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Abstract. In this article, linear operators satisfying anti-commutation relations are considered. It is proven that an anti-commutative type of the Glimm-Jaffe-Nelson commutator theorem follows.

1 Introduction and Main Theorem

In this article we consider the self-adjointness of linear operators satisfying anti-commutation relations. For criteria on the self-adjointness of the symmetric operator satisfying a commutation relation, there is the Glimm-Jaffe-Nelson commutator theorem. Refer to [2, 3, 4] of the original papers, and see also e.g. ([1]; Theorem 2.32, [6]; Theorem X.36). We investigate an anti-commutative type of the commutator theorem.

Theorem 1

Let H be a symmetric operator and A be a strictly positive self-adjoint operator i.e. there exists a constant $\delta_A > 0$ such that $(\Psi, A\Psi) \geq \delta_A(\Psi, \Psi)$. We assume the following conditions :

(C.1) There exists a core \mathcal{D}_0 of A such that $\mathcal{D}_0 \subset \mathcal{D}(H)$ where $\mathcal{D}(H)$ denotes the domain of H .

(C.2) There exists a constant $a > 0$ such that for all $\Psi \in \mathcal{D}_0$,

$$\|H\Psi\| \leq a\|A\Psi\|. \quad (1)$$

(C.3) There exists a constant $b > 0$ such that for all $\Psi \in \mathcal{D}_0$,

$$\left| (H\Psi, A\Psi) + (A\Psi, H\Psi) \right| \leq b\|A^{1/2}\Psi\|^2. \quad (2)$$

Then H is essentially self-adjoint on \mathcal{D}_0 .

Remark 1 Let X and Y be symmetric operators on a Hilbert space \mathcal{X} . Then the real part and the imaginary part of the inner product $(X\Psi, Y\Psi)$ for $\Psi \in \mathcal{D}(XY) \cap \mathcal{D}(YX)$ is expressed by using the commutator and the anti-commutator :

$$\operatorname{Re}(X\Psi, Y\Psi) = \frac{(\Psi, \{X, Y\}\Psi)}{2}, \quad (3)$$

and

$$\operatorname{Im}(X\Psi, Y\Psi) = \frac{(\Psi, [X, Y]\Psi)}{2}, \quad (4)$$

where $\{X, Y\} = XY + YX$ and $[X, Y] = XY - YX$. On the proof of the Glimm-Jaffe-Nelson commutator theorem, the estimate of a type of (4) is investigated. In this article we pply the estimate of a type of (3).

Remark 2 Since $\|A^{1/2}\Phi\| \leq \|A\Phi\| + \|\Phi\|$ follows for $\Phi \in \mathcal{D}(A)$, we see that D_0 is a core of $A^{1/2}$. Then from (C.2) and (C.3), it is seen that $\mathcal{D}(A) \subset \mathcal{D}(\overline{H_{\upharpoonright \mathcal{D}_0}})$ and for $\Phi \in \mathcal{D}(A)$,

$$\left| (\overline{H_{\upharpoonright \mathcal{D}_0}}\Phi, A\Phi) + (A\Phi, \overline{H_{\upharpoonright \mathcal{D}_0}}\Phi) \right| \leq b\|A^{1/2}\Phi\|^2 \quad (5)$$

hold, where \overline{X} denotes the closure of a operator X .

(Proof of Theorem 1)

Let $c \in \mathbf{R}$ be a real number satisfying $c > \frac{b}{2}$, and let us set $z = c + i \in \mathbf{C}$. For a closable operator X , it is seen that $\overline{X} = (X^*)^*$ and $(\overline{X})^* = X^*$ follow. Then from the general theorem ([6], Theorem X.1) on the essential-self-adjointness of closed symmetric operator, it is enough to show that $\dim \ker \left(H_{\upharpoonright \mathcal{D}_0}^* + z \right) = \dim \ker \left(H_{\upharpoonright \mathcal{D}_0}^* + z^* \right) = 0$. Let $\Psi \in \mathcal{D}_0$. It is noted that $A^{-\alpha}$, $\alpha > 0$ is bounded, since A is a strictly positive. Then by using $\operatorname{Re}(f, g) = \frac{(f, g) + (g, f)}{2}$ and $(H_{\upharpoonright \mathcal{D}_0}^*)^* = \overline{H_{\upharpoonright \mathcal{D}_0}}$, we see that for $\Xi = A^{-1}\Psi \in \mathcal{D}(A)$,

$$\operatorname{Re} \left(\Xi, (H_{\upharpoonright \mathcal{D}_0}^* + z)\Psi \right) = \frac{1}{2} \left((\overline{H_{\upharpoonright \mathcal{D}_0}}\Xi, A\Xi) + (A\Xi, \overline{H_{\upharpoonright \mathcal{D}_0}}\Xi) \right) + c\|A^{1/2}\Xi\|^2.$$

Then by using (5) and $\|A^{1/2}\Xi\| = \|A^{-1/2}\Psi\|$, we have

$$\left| \operatorname{Re} \left(\Xi, (H_{\upharpoonright \mathcal{D}_0}^* + z)\Psi \right) \right| \geq \left(c - \frac{b}{2} \right) \|A^{-1/2}\Psi\|. \quad (6)$$

Now let us assume that $\Psi \in \ker \left(H_{\upharpoonright \mathcal{D}_0}^* + z \right)$. Then, we see from (6) that $A^{-1/2}\Psi = 0$. Then by acting $A^{1/2}$ the both side of this equation, we obtain $\Psi = 0$. Hence $\ker \left(H_{\upharpoonright \mathcal{D}_0}^* + z \right) = \{0\}$. Similarly, we can prove that $\ker \left(H_{\upharpoonright \mathcal{D}_0}^* + z^* \right) = \{0\}$. Thus the proof is obtained. ■

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