac{1}{2} \sum_{T=(p, q, r)}\left[\left\langle d \xi_{\gamma}\left(\delta \nu_{T}\right) \times(r-q), v_{p}\right\rangle+\left\langle d \xi_{\gamma}\left(\delta \nu_{T}\right) \times(p-r), v_{q}\right\rangle+\left\langle d \xi_{\gamma}\left(\delta \nu_{T}\right) \times(q-p), v_{r}\right\rangle\right. \\
& \left.\quad+\left\langle\xi_{\gamma}\left(\nu_{T}\right) \times\left(v_{r}-v_{q}\right), v_{p}\right\rangle+\left\langle\xi_{\gamma}\left(\nu_{T}\right) \times\left(v_{p}-v_{r}\right), v_{q}\right\rangle+\left\langle\xi_{\gamma}\left(\nu_{T}\right) \times\left(v_{q}-v_{p}\right), v_{r}\right\rangle\right] \\
= & \frac{1}{2} \sum_{T=(p, q, r)}\left(I_{T}+I I_{T}\right) .
\end{aligned}
$$

A simple calculation shows that

$$
I I_{T}=-2\left\langle\xi_{\gamma}\left(\nu_{T}\right), v_{p} \times v_{q}+v_{q} \times v_{r}+v_{r} \times v_{p}\right\rangle
$$

For a triangle $T=(p, q, r)$ if we put
$A_{T}:=(r-q) \times v_{p}+(p-r) \times v_{q}+(q-p) \times v_{r}=(p-r) \times\left(v_{q}-v_{p}\right)+(q-p) \times\left(v_{r}-v_{p}\right)$,
then we have

$$
\begin{aligned}
\delta \nu_{T} & =\frac{1}{2 \operatorname{Area}(T)}\left(A_{T}-\left\langle A_{T}, \nu_{T}\right\rangle \nu_{T}\right) \\
I_{T} & =\frac{1}{2 \operatorname{Area}(T)}\left\langle d \xi_{\gamma}\left(A_{T}-\left\langle A_{T}, \nu_{T}\right\rangle \nu_{T}\right), A_{T}\right\rangle=2 \operatorname{Area}(T)\left\langle d \xi_{\gamma}\left(\delta \nu_{T}\right), \delta \nu_{T}\right\rangle
\end{aligned}
$$

The last equality follows from the facts the first variation of the unit normal $\delta \nu_{T}$ is tangent to the triangle and $d \xi_{\gamma}\left(\delta \nu_{T}\right)$ is perpendicular to $\nu_{T}$ (since $d \xi_{\gamma}$ is the endomorphism of the tangent space).

Proposition 2.4.15. The bilinear form ${ }^{t} \vec{v} Q^{\gamma} \vec{v}$ can be represented as follows:

$$
{ }^{t} \vec{v} Q^{\gamma} \vec{v}=\sum_{T=(p, q, r)} \frac{\left\langle d \xi_{\gamma}\left(A_{T}-\left\langle A_{T}, \nu_{T}\right\rangle \nu_{T}\right), A_{T}\right\rangle}{4 \operatorname{Area}(T)}-\left\langle\xi_{\gamma}\left(\nu_{T}\right), v_{p} \times v_{q}+v_{q} \times v_{r}+v_{r} \times v_{p}\right\rangle,
$$

where we put $A_{T}:=(r-q) \times v_{p}+(p-r) \times v_{q}+(q-p) \times v_{r}$.
Although we cannot achieve to determine the explicit representation of $Q^{\gamma}$, we have the following facts: If we consider the matrix $Q^{\gamma}$ as an $n \times n$ grid with a $3 \times 3$ entry $Q_{p, q}^{\gamma}$ for each pair of interior vertices $p, q$ of $M_{h}$, then
(1) $Q_{p, q}^{\gamma}=0$ when the vertices $p$ and $p$ are not adjacent.
(2) For the diagonal entries $Q_{p, p}^{\gamma}$ we have

$$
\begin{aligned}
{ }^{t} v_{p} Q_{p, p}^{\gamma} v_{p} & =\sum_{T=(p, q, r) \in \operatorname{star}(p)} \frac{1}{4 \operatorname{Area}(T)}\left\langle d \xi_{\gamma}\left(A_{T}-\left\langle A_{T}, \nu_{T}\right\rangle \nu_{T}\right), A_{T}\right\rangle \\
& =\sum_{T=(p, q, r) \in \operatorname{star}(p)} \frac{1}{4 \operatorname{Area}(T)}\left\langle d \xi_{\gamma}\left(A_{T}-\left\langle A_{T}, \nu_{T}\right\rangle \nu_{T}\right), A_{T}-\left\langle A_{T}, \nu_{T}\right\rangle \nu_{T}\right\rangle,
\end{aligned}
$$

where we put $A_{T}=(r-q) \times v_{p}$.
(3) For adjacent vertices $p, q$ if we denote the adjacent triangles $T_{1}=\left(p, r_{1}, q\right)$ and $T_{2}=\left(p, q, r_{2}\right)$, then

$$
\begin{aligned}
&{ }^{t} v_{p} Q_{p, q}^{\gamma} v_{q}= \frac{\left\langle d \xi_{\gamma}\left(A_{1}-A_{1}\left\langle A_{1}, \nu_{1}\right\rangle \nu_{1}\right), A_{1}\right\rangle}{4 \operatorname{Area}\left(T_{1}\right)}+\frac{\left\langle d \xi_{\gamma}\left(A_{2}-A_{2}\left\langle A_{2}, \nu_{2}\right\rangle \nu_{1}\right), A_{2}\right\rangle}{4 \operatorname{Area}\left(T_{2}\right)} \\
&-\left\langle\xi_{\gamma}\left(\nu_{1}\right)-\xi_{\gamma}\left(\nu_{2}\right), v_{p} \times v_{q}\right\rangle,
\end{aligned}
$$

where we put $A_{1}=\left(r_{2}-q\right) \times v_{p}+\left(p-r_{2}\right) \times v_{q}, A_{2}=\left(q-r_{1}\right) \times v_{p}+\left(r_{1}-p\right) \times v_{q}$, $\nu_{1}:=\nu_{T_{1}}$ and $\nu_{2}:=\nu_{T_{2}}$.

Moreover, we still have the following stability result.
Corollary 2.4.16. Let $\gamma: S^{2} \rightarrow \mathbb{R}_{>0}$ be of class $C^{2}$ and convex. If a discrete CAMC surface $M_{h}$ has only one interior vertex, then it is stable.

Proof. The single interior vertex is denoted by $p$, and $\operatorname{star}(p)=M_{h}$. Then $Q^{\gamma}=Q_{p, p}^{\gamma}$ and $Q^{V}=Q_{p, p}^{V}$ are $3 \times 3$ matrices. By Propsition 2.4.14, $Q^{V}=0$ and we have

$$
{ }^{t} v_{p} Q_{p, p}^{\gamma} v_{p}=\frac{1}{4} \sum_{T=(p, q, r) \in \operatorname{star}(p)} \frac{1}{\operatorname{Area}(T)}\left\langle d \xi_{\gamma}\left(A_{T}-\left\langle A_{T}, \nu_{T}\right\rangle \nu_{T}\right), A_{T}-\left\langle A_{T}, \nu_{T}\right\rangle \nu_{T}\right\rangle .
$$

for any vector $v_{p} \in \mathbb{R}^{3}$, where we put $A_{T}:=(r-q) \times v_{p}$. Since the matrix $d \xi_{\gamma}$ is positive semidefinite on the tangent space by the assumption, we have $\delta^{2} \mathscr{F}_{\gamma} \geq 0$ for all admissible variations.

Remark. This is a generalization of the stability result in [37], Corollary 5.1. We also remark that the convexity condition for the energy density $\gamma: S^{2} \rightarrow \mathbb{R}_{>0}$ seems to be natural requirement for the stability (cf. [31], [25], [32] for strictly convex case, [24] for convex case).

## Appendix

## Examples of the anisotropic energy

The following contents are contained in [39] and we will reformulate with our notation.
In the following we assume that $X$ has a nonparametric form, i.e., $M \subset \mathbb{R}^{n}$ is a domain and $X(x)=(x, \varphi(x))$ for some smooth function $\varphi: M \rightarrow \mathbb{R}$. In this case, the unit normal vector field $\nu_{X}$ along $X$ and the area element can be written as follows:

$$
\nu_{X}=\frac{1}{\sqrt{1+\|\nabla \varphi\|^{2}}}(\nabla \varphi,-1), \quad d A_{X}=\sqrt{1+\|\nabla \varphi\|^{2}} d x
$$

We also assume that the image of the Cahn-Hoffman map $\widetilde{W}_{\gamma}$ has a nonparametric form. In this case, a local coordinate of $\widetilde{W}_{\gamma}$ is defined by using the projection

$$
\xi_{\gamma}(\nu)=\left(y_{1}, \ldots, y_{n}, y_{n+1}\right) \mapsto\left(y_{1}, \ldots, y_{n}\right) .
$$

By using this notation, we can write the anisotropic shape operator and the anisotropic mean curvature as

$$
-d \widetilde{\xi}_{\gamma}=-\left(\frac{\partial y_{i}}{\partial x_{j}}\right)_{i, j=1, \ldots, n}, \quad \Lambda=-\frac{1}{n} \sum_{i=1}^{n} \frac{\partial y_{i}}{\partial x_{i}}
$$

Example 2.4.17 (isotropic case). If we take $\Omega=\left\{\nu=\left(\nu_{1}, \ldots, \nu_{n+1}\right) \in S^{n} \mid \nu_{n+1}<0\right\}$ and $\gamma(\nu) \equiv 1$, then we have $\xi_{\gamma}(\nu)=\nu$ for $\nu \in \Omega, \widetilde{W}_{\gamma}=\Omega$ and

$$
\widetilde{\xi}_{\gamma}=\xi_{\gamma} \circ \nu_{X}=\frac{1}{\sqrt{1+\|\nabla \varphi\|^{2}}}(\nabla \varphi,-1), \quad \Lambda=-\frac{1}{n} \sum_{i=1}^{n} \frac{\partial}{\partial x_{i}}\left(\frac{\varphi_{i}}{\sqrt{1+\|\nabla \varphi\|^{2}}}\right)
$$

where we write $\varphi_{i}=\partial \varphi / \partial x_{i}$. Therefore $X$ is CAMC if and only if $X$ is a CMC hypersurface in $\mathbb{R}^{n+1}$ in the usual sense.

Example 2.4.18 (hyperbolic case). If we take $\Omega=\left\{\nu \in S^{n} \mid \nu_{n+1}<-1 / \sqrt{2}\right\}$ and $\gamma: \Omega \rightarrow \mathbb{R}_{>0}$ to be

$$
\gamma(\nu):=\sqrt{\nu_{n+1}^{2}-\nu_{1}^{2}-\cdots-\nu_{n}^{2}}=\sqrt{-1+2 \nu_{n+1}^{2}}
$$

Then we have

$$
\xi_{\gamma}(\nu)=\frac{1}{\gamma(\nu)}\left(-\nu_{1}, \ldots,-\nu_{n}, \nu_{n+1}\right)
$$

and the image of the Cahn-Hoffman map becomes the lower sheet of the two-sheeted hyperboloid $y_{n+1}^{2}-\left(y_{1}^{2}+\cdots+y_{n}^{2}\right)=1$. Since

$$
\int_{M} \gamma\left(\nu_{X}\right) d A_{X}=\int_{M} \sqrt{\frac{1-\|\nabla \varphi\|^{2}}{1+\|\nabla \varphi\|^{2}}} \cdot \sqrt{1+\|\nabla \varphi\|^{2}} d x=\int_{M} \sqrt{1-\|\nabla \varphi\|^{2}} d x
$$

the anisotropic surface energy coincides with the $n$-dimensional volume of $X$ in the Lorentz-Minkowski space $\mathbb{R}^{n, 1}$. Moreover we have

$$
\widetilde{\xi}_{\gamma}=\xi_{\gamma} \circ \nu_{X}=\frac{1}{\sqrt{1-\|\nabla \varphi\|^{2}}}(-\nabla \varphi,-1), \quad \Lambda=\frac{1}{n} \sum_{i=1}^{n} \frac{\partial}{\partial x_{i}}\left(\frac{\varphi_{i}}{\sqrt{1-\|\nabla \varphi\|^{2}}}\right)
$$

Therefore $X$ is CAMC if and only if $X$ is a spacelike constant mean curvature hypersurface in the Lorentz-Minkowski space $\mathbb{R}_{1}^{n}$.

More generally we have the following result:
Theorem 2.4.19 ([17], Theorem 1). Set $\Omega_{1}=\left\{\nu=\left(\nu_{1}, \nu_{2}, \nu_{3}\right) \in S^{2}| | \nu_{3} \mid>1 / \sqrt{2}\right\}$, $\Omega_{2}=\left\{\nu=\left(\nu_{1}, \nu_{2}, \nu_{3}\right) \in S^{2}| | \nu_{3} \mid<1 / \sqrt{2}\right\}$. Define a function $\gamma: S^{2} \rightarrow \mathbb{R}$ as

$$
\gamma\left(\nu_{1}, \nu_{2}, \nu_{3}\right):=\sqrt{\left|\nu_{3}^{2}-\nu_{1}^{2}-\nu_{2}^{2}\right|}=\sqrt{\left|2 \nu_{3}^{2}-1\right|}
$$

Then, an immersion $X: M \rightarrow \mathbb{R}^{3}$ with Gauss image $\nu(M) \subset \Omega_{1} \cup \Omega_{2}$ is anisotropic zero mean curvature if and only if the mean curvature of $X$ is zero as an immersed surface in $\mathbb{R}_{1}^{3}$.

Example 2.4.20 (parabolic case). If we take $\Omega=\left\{\nu \in S^{n} \mid \nu_{n+1}<0\right\}$ and $\gamma: \Omega \rightarrow \mathbb{R}_{>0}$ to be

$$
\gamma(\nu)=-\frac{1}{2} \frac{\nu_{1}^{2}+\cdots+\nu_{n}^{2}}{\nu_{n+1}}=\frac{1}{2}\left(\nu_{n+1}-\frac{1}{\nu_{n+1}}\right)
$$

Then we have

$$
\xi_{\gamma}(\nu)=\left(-\frac{\nu_{1}}{\nu_{n+1}}, \ldots,-\frac{\nu_{n}}{\nu_{n+1}}, \frac{1-\nu_{n+1}^{2}}{2 \nu_{n+1}^{2}}\right)
$$

and the image of the Cahn-Hoffman map of $\gamma$ becomes the paraboloid $y_{n+1}=(1 / 2)\left(y_{1}^{2}+\right.$ $\cdots+y_{n}^{2}$ ). Since we have

$$
\int_{M} \gamma\left(\nu_{X}\right) d A_{X}=\int_{M} \frac{\|\nabla \varphi\|^{2}}{2 \sqrt{1+\|\nabla \varphi\|^{2}}} d A_{X}=\frac{1}{2} \int_{M}\|\nabla \varphi\|^{2} d x
$$

the anisotropic surface energy coincides with the Dirichlet energy of $\varphi$. Moreover we have

$$
\widetilde{\xi}_{\gamma}=\xi_{\gamma} \circ \nu_{X}=\left(\nabla \varphi,\|\nabla \varphi\|^{2} / 2\right), \quad \Lambda=-\frac{1}{n} \sum_{i=1}^{n} \frac{\partial^{2} \varphi}{\partial x_{i}^{2}}, \quad H_{2}^{\gamma}=\varphi_{11} \varphi_{22}-\varphi_{12}^{2} .
$$

Therefore $X$ has zero anisotropic mean curvature if and only if $\varphi$ is harmonic on $M$. Note that this geometry is called the isotropic geometry [33].

In the paper [39], it is remarked that the relation between the energy density and its Wulff shape by using the Legendre transformation. We will give its reformulation in our notation.

Let us consider a nonparametric case and we use the same notation as before. Let $S_{-}^{n}:=\left\{\nu=\left(\nu_{1}, \ldots, \nu_{n+1}\right) \in S^{n} \mid \nu_{n+1}<0\right\}$ be the lower hemisphere of $S^{n}, \Omega \subset S_{-}^{n}$ be a domain and $\gamma: \Omega \rightarrow \mathbb{R}_{>0}$ be a positive-valued $C^{2}$ function. Moreover we assume that the image of the Cahn-Hoffman map $\xi_{\gamma}$, denoted by $M_{0}$, can be represented as a graph $M_{0}=\left\{y \in \mathbb{R}^{n} \mid w=G(y)\right\}$ for some function $G$.

Let us recall the definition of the Legendre transformation. The Legendre transformation of a function $G(y)=G\left(y_{1}, \ldots, y_{n}\right)$ is given by a function

$$
v=G^{*}\left(p_{1}, \ldots, p_{n}\right), \quad p_{i}=\frac{\partial G}{\partial y_{i}}, \quad G^{*}:=-G+\sum_{i} y_{i} p_{i}
$$

We assume that $\operatorname{det}\left(G_{y_{i} y_{j}}\right) \neq 0$ to make the map $y \mapsto p$ diffeomorphism. By the definition of the anisotropic Gauss map we have $p_{i}=\partial G / \partial y_{i}=\partial \varphi / \partial x_{i}$. Therefore it follows from the representation of the unit normal that we have

$$
G^{*}\left(p_{1}, \ldots, p_{n}\right)=G^{*}\left(-\nu_{1} / \nu_{n+1}, \ldots,-\nu_{n} / \nu_{n+1}\right)
$$

and

$$
G^{*}(p) d x=-\nu_{n+1} G^{*}\left(-\nu_{1} / \nu_{n+1}, \ldots,-\nu_{n} / \nu_{n+1}\right) d A_{X} .
$$

We will prove the coefficient of $d A_{X}$ coincide with the energy density $\gamma$ we gave first. If we choose the unit normal of $M_{0}$ as

$$
\nu=\left(\frac{\nabla G}{\sqrt{1+|\nabla G|^{2}}},-\frac{1}{\sqrt{1+|\nabla G|^{2}}}\right)=\left(\frac{\nabla \varphi}{\sqrt{1+|\nabla \varphi|^{2}}},-\frac{1}{\sqrt{1+|\nabla \varphi|^{2}}}\right)
$$

then the support function of $M_{0}$ becomes

$$
\begin{aligned}
\langle(y, G(y)), \nu\rangle_{\mathbb{R}^{n+1}} & =\frac{\langle\nabla G, y\rangle_{\mathbb{R}^{n}}-G}{\sqrt{1+|\nabla \varphi|^{2}}} \\
& =-\nu_{n+1} G^{*}(p)=-\nu_{n+1} G^{*}\left(-\nu_{1} / \nu_{n+1}, \ldots, \nu_{n} / \nu_{n+1}\right)
\end{aligned}
$$

On the other hand if we compute the support function of $M_{0}$ as the image of the Cahn-Hoffman map we have $\left\langle\xi_{\gamma}(\nu), \nu\right\rangle=\gamma(\nu)$, hence we have the desired result.

In this sense, the energy appeared in [39] (Proposition 5) can be regarded as the anisotropic energy we defined. As he mentioned there, we have the following first variation formula

$$
\delta \int G^{*}(p) d x=n \int \Lambda \delta \varphi d x
$$

for any support compact variation of $\varphi$.

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