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<https://hdl.handle.net/2324/1462320>

出版情報 : MI Preprint Series. 2014-09, 2014-09. 九州大学大学院数理学研究院
バージョン :
権利関係 :



MI Preprint Series

Mathematics for Industry
Kyushu University

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MI 2014-9

(Received September 1, 2014)

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Spectral properties of the linearized semigroup of the compressible Navier-Stokes equation on a periodic layer

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Abstract

The linearized problem for the compressible Navier-Stokes equation around a given constant state is considered in a periodic layer of \mathbb{R}^n with $n \geq 2$, and spectral properties of the linearized semigroup is investigated. It is shown that the linearized operator generates a C_0 -semigroup in L^2 over the periodic layer and the time-asymptotic leading part of the semigroup is given by a C_0 -semigroup generated by an $n - 1$ dimensional elliptic operator with constant coefficients that are determined by solutions of a Stokes system over the basic period domain.

1 Introduction

This paper is concerned with the initial boundary value problem for the following compressible Navier-Stokes equation in a periodic layer Ω :

$$\left\{ \begin{array}{l} \partial_t \rho + \operatorname{div}(\rho u) = 0, \\ \rho(\partial_t v + v \cdot \nabla v) - \mu \Delta v - (\mu + \mu') \nabla \operatorname{div} v + \nabla(P(\rho)) = 0, \\ v|_{\partial\Omega} = 0, \\ (\rho, v)|_{t=0} = (\rho_0, v_0). \end{array} \right. \quad (1.1)$$

Here $\rho = \rho(x, t)$ and $v = {}^\top(v^1(x, t), \dots, v^n(x, t))$ denote the unknown density and velocity, respectively, at time t and position x ; Ω is a periodic layer defined by:

$$\Omega := \{x = (x', x_n); x' \in \mathbb{R}^{n-1}, \omega_1(x') < x_n < \omega_2(x')\},$$

where ω_1 and ω_2 are nonconstant and smooth functions of x' satisfying the periodicity conditions $\omega_j(x' + \frac{2\pi}{\alpha_k} \mathbf{e}'_k) = \omega_j(x')$ ($j = 1, 2$; $k = 1, \dots, n-1$) with constants $\alpha_k > 0$ and $\mathbf{e}'_k = {}^\top(0, \dots, \overset{k}{1}, \dots, 0) \in \mathbb{R}^{n-1}$; μ and μ' are the viscosity coefficients that are constants satisfying

$$\mu > 0, \quad \frac{2}{n}\mu + \mu' \geq 0;$$

P is the pressure for which we assume that P is a smooth function of ρ that satisfies

$$P'(\rho_*) > 0$$

for a given positive constant ρ_* . Here and in what follows, ${}^\top \cdot$ stands for the transposition.

We are interested in the large time behavior of solutions to (1.1) around the constant equilibrium $u_s = {}^\top(\rho_*, 0)$. To establish a detailed asymptotic description of large time behavior, we here study spectral properties of the linearized semigroup for (1.1) around u_s as a first step of the analysis.

The system of equations for the perturbation is written as

$$\begin{cases} \partial_t \phi + \gamma \operatorname{div} w = f^0, \\ \partial_t w - \nu \Delta w - \tilde{\nu} \nabla \operatorname{div} w + \gamma \nabla \phi = f, \\ w|_{\partial\Omega} = 0, \\ u|_{t=0} = u_0 = (\phi_0, w_0). \end{cases} \quad (1.2)$$

Here $u = {}^\top(\phi, w)$ with $\phi = \frac{1}{\rho_*}(\rho - \rho_*)$ and $w = \frac{1}{\gamma}v$ denotes the (scaled) perturbation from $u_s = {}^\top(\rho_*, 0)$; ν , $\tilde{\nu}$ and γ are parameters given by

$$\nu = \frac{\mu}{\rho_*}, \quad \tilde{\nu} = \frac{\mu + \mu'}{\rho_*}, \quad \gamma = \sqrt{P'(\rho_*)};$$

and f^0 and f denote the nonlinearities

$$f^0 = -\gamma \operatorname{div}(\phi w),$$

$$f = \frac{\gamma \phi}{\rho_*(1 + \phi)} \{ \nu \Delta w + \tilde{\nu} \nabla \operatorname{div} w \} - \gamma^2 w \cdot \nabla w - \left\{ \frac{1}{\rho_*(1 + \phi)} \nabla(P(\rho_*(1 + \phi))) - \frac{P'(\rho_*)}{\rho_*} \nabla \phi \right\}.$$

Large time behavior of solutions to the compressible Navier-Stokes equations has been extensively studied since the pioneering works by Matsumura-Nishida [16, 17, 18]. See, e.g., [5, 10, 11, 13, 14, 15, 20] and references therein. In [7, 8, 9], the stability of u_s was studied when the underlying domain is an n dimensional infinite layer:

$$\mathbb{R}^{n-1} \times (0, 1) = \{x = (x', x_n) ; x' = (x_1, \dots, x_{n-1}) \in \mathbb{R}^{n-1}, 0 < x_n < 1\}.$$

It was proved that u_s is stable under sufficiently small initial perturbations and the L^2 norm of the perturbation decays in the order $t^{-\frac{n-1}{4}}$ as $t \rightarrow \infty$. Furthermore, it

was shown that the perturbation behaves like a solution of an $n - 1$ dimensional heat equation. In this paper we extend the results on the asymptotic behavior of the linearized semigroup for (1.2) obtained in [7, 8] to the case of the periodic layer Ω . We will prove that the linearized semigroup behaves as $t \rightarrow \infty$ like a semigroup generated by an $n - 1$ dimensional elliptic operator with constant coefficients. More precisely, we consider the linear problem

$$\partial_t u + Lu = 0, \quad u|_{t=0} = u_0, \quad (1.3)$$

where $u = {}^\top(\phi, w)$ is the unknown; $u_0 = {}^\top(\phi_0, w_0)$ is a given initial datum; and L is the operator of the form

$$L = \begin{pmatrix} 0 & \gamma \operatorname{div} \\ \gamma \nabla & -\nu \Delta - \tilde{\nu} \nabla \operatorname{div} \end{pmatrix}.$$

It is shown that $-L$ generates a contraction C_0 -semigroup e^{-tL} on $L^2(\Omega)$ and e^{-tL} is decomposed as

$$e^{-tL} = e^{-tL} \Pi + e^{-tL} (I - \Pi).$$

Here Π is a bounded projection on $L^2(\Omega)$; and it holds that

$$\|e^{-tL} \Pi u_0\|_{L^2(\Omega)} \leq C(1+t)^{-\frac{n-1}{4}} \|u_0\|_{L^1(\Omega)},$$

$$\|e^{-tL} (I - \Pi) u_0\|_{L^2(\Omega)} \leq C e^{-\beta t} \|u_0\|_{L^2(\Omega)}$$

and

$$\|e^{-tL} \Pi u_0 - [e^{-tH} \sigma_0] u^{(0)}\|_{L^2(\Omega)} \leq C t^{-\frac{n-1}{4} - \frac{1}{2}} \|u_0\|_{L^1(\Omega)}, \quad (1.4)$$

where β is a positive constant; and e^{-tH} is a C_0 -semigroup in $L^2(\mathbb{R}^{n-1})$ generated by the operator $-H$:

$$H\sigma = -\frac{\gamma^2}{\nu} \sum_{i,j=1}^{n-1} a_{ij} \partial_{x_i} \partial_{x_j} \sigma.$$

Here (a_{ij}) is a positive definite symmetric matrix with constant components; and σ_0 and $u^{(0)}$ are given by

$$\sigma_0 = \frac{|Q|}{|\Omega_{per}|} \int_{\omega_1(x')}^{\omega_2(x')} \phi_0(x', x_n) dx_n, \quad u^{(0)} = {}^\top(1, 0),$$

where Ω_{per} is the basic period domain given by

$$\Omega_{per} = \{x = (x', x_n); x' \in Q, \omega_1(x') < x_n < \omega_2(x')\}$$

with the basic period cell $Q = \prod_{j=1}^{n-1} [-\frac{\pi}{\alpha_j}, \frac{\pi}{\alpha_j})$. Here and in what follows, for a bounded domain D , $|D|$ denotes the volume of D . We note that the matrix (a_{ij}) is given by

$$a_{ij} = \frac{1}{|\Omega_{per}|} (\nabla w^{(i)}, \nabla w^{(j)})_{L^2(\Omega_{per})},$$

where $w^{(k)}$ ($k = 1, \dots, n-1$) are functions Q -periodic in x' satisfying the following Stokes system:

$$\begin{cases} \operatorname{div} w^{(k)} = 0, \\ -\Delta w^{(k)} + \nabla \phi^{(k)} = \mathbf{e}_k, \\ w^{(k)}|_{x_n=\omega_1(x'), \omega_2(x')} = 0 \end{cases}$$

for some $\phi^{(k)} = \phi^{(k)}(x', x_n)$ being Q -periodic in x' . Here $\mathbf{e}_k = {}^\top(0, \dots, \overset{k}{1}, \dots, 0) \in \mathbb{R}^n$.

We will prove our results as follows. In the case of infinite layers analyzed in [7, 8, 9], the spectral properties of the linearized semigroup were investigated by using the Fourier transform in $x' \in \mathbb{R}^{n-1}$. In the case of the periodic layer Ω , the Fourier transform does not work well any longer, instead, we employ the Bloch wave decomposition which transforms the linearized problem (1.3) on Ω to the problem $\partial_t u + L_{\eta'} u = 0$ on Ω_{per} under Q -periodic boundary conditions in x' . Here $L_{\eta'}$ is the linear operator obtained by replacing the partial derivatives ∂_{x_j} ($j = 1, \dots, n-1$) in L by $\partial_{x_j} + i\eta_j$ with parameter $\eta' = (\eta_1, \dots, \eta_{n-1}) \in Q^*$, where Q^* is the dual cell defined by $Q^* = \prod_{j=1}^{n-1} [-\frac{\alpha_j}{2}, \frac{\alpha_j}{2})$.

When $|\eta'| \ll 1$, the operator $L_{\eta'}$ can be regarded as a perturbation of L_0 ; and the analytic perturbation theory is applied to show that

$$\rho(-L_{\eta'}) \supset \{\operatorname{Re} \lambda > -\beta_0\} \setminus \{\lambda_{\eta'}\} \text{ for some } \beta_0 > 0,$$

$$\sigma(-L_{\eta'}) \cap \{|\lambda| < \frac{\beta_0}{2}\} = \{\lambda_{\eta'}\},$$

where

$$\lambda_{\eta'} = -\frac{\gamma^2}{\nu} \sum_{i,j=1}^{n-1} a_{ij} \eta_i \eta_j + O(|\eta'|^3)$$

as $\eta' \rightarrow 0$. It then follows that this part of e^{-tL} behaves as in (1.4). As for the remaining part of η' , we establish some estimates for a modified Stokes system (see section 4.3); and based on the established estimates we prove by an energy method that if $|\eta'| \geq r_0$ ($\eta' \in Q^*$), then

$$\rho(-L_{\eta'}) \supset \{\operatorname{Re} \lambda \geq -\beta_1\} \text{ for some } \beta_1 > 0,$$

and hence, this part of e^{-tL} decays exponentially. We note that we consider the linearized operator L as an operator on L^2 as in [6], which is in contrast to [7, 8] where the underlying space for the linearized operator is $H^1 \times L^2$. The L^2 setting will be useful for the stability analysis of stationary flows with nonzero velocity fields.

This paper is organized as follows. In section 2 we introduce some notations, function spaces and state some properties of the Bloch wave decomposition. In section 3 we state the main result of this paper. The proof of the main result is given in sections 4–5. In section 6 we give an outline of the proof of a lemma used in section 4.3.

2 Preliminaries

In this section we introduce the notations, function spaces and operators which will be used in this paper.

For a domain D and $1 \leq p \leq \infty$, the Lebesgue space over D is denoted by $L^p(D)$ and its norm is denoted by $\|\cdot\|_{L^p(D)}$. The symbol $W^{l,p}(D)$ stands for the l th order L^p Sobolev space and its norm is denoted by $\|\cdot\|_{W^{l,p}(D)}$. When $p = 2$, we denote $W^{l,2}(D)$ by $H^l(D)$ and its norm is denoted by $\|\cdot\|_{H^l(D)}$. We denote by $C_0^l(D)$ the set of all C^l function whose support is compact in D . The completion of $C_0^l(D)$ in $W^{l,p}(D)$ is denoted by $W_0^{l,p}(D)$. In particular, we write $W_0^{l,2}(D)$ as $H_0^l(D)$.

We simply denote by $L^p(D)$ the set of all vector fields $W = {}^\top(w^1, \dots, w^n)$ on D whose components w^j ($j = 1, \dots, n$) belong to $L^p(D)$ and the norm is also denoted by $\|\cdot\|_{L^p(D)}$ if no confusion will occur. Similarly, the symbols $W^{l,p}(D)$ and $H^l(D)$ are used for vector fields.

For $u = {}^\top(\phi, w)$ with $\phi \in W^{k,p}(D)$ and $w = {}^\top(w^1, \dots, w^n) \in W^{l,q}(D)$, we define the norm $\|u\|_{W^{k,p}(D) \times W^{l,q}(D)}$ by

$$\|u\|_{W^{k,p}(D) \times W^{l,q}(D)} = \|\phi\|_{W^{k,p}(D)} + \|w\|_{W^{l,q}(D)}.$$

We define the sets Q , Q^* , Ω_{per} , $\Sigma_{j,\pm}$ ($j = 1, \dots, n-1$) and Σ_n as follows:

$$\begin{aligned} Q &:= \prod_{i=1}^{n-1} \left[-\frac{\pi}{\alpha_i}, \frac{\pi}{\alpha_i}\right), \quad Q^* := \prod_{i=1}^{n-1} \left[-\frac{\alpha_i}{2}, \frac{\alpha_i}{2}\right), \\ \Omega_{per} &:= \{x = (x', x_n); x' \in Q, \omega_1(x') < x_n < \omega_2(x')\}, \\ \Sigma_{j,\pm} &:= \left\{x \in \overline{\Omega_{per}}; x_j = \pm \frac{\pi}{\alpha_j}\right\}, \\ \Sigma_n &:= \{x \in \partial\Omega; x' \in Q, x_n = \omega_j(x') \text{ } j = 1, 2\}. \end{aligned}$$

In the case $D = \Omega_{per}$, we simply write $L^p(\Omega_{per})$ as L^p , and likewise, $W^{k,p}(\Omega_{per})$, $H^l(\Omega_{per})$ as $W^{k,p}$, H^l , respectively. Similarly, the norms are also abbreviated to $\|\cdot\|_{H^l}$, $\|\cdot\|_{W^{k,p}}$, and, in particular, we write $\|\cdot\|_{L^p(\Omega_{per})}$ as $\|\cdot\|_p$.

The inner product of $L^2(D)$ is defined by

$$(f, g)_{L^2(D)} = \int_D f(x) \overline{g(x)} dx, \quad f, g \in L^2(D).$$

When $D = \Omega_{per}$, we abbreviate it to (f, g) . The dual space of $H_0^1(D)$ is denoted by $H^{-1}(D)$, and the pairing between $H^{-1}(D)$ and $H_0^1(D)$ is written as $[\cdot, \cdot]$. For $f \in L^2(\Omega_{per})$, its mean value over Ω_{per} is denoted by $\langle f \rangle$, i.e.,

$$\langle f \rangle = (f, 1) = \frac{1}{|\Omega_{per}|} \int_{\Omega_{per}} f(x) dx.$$

We often write $x \in \Omega$ as $x = {}^\top(x', x_n)$, $x' = {}^\top(x_1, \dots, x_{n-1}) \in \mathbb{R}^{n-1}$. The partial derivatives of a function u are denoted by ∂_{x_j} , $\partial_{x_j} \partial_{x_k}$, and so on.

We will work in spaces of functions Q -periodic in x' , and so, we introduce the function spaces $L_{per}^2(\Omega_{per})$, $C_{per}^\infty(\overline{\Omega_{per}})$, $C_{0,per}^\infty(\Omega_{per})$, $H_{per}^l(\Omega_{per})$, $H_{0,per}^l(\Omega_{per})$ that are defined by

$$L_{per}^2(\Omega_{per}) = \{u|_{\Omega_{per}}; u \in L_{loc}^2(\overline{\Omega}), u(x' + \frac{2\pi}{\alpha_j} e'_j, x_n) = u(x', x_n), (x', x_n) \in \Omega, 1 \leq j \leq n-1\},$$

$$C_{per}^\infty(\overline{\Omega_{per}}) = \{u|_{\overline{\Omega_{per}}}; u \in C^\infty(\overline{\Omega}), u(x' + \frac{2\pi}{\alpha_j} e'_j, x_n) = u(x', x_n), (x', x_n) \in \Omega, 1 \leq j \leq n-1\},$$

$$C_{0,per}^\infty(\Omega_{per}) = \{u \in C_{per}^\infty(\overline{\Omega_{per}}); u = 0 \text{ in a neighborhood of } \partial\Omega\},$$

$$H_{per}^l(\Omega_{per}) = \text{the closure of } C_{per}^\infty(\overline{\Omega_{per}}) \text{ in } H^l(\Omega_{per}),$$

$$H_{0,per}^l(\Omega_{per}) = \text{the closure of } C_{0,per}^\infty(\Omega_{per}) \text{ in } H^l(\Omega_{per}).$$

Observe that $L_{per}^2(\Omega_{per})$ can be identified with $L^2(\Omega_{per})$, and that

$$H_{per}^l(\Omega_{per}) = \{u \in H^l(\Omega_{per}); \partial_{x'}^\beta u|_{\Sigma_{j,-}} = \partial_{x'}^\beta u|_{\Sigma_{j,+}}, 1 \leq j \leq n-1, |\beta| \leq l-1\},$$

$$H_{0,per}^1(\Omega_{per}) = \{u \in H_{per}^1(\Omega_{per}); u|_{\Sigma_n} = 0\}.$$

We also set

$$L_{*,per}^2(\Omega_{per}) = \{f \in L_{per}^2(\Omega_{per}); \langle f \rangle = 0\}$$

and

$$H_{*,per}^l(\Omega_{per}) = H_{per}^l(\Omega_{per}) \cap L_{*,per}^2(\Omega_{per}).$$

For $\eta' \in \mathbb{R}^{n-1}$ we denote

$$\tilde{\eta}' = {}^\top(\eta', 0) \in \mathbb{R}^n,$$

and $\nabla_{\eta'}$ is defined by

$$\nabla_{\eta'} = \nabla + i\tilde{\eta}'.$$

We further introduce the following notations

$$\nabla'_{\eta'} = \nabla' + i\eta', \quad \Delta_{\eta'} = \nabla_{\eta'} \cdot \nabla_{\eta'}, \quad \text{div}_{\eta'} w = \nabla_{\eta'} \cdot w - i \langle \tilde{\eta}' \cdot w \rangle.$$

Here ∇' denotes the gradient with respect to $x' = (x_1, \dots, x_{n-1}) \in \mathbb{R}^{n-1}$. We note that $\Delta_{\eta'} = \nabla'_{\eta'} \cdot \nabla'_{\eta'} + \partial_{x_n}^2$.

We next introduce some operators. We denote by \mathbb{P}_0 and $\tilde{\mathbb{P}}$ the following $(n+1) \times (n+1)$ diagonal matrices:

$$\mathbb{P}_0 = \text{diag}(1, 0, \dots, 0), \quad \tilde{\mathbb{P}} = \text{diag}(0, 1, \dots, 1).$$

Note that $\mathbb{P}_0 u = {}^\top(\phi, 0)$ and $\tilde{\mathbb{P}} u = {}^\top(0, w)$ for $u = {}^\top(\phi, w)$ with $w = {}^\top(w_1, \dots, w_n)$.

We denote the kernel and range of an operator A by $\text{Ker } A$ and $R(A)$, respectively.

For a function $f = f(x')$ ($x' \in \mathbb{R}^{n-1}$), we denote its Fourier transform by \hat{f} or $\mathcal{F}[f]$:

$$\hat{f}(\xi') = \mathcal{F}[f](\xi') = \int_{\mathbb{R}^{n-1}} f(x') e^{-i\xi' \cdot x'} dx' \quad (\xi' \in \mathbb{R}^{n-1}).$$

The inverse Fourier transform \mathcal{F}^{-1} is defined by

$$\mathcal{F}^{-1}[f](x') = (2\pi)^{-(n-1)} \int_{\mathbb{R}^{n-1}} f(\xi') e^{i\xi' \cdot x'} d\xi' \quad (x' \in \mathbb{R}^{n-1}).$$

We next introduce the Bloch wave decomposition. Let $\mathcal{S}(\mathbb{R}^{n-1})$ denote the Schwartz space on \mathbb{R}^{n-1} .

Definition 2.1. We define the operator T by $(T\varphi)(x', \eta')$ ($x' \in \mathbb{R}^{n-1}$, $\eta' \in \mathbb{R}^{n-1}$) for $\varphi \in \mathcal{S}(\mathbb{R}^{n-1})$, where

$$\begin{aligned} (T\varphi)(x', \eta') &= \frac{1}{(2\pi)^{\frac{n-1}{2}} |Q|^{\frac{1}{2}}} \sum_{(k_1, \dots, k_{n-1}) \in \mathbb{Z}^{n-1}} \hat{\varphi}(\eta' + \sum_{j=1}^{n-1} k_j \alpha_j \mathbf{e}'_j) e^{i \sum_{j=1}^{n-1} k_j \alpha_j x_j} \\ &= \frac{1}{|Q^*|^{\frac{1}{2}}} \sum_{(l_1, \dots, l_{n-1}) \in \mathbb{Z}^{n-1}} \varphi(x' + \sum_{j=1}^{n-1} l_j \frac{2\pi}{\alpha_j} \mathbf{e}'_j) e^{-i\eta' \cdot (x' + \sum_{j=1}^{n-1} l_j \frac{2\pi}{\alpha_j} \mathbf{e}'_j)}. \end{aligned} \quad (2.1)$$

We also define the operator U as follows. For a function $\varphi(x', \eta') \in C^\infty(\mathbb{R}^{n-1} \times \mathbb{R}^{n-1})$ such that $\varphi(x', \eta')$ is Q -periodic in x' and $\varphi(x', \eta')e^{i\eta' \cdot x'}$ is Q^* -periodic in η' , we define $(U\varphi)(x')$ ($x' \in \mathbb{R}^{n-1}$) by

$$(U\varphi)(x') = \frac{1}{|Q^*|^{\frac{1}{2}}} \int_{Q^*} \varphi(\eta', x') e^{i\eta' \cdot x'} d\eta'. \quad (2.2)$$

Note that $\varphi(x, \eta' + \alpha_j \mathbf{e}'_j) = \varphi(x', \eta') e^{-i\alpha_j \mathbf{e}'_j \cdot x'}$ ($j = 1, \dots, n-1$).

The operators T and U have the following properties. See, e.g., [21] for the details.

- Proposition 2.2.** (i) $(T\varphi)(x', \eta')$ is Q -periodic in x' and $(T\varphi)(x', \eta')e^{i\eta' \cdot x'}$ is Q^* -periodic in η' .
- (ii) T is uniquely extended to an isometric operator from $L^2(\mathbb{R}^{n-1})$ to $L^2(Q^*; L^2(Q))$.
- (iii) U is the inverse operator of T .
- (iv) Let ψ be Q -periodic in x' . Then it holds that $T(\psi\varphi) = \psi T(\varphi)$.
- (v) $T(\partial_{x_j} \varphi) = (\partial_{x_j} + i\eta_j)T\varphi$ ($j = 1, \dots, n-1$) and T defines an isomorphism from $H^l(\mathbb{R}^{n-1})$ to $L^2(Q^*; H_{per}^l(Q))$. (Here $H_{per}^l(Q)$ denotes the space of Q -periodic functions belonging to $H^l(Q)$, as in the case of $H_{per}^l(\Omega_{per})$.)

We next consider T as an operator acting on functions in $H^l(\Omega)$. Let $y = \Phi(x)$ be the following transformation

$$\begin{cases} y' = x', \\ y_n = \frac{1}{\omega_2(x') - \omega_1(x')} (x_n - \omega_1(x')). \end{cases}$$

Then Φ is a diffeomorphism from Ω to $\mathbb{R}^{n-1} \times (0, 1)$ and Φ transforms Q -periodic functions on Ω to those on $\mathbb{R}^{n-1} \times (0, 1)$. We denote the inverse transform of Φ by Ψ and we define the operators Φ^* and Ψ^* by $[\Phi^* u](x) = u(\Phi(x))$ and $[\Psi^* u](y) = u(\Psi(y))$, respectively. Then Φ^* is an isomorphism from $H^l(\Omega)$ to $H^l(\mathbb{R}^{n-1} \times (0, 1))$, and likewise, from $H_{per}^l(\Omega_{per})$ to $H_{per}^l(Q \times (0, 1))$, where $H_{per}^l(Q \times (0, 1))$ denotes the space of Q -periodic functions belonging to $H^l(Q \times (0, 1))$.

It is not difficult to see that Proposition 2.2 holds with $H^l(\mathbb{R}^{n-1})$ replaced by $H^l(\mathbb{R}^{n-1} \times (0, 1))$, and likewise, with $H_{per}^l(Q)$ replaced by $H_{per}^l(Q \times (0, 1))$. It then follows that $\Phi^*T\Psi^*$ is an isomorphism from $H^l(\Omega)$ to $L^2(Q^*; H_{per}^l(\Omega_{per}))$. Using the second expression of T in Definition 2.1 and the periodicity of ω_j ($j = 1, 2$), one can see that $\Phi^*T\Psi^*u = Tu$ for functions u on Ω . Therefore, we will write $\Phi^*T\Psi^*u$ as Tu if no confusion will occur.

3 Main results

In this section we state the main results of this paper.

Let us consider the following linear problem

$$\partial_t u + Lu = 0, \quad u = {}^\top(\phi, w). \quad (3.1)$$

Here L is the operator on $L^2(\Omega)$ given by

$$L = \begin{pmatrix} 0 & \gamma \operatorname{div} \\ \gamma \nabla & -\nu \Delta - \tilde{\nu} \nabla \operatorname{div} \end{pmatrix} \quad (3.2)$$

with domain

$$D(L) = \{u = {}^\top(\phi, w) \in L^2(\Omega); w \in H_0^1(\Omega), Lu \in L^2(\Omega)\}. \quad (3.3)$$

Our main issue is to investigate the spectral properties of the semigroup generated by $-L$. We first state that it generates a contraction semigroup.

Theorem 3.1. *The operator $-L$ generates a contraction C_0 -semigroup e^{-tL} on $L^2(\Omega)$ and it holds that*

$$\|e^{-tL}u_0\|_{L^2(\Omega)} \leq \|u_0\|_{L^2(\Omega)} \quad (u_0 \in L^2(\Omega)).$$

The semigroup e^{-tL} has the following properties.

Theorem 3.2. *There is a bounded projection $\Pi : L^2(\Omega) \rightarrow L^2(\Omega)$ satisfying $\Pi L \subset L\Pi$ and $\Pi e^{-tL} = e^{-tL}\Pi$, and there hold the following estimates uniformly for $t > 0$ and $u_0 \in L^1(\Omega) \cap L^2(\Omega)$:*

- (i) $\|e^{-tL}\Pi u_0\|_{L^2(\Omega)} \leq C(1+t)^{-\frac{n-1}{4}}\|u_0\|_{L^1(\Omega)},$
- (ii) $\|e^{-tL}(I - \Pi)u_0\|_{L^2(\Omega)} \leq Ce^{-\beta t}\|u_0\|_{L^2(\Omega)},$
- (iii) $\|e^{-tL}\Pi u_0 - [e^{-tH}\sigma_0]u^{(0)}\|_{L^2(\Omega)} \leq Ct^{-\frac{n-1}{4}-\frac{1}{2}}\|u_0\|_{L^1(\Omega)}.$

Here β is a positive constant; e^{-tH} is the C_0 -semigroup in $L^2(\mathbb{R}^{n-1})$ generated by the operator $-H$ defined by

$$H\sigma = -\frac{\gamma^2}{\nu} \sum_{i,j=1}^{n-1} a_{ij} \partial_{x_i} \partial_{x_j} \sigma \quad (\sigma \in D(H))$$

with domain $D(H) = H^2(\mathbb{R}^{n-1})$; and σ_0 and $u^{(0)}$ are given as follows:

$$\sigma_0 = \frac{|Q|}{|\Omega_{per}|} \int_{\omega_1(x')}^{\omega_2(x')} \phi_0(x', x_n) dx_n, \quad u^{(0)} = {}^\top(1, 0).$$

Here a_{ij} satisfies that

$$\sum_{i,j=1}^{n-1} a_{ij} \xi_i \xi_j \geq \kappa_0 |\xi'|^2 \quad (\xi' = (\xi_1, \dots, \xi_{n-1}) \in \mathbb{R}^{n-1})$$

with a constant $\kappa_0 > 0$ independent of ξ' . Furthermore, a_{ij} is given as $a_{ij} = (\nabla w^{(i)}, \nabla w^{(j)})_{L^2(\Omega_{per})}$ with ${}^\top(\phi^{(k)}, w^{(k)})$ ($k = 1, \dots, n-1$) satisfying the following Stokes system in Ω_{per} :

$$\begin{cases} \operatorname{div} w^{(k)} = 0, \\ -\Delta w^{(k)} + \nabla \phi^{(k)} = \mathbf{e}'_k, \\ w^{(k)}|_{\Sigma_{j,+}} = w^{(k)}|_{\Sigma_{j,-}}, \quad \phi^{(k)}|_{\Sigma_{j,+}} = \phi^{(k)}|_{\Sigma_{j,-}}, \quad w^{(k)}|_{\Sigma_n} = 0, \quad \langle \phi^{(k)} \rangle = 0 \end{cases} \quad (3.4)$$

for some $\phi^{(k)}$.

The proof of Theorem 3.2 will be given in sections 4 and 5. To prove Theorem 3.2, we will consider the resolvent problem $\lambda u + L_{\eta'} u = f$ on $L^2_{per}(\Omega_{per})$ with parameter $\eta' \in Q^*$. In the case of $|\eta'| \leq r_0$ for some small $r_0 > 0$, we regard $L_{\eta'}$ as a perturbation of L_0 and apply the analytic perturbation theory to study the spectrum of $-L_{\eta'}$. For $\eta' \in Q^*$ with $|\eta'| \geq r_0$, we establish estimates for a modified Stokes system and apply an energy method. Based on the analysis for $-L_{\eta'}$, we give a proof of Theorem 3.2.

4 Spectral properties of $L_{\eta'}$

In this section we investigate spectral properties of $L_{\eta'}$.

4.1 Formulation

Let us consider the resolvent problem for (3.1)

$$(\lambda + L)u = f, \quad u \in D(L). \quad (4.1)$$

Here $\lambda \in \mathbb{C}$ is a resolvent parameter.

Applying Ψ^* to (4.1), we have

$$(\lambda + \Psi^* L) \Psi^* u = \Psi^* f \quad \text{in } \mathbb{R}^{n-1} \times (0, 1). \quad (4.2)$$

Here $\Psi^* L$ is the differential operator of the form

$$\Psi^* L = \begin{pmatrix} 0 & \sum_{j=1}^n l_{12}^j(y', y_n) \partial_{y_j} \\ \sum_{j=1}^n l_{21}^j(y', y_n) \partial_{y_j} & \sum_{j,k=1}^n l_{22}^{j,k}(y', y_n) \partial_{y_j} \partial_{y_k} + \sum_{j=1}^n l_{22}^j(y', y_n) \partial_{y_j} \end{pmatrix}$$

with some l_{pq}^j and l_{pq}^{jk} ($p, q = 1, 2$) being Q -periodic in y' . We next apply T to (4.2). It then follows from Proposition 2.2 (i), (iv) and (v) that (4.2) is transformed into the following problem on $Q \times (0, 1)$:

$$(\lambda + \Psi^* L_{\eta'}) T \Psi^* u = T \Psi^* f \quad (\eta' \in Q^*) \quad (4.3)$$

with Q -periodic boundary condition in y' . Applying Φ^* to (4.3) we arrive at

$$(\lambda + L_{\eta'}) T u = T f \quad \text{on } \Omega_{per} \quad (4.4)$$

with the dual parameter $\eta' \in Q^*$, where $L_{\eta'}$ is the operator on $L_{per}^2(\Omega_{per})$ of the form

$$L_{\eta'} := \begin{pmatrix} 0 & \gamma^\top \nabla_{\eta'} \\ \gamma \nabla_{\eta'} & -\nu \Delta_{\eta'} - \tilde{\nu} \nabla_{\eta'}^\top \nabla_{\eta'} \end{pmatrix}$$

with domain $D(L_{\eta'})$

$$D(L_{\eta'}) = \{u = {}^\top(\phi, w) \in L_{per}^2(\Omega_{per}); L_{\eta'} u \in L_{per}^2(\Omega_{per}), w \in H_{0,per}^1(\Omega_{per})\}.$$

It is not difficult to see that $D(L_{\eta'}) = D(L_0)$ for all $\eta' \in Q^*$ and that $L_{\eta'}$ is a closed operator on $L_{per}^2(\Omega_{per})$.

If $\lambda \in \rho(-L_{\eta'})$, then, by (4.4), u is written as

$$u = U(\lambda + L_{\eta'})^{-1} T f.$$

Therefore, to investigate the resolvent of $-L$, we will consider the problem for $-L_{\eta'}$:

$$\lambda u + L_{\eta'} u = f, \quad u \in D(L_0). \quad (4.5)$$

Before going further, we also introduce the adjoint operator of $L_{\eta'}$. We define the operator $L_{\eta'}^*$ by

$$L_{\eta'}^* := \begin{pmatrix} 0 & -\gamma^\top \nabla_{\eta'} \\ -\gamma \nabla_{\eta'} & -\nu \Delta_{\eta'} - \tilde{\nu} \nabla_{\eta'}^\top \nabla_{\eta'} \end{pmatrix}$$

with domain

$$D(L_{\eta'}^*) = \{u = {}^\top(\phi, w) \in L_{per}^2(\Omega_{per}); L_{\eta'}^* u \in L_{per}^2(\Omega_{per}), w \in H_{0,per}^1(\Omega_{per})\}.$$

One can see that $D(L_{\eta'}^*) = D(L_0^*)$ for all $\eta' \in Q^*$ and that $L_{\eta'}^*$ is the adjoint operator of $L_{\eta'}$.

4.2 The case $|\eta'| \leq r_0$

In this subsection we consider (4.5) with $|\eta'| \leq r_0$ for some sufficiently small $r_0 > 0$. It is convenient to write $L_{\eta'}$ in the form

$$L_{\eta'} := L_0 + \sum_{j=1}^{n-1} \eta_j L_j^{(1)} + \sum_{j,k=1}^{n-1} \eta_j \eta_k L_{jk}^{(2)},$$

where

$$\begin{aligned} L_0 &:= \begin{pmatrix} 0 & \gamma \operatorname{div} \\ \gamma \nabla & -\nu \Delta - \tilde{\nu} \nabla \operatorname{div} \end{pmatrix}, \\ L_j^{(1)} &:= i \begin{pmatrix} 0 & \gamma^\top \mathbf{e}_j \\ \gamma \mathbf{e}_j & -2\nu \partial_{x_j} - \tilde{\nu} \mathbf{e}_j \operatorname{div} - \tilde{\nu} \nabla (\cdot^\top \mathbf{e}_j) \end{pmatrix}, \\ L_{jk}^{(2)} &:= \begin{pmatrix} 0 & 0 \\ 0 & -\nu \delta_{jk} - \tilde{\nu} \mathbf{e}_j^T \mathbf{e}_k \end{pmatrix}. \end{aligned}$$

We also set

$$M_{\eta'} = \sum_{j=1}^{n-1} \eta_j L_j^{(1)} + \sum_{j,k=1}^{n-1} \eta_j \eta_k L_{jk}^{(2)},$$

namely,

$$M_{\eta'} := \begin{pmatrix} 0 & i\gamma^\top \tilde{\eta}' \\ i\gamma \tilde{\eta}' & \nu(|\eta'|^2 - 2i\tilde{\eta}' \cdot \nabla) - i\tilde{\nu} \tilde{\eta}'^\top (\nabla + i\tilde{\eta}') - i\tilde{\nu} \nabla^\top \tilde{\eta}' \end{pmatrix}.$$

Similarly, we write $L_{\eta'}^*$ as

$$L_{\eta'}^* := L_0^* + \sum_{j=1}^{n-1} \eta_j L_j^{(1)*} + \sum_{j,k=1}^{n-1} \eta_j \eta_k L_{jk}^{(2)*},$$

where

$$\begin{aligned} L_0^* &:= \begin{pmatrix} 0 & -\gamma \operatorname{div} \\ -\gamma \nabla & -\nu \Delta - \tilde{\nu} \nabla \operatorname{div} \end{pmatrix}, \\ L_j^{(1)*} &:= i \begin{pmatrix} 0 & -\gamma^\top \mathbf{e}_j \\ -\gamma \mathbf{e}_j & -2\nu \partial_{x_j} - \tilde{\nu} \mathbf{e}_j \operatorname{div} - \tilde{\nu} \nabla (\cdot^\top \mathbf{e}_j) \end{pmatrix}, \\ L_{jk}^{(2)*} &:= \begin{pmatrix} 0 & 0 \\ 0 & -\nu \delta_{jk} - \tilde{\nu} \mathbf{e}_j^T \mathbf{e}_k \end{pmatrix}. \end{aligned}$$

We begin with the resolvent estimates for the case $\eta' = 0$ which implies the generation of a contraction semigroup e^{-tL_0} .

In what follows we write

$$X = L_{per}^2(\Omega_{per})$$

for simplicity of notation.

Proposition 4.1. *It holds that $\{\lambda; \operatorname{Re} \lambda > 0\} \subset \rho(-L_0)$, and if $\operatorname{Re} \lambda > 0$, then*

$$\begin{aligned} \|(\lambda + L_0)^{-1} f\|_2 &\leq \frac{1}{\operatorname{Re} \lambda} \|f\|_2, \\ \|\nabla \tilde{\mathbb{P}}(\lambda + L_0)^{-1} f\|_2 &\leq \frac{1}{(\nu \operatorname{Re} \lambda)^{\frac{1}{2}}} \|f\|_2. \end{aligned}$$

The same conclusion also holds for the adjoint operator L_0^ .*

Proof. Let $\operatorname{Re} \lambda > 0$. Since

$$\operatorname{Re}((\lambda + L_0)u, u) = \nu \|\nabla w\|_2^2 + \tilde{\nu} \|\operatorname{div} w\|_2^2 + \operatorname{Re} \lambda \|u\|_2^2, \quad (4.6)$$

we see that if $(\lambda + L_0)u = 0$, then $u = 0$, and so, $\lambda + L_0$ is injective when $\operatorname{Re} \lambda > 0$. Observe also that if $\operatorname{Re} \lambda > 0$, then

$$\|u\|_2 \leq \frac{1}{\operatorname{Re} \lambda} \|(\lambda + L_0)u\|_2, \quad (4.7)$$

$$\|\nabla w\|_2 \leq \frac{1}{(\nu \operatorname{Re} \lambda)^{\frac{1}{2}}} \|(\lambda + L_0)u\|_2. \quad (4.8)$$

It follows from (4.7) that $R(\lambda + L_0)$ is a closed subspace of X . We note that these inequalities also holds with L_0 replaced by L_0^* . Let $v \in R(\lambda + L_0)^\perp$. Then, since $((\lambda + L_0)u, v) = 0$ for all $u \in D(L_0)$, we see that $v \in D(L_0^*)$ and $(\lambda + L_0^*)v = 0$. This, together with (4.7) with L_0 replaced by L_0^* , implies that $v = 0$. We thus conclude that $R(\lambda + L_0) = X$ and, hence, $\lambda + L_0$ is surjective. This completes the proof. \square

The following estimates show that $-L_{\eta'}$ also generates a contraction semigroup.

Proposition 4.2. *It holds that $\{\lambda : \operatorname{Re} \lambda > 0\} \subset \rho(-L_{\eta'})$ and the following estimates hold for $\operatorname{Re} \lambda > 0$:*

$$\|(\lambda + L_{\eta'})^{-1}f\|_2 \leq \frac{1}{\operatorname{Re} \lambda} \|f\|_2,$$

$$\|\nabla \tilde{\mathbb{P}}(\lambda + L_{\eta'})^{-1}f\|_2 \leq \frac{C}{(\nu)^{\frac{1}{2}}} \left(\frac{1}{(\operatorname{Re} \lambda)^{\frac{1}{2}}} + \frac{1}{\operatorname{Re} \lambda} \right) \|f\|_2.$$

The same conclusion also holds for the adjoint operator $L_{\eta'}^$*

Proof. We have

$$\operatorname{Re}((\lambda + L_{\eta'})u, u) = \operatorname{Re} \lambda \|u\|_2^2 + \nu \|\nabla_{\eta'} w\|_2^2 + \tilde{\nu} \|\nabla_{\eta'} \cdot w\|_2^2.$$

It then follows that if $\operatorname{Re} \lambda > 0$, then

$$\|u\|_2 \leq \frac{1}{\operatorname{Re} \lambda} \|(\lambda + L_{\eta'})u\|_2.$$

We also have

$$\operatorname{Re}((\lambda + L_{\eta'})u, u) \geq \nu \|\nabla w\|_2^2 + (\operatorname{Re} \lambda - C) \|u\|_2^2$$

for a constant $C > 0$ uniformly for $\eta' \in Q^*$. Therefore, we deduce that

$$\nu \|\nabla w\|_2^2 \leq \|(\lambda + L_{\eta'})u\|_2 \|u\|_2 + C \|u\|_2^2 \leq C \left(\frac{1}{\operatorname{Re} \lambda} + \frac{1}{(\operatorname{Re} \lambda)^2} \right) \|(\lambda + L_{\eta'})u\|_2^2,$$

which gives

$$\|\nabla w\|_2 \leq \frac{C}{\nu^{\frac{1}{2}}} \left(\frac{1}{(\operatorname{Re} \lambda)^{\frac{1}{2}}} + \frac{1}{\operatorname{Re} \lambda} \right) \|(\lambda + L_{\eta'})u\|_2.$$

As in the proof of Proposition 4.1, one can now obtain the desired results. This completes the proof. \square

We next show that $\lambda = 0$ is a simple eigenvalue of $-L_0$.

Proposition 4.3. *There exists a constant $\beta_0 > 0$ such that $\rho(-L_0) \supset \{\lambda; \operatorname{Re} \lambda > -\beta_0, \lambda \neq 0\}$. Furthermore, $\lambda = 0$ is a simple eigenvalue of $-L_0$ and it holds that for λ satisfying $\operatorname{Re} \lambda > -\beta_0$ and $\lambda \neq 0$,*

$$(\lambda + L_0)^{-1}f = \frac{1}{\lambda}\Pi^{(0)}f + S_\lambda(I - \Pi^{(0)})f,$$

and there hold the following estimates uniformly for λ satisfying $\operatorname{Re} \lambda > -\beta_0$:

$$\|S_\lambda(I - \Pi^{(0)})f\|_2 \leq \frac{C}{\operatorname{Re} \lambda + \beta_0} \|f\|_2,$$

$$\|\nabla \tilde{\mathbb{P}} S_\lambda(I - \Pi^{(0)})f\|_2 \leq \frac{C}{(\operatorname{Re} \lambda + \beta_0)^{\frac{1}{2}}} \|f\|_2.$$

Here $\Pi^{(0)}$ is the eigenprojection for the eigenvalue $\lambda = 0$ defined by

$$\Pi^{(0)}u = (u, u^{(0)*})u^{(0)} = \langle \phi \rangle u^{(0)}, \quad u = {}^\top(\phi, w),$$

where

$$u^{(0)} = {}^\top(1, 0), \quad u^{(0)*} = \frac{1}{|\Omega_{per}|} {}^\top(1, 0),$$

and S_λ is the operator defined by

$$S_\lambda := [(I - \Pi^{(0)})(\lambda + L_0)(I - \Pi^{(0)})]^{-1}.$$

The same conclusion also holds with $L_0, S_\lambda, \Pi^{(0)}$ replaced by $L_0^*, S_\lambda^*, \Pi^{(0)*}$, respectively where

$$S_\lambda^* := [(I - \Pi^{(0)*})(\lambda + L_0^*)(I - \Pi^{(0)*})]^{-1}, \quad \Pi^{(0)*}u = (u, u^{(0)})u^{(0)*}.$$

We give a proof of Proposition 4.3 only for L_0 since the case of L_0^* can be treated similarly. To prove Proposition 4.3, we prepare the following two lemmas.

Lemma 4.4. *It holds that $\operatorname{Ker} L_0 = \operatorname{span}\{u^{(0)}\}$ and $\Pi^{(0)}$ is a bounded projection on $L^2(\Omega_{per})$ that satisfies $\Pi^{(0)}X = \operatorname{Ker} L_0$, $\Pi^{(0)}L_0 \subset L_0\Pi^{(0)} = 0$.*

Proof of Lemma 4.4. Let $L_0u = 0$. It then follows from (4.6) that $\nabla w = 0$, and hence, $\nabla \phi = 0$. This implies that $w = 0$ and $\phi = \text{constant}$. This shows that $\operatorname{Ker} L_0 = \operatorname{span}\{u^{(0)}\}$. Clearly, $\Pi^{(0)}$ is a bounded projection onto $\operatorname{Ker} L_0$. For $u = {}^\top(\phi, w)$, we have $L_0\Pi^{(0)}u = {}^\top(0, \gamma \nabla \langle \phi \rangle) = 0$. On the other hand, for $u \in D(L_0)$, we have $\Pi^{(0)}L_0u = \langle \gamma \operatorname{div} w \rangle u^{(0)} = 0$. We thus conclude that $\Pi^{(0)}L_0 \subset L\Pi^{(0)} = 0$. This completes the proof. \square

Lemma 4.5. *It holds that $\rho(-L_0|_{(I - \Pi^{(0)})X}) \supset \{\lambda; \operatorname{Re} \lambda > -\beta_0\}$ with a positive constant β_0 , and the estimates for S_λ in Proposition 4.3 hold true.*

Proof of Lemma 4.5. In the proof we set $\mathcal{A} = -L_0|_{(I-\Pi^{(0)})X}$. Let us consider $\lambda u + \mathcal{A}u = f$. It is known that there exists a bounded linear operator $\mathcal{B} : L_*^2(\Omega_{per}) \rightarrow H_{0,per}^1(\Omega_{per})$ such that for any $g \in L_{*,per}^2(\Omega_{per})$ it holds that $\operatorname{div} \mathcal{B}g = g$ and $\|\nabla \mathcal{B}g\|_2 \leq c_0 \|g\|_2$ for some constant $c_0 > 0$. See [1, 2, 4] for the details.

We follow the argument in [6]. We introduce a new inner product

$$((u_1, u_2)) = (u_1, u_2) - \delta\{(w_1, \mathcal{B}\phi_2) + (\mathcal{B}\phi_1, w_2)\}$$

for $u_j = {}^\top(\phi_j, w_j)$ ($j = 1, 2$) with a constant $\delta > 0$ to be determined later. This pairing $((u_1, u_2))$ defines an inner product on $L_{*,per}^2(\Omega_{per}) \times L^2(\Omega_{per})$ if $\delta > 0$ is sufficiently small. In fact, using the Poincaré inequality: $\|w\|_2 \leq c_1 \|\nabla w\|_2$, we see that there exists a constant $C > 0$ such that

$$((u, u)) = \|u\|_2^2 - \delta\{(w, \mathcal{B}\phi) + (\mathcal{B}\phi, w)\} \geq (1 - \delta c_0 c_1) \|u\|_2^2$$

and $((u, u)) \leq (1 + \delta c_0 c_1) \|u\|_2^2$. Therefore, $((\cdot, \cdot))$ is an inner product and the norm defined by $((\cdot, \cdot))$ is equivalent to the norm $\|\cdot\|_2$ if $\delta > 0$ is taken sufficiently small.

We denote $\mathcal{A}u = {}^\top(\mathcal{A}_1 u, \mathcal{A}_2 u)$. Note that $\int_{\Omega_{per}} \mathcal{A}_1 u dx = 0$. We see that

$$\begin{aligned} ((\mathcal{A}u, u)) &= (L_0 u, u) - \delta\{(\mathcal{A}_2 u, \mathcal{B}\phi) + (\mathcal{B}(\mathcal{A}_1 u), w)\} \\ &\geq \nu \|\nabla w\|_2^2 + \tilde{\nu} \|\operatorname{div} w\|_2^2 + \frac{1}{2} \delta \gamma \|\phi\|_2^2 - \delta \left\{ \left(\frac{\nu^2 c_0^2}{\gamma} + \gamma c_1^2 \right) \|\nabla w\|_2^2 + \frac{\tilde{\nu}^2}{\gamma} \|\operatorname{div} w\|_2^2 \right\} \\ &\geq \frac{1}{2} \nu \|\nabla w\|_2^2 + \frac{1}{2} \tilde{\nu} \|\operatorname{div} w\|_2^2 + \frac{1}{2} \delta \gamma \|\phi\|_2^2 \end{aligned}$$

if $\delta > 0$ is taken suitably small. Therefore, we have

$$(1 - \delta c_0 c_1) \operatorname{Re} \lambda \|u\|_2 + \frac{1}{2} \nu \|\nabla w\|_2^2 + \frac{1}{2} \tilde{\nu} \|\operatorname{div} w\|_2^2 + \frac{1}{2} \delta \gamma \|\phi\|_2^2 \leq \operatorname{Re}((f, u)) \leq C \|f\|_2 \|u\|_2.$$

Setting $\beta_0 = \frac{1}{2(1-\delta c_0 c_1)} \min\{\delta \gamma, \frac{\nu}{2c_1^2}\}$ we find by the Poincaré inequality that

$$(\operatorname{Re} \lambda + \beta_0) \|u\|_2 \leq C \|f\|_2.$$

We thus conclude that if $\operatorname{Re} \lambda + \beta_0 > 0$, then

$$\|u\| \leq \frac{C}{\operatorname{Re} \lambda + \beta_0} \|f\|_2$$

and

$$\|\nabla w\|_2 \leq \frac{C}{(\operatorname{Re} \lambda + \beta_0)^{\frac{1}{2}}} \|f\|_2.$$

These estimates, together with Proposition 4.1, yield the desired results. This completes the proof. \square

We are now in a position to prove Proposition 4.3.

Proof of Proposition 4.3. We set $X_0 = \Pi^{(0)}X$ and $X_1 = (I - \Pi^{(0)})X$. By Lemma 4.4, we have $X = X_0 \oplus X_1$ and $\rho(-L_0|_{X_0}) = \{\lambda : \lambda \neq 0\}$. This, together with Lemma 4.5, shows that $\{\lambda ; \operatorname{Re} \lambda > -\beta_0, \lambda \neq 0\} \subset \rho(-L_0)$,

$$(\lambda + L_0)^{-1}f = \frac{1}{\lambda}\Pi^{(0)}f + S_\lambda(I - \Pi^{(0)})f,$$

and S_λ satisfies the desired estimates. This completes the proof. \square

We next derive the resolvent estimates for $-L_{\eta'}$ with $|\eta'| \leq r_0$.

Theorem 4.6. *There exists a constant $r_0 > 0$ such that if $\eta' \in Q^*$ satisfies $|\eta'| \leq r_0$, then*

$$\Sigma_1 := \{\lambda ; \operatorname{Re} \lambda \geq -\frac{3}{4}\beta_0\} \cap \{\lambda ; |\lambda| \geq \frac{\beta_0}{2}\} \subset \rho(-L_{\eta'}),$$

and the following estimates hold uniformly for $\lambda \in \Sigma_1$:

$$\|(\lambda + L_{\eta'})^{-1}f\|_2 \leq \frac{C}{\operatorname{Re} \lambda + \beta_0} \|f\|_2,$$

$$\|\nabla \tilde{\mathbb{P}}(\lambda + L_{\eta'})^{-1}f\|_2 \leq \frac{C}{(\operatorname{Re} \lambda + \beta_0)^{\frac{1}{2}}} \|f\|_2.$$

The same conclusion also holds with $L_{\eta'}$ replaced by $L_{\eta'}^*$.

Proof. Let $\lambda \in \Sigma_1$. By Proposition 4.3, we see that

$$\|(\lambda + L_0)^{-1}f\|_2 + \|\nabla \tilde{\mathbb{P}}(\lambda + L_0)^{-1}f\|_2 \leq C_1 \|f\|_2 \quad (4.9)$$

uniformly for $\lambda \in \Sigma_1$. Here C_1 is a constant depending only on β_0 . It then follows that

$$\|L_j^{(1)}u\|_2 \leq C\{\|w\|_2 + \|\nabla w\|_2 + \|\phi\|_2\} \leq CC_1\|(\lambda + L_0)u\|_2, \quad (4.10)$$

$$\|L_{jk}^{(2)}u\|_2 \leq C\|w\|_2 \leq CC_1\|(\lambda + L_0)u\|_2 \quad (4.11)$$

uniformly for $\lambda \in \Sigma_1$ and $u \in D(L_0)$. We thus obtain

$$\|M_{\eta'}(\lambda + L_0)^{-1}f\|_2 \leq CC_1|\eta'| \|f\|_2 \quad (\lambda \in \Sigma_1)$$

uniformly for $\lambda \in \Sigma_1$ and $f \in X$. Therefore, if $r_0 > 0$ is a constant satisfying $r_0 < \frac{1}{CC_1}$, then $\lambda \in \rho(-L_{\eta'})$ for $|\eta'| \leq r_0$ and it holds that

$$(\lambda + L_{\eta'})^{-1} = (\lambda + L_0)^{-1} \sum_{N=0}^{\infty} (-1)^N (M_{\eta'}(\lambda + L_0)^{-1})^N,$$

$$\|(\lambda + L_{\eta'})^{-1}f\|_2 \leq \frac{C}{\operatorname{Re} \lambda + \beta_0} \sum_{N=0}^{\infty} \|M_{\eta'}(\lambda + L_0)^{-1}\|^N \|f\|_2 \leq \frac{C}{\operatorname{Re} \lambda + \beta_0} \|f\|_2.$$

Similarly, $\|\nabla \tilde{\mathbb{P}}(\lambda + L_{\eta'})^{-1}f\|_2$ can be estimated as

$$\|\nabla \tilde{\mathbb{P}}(\lambda + L_{\eta'})^{-1}f\|_2 \leq \frac{C}{(\operatorname{Re} \lambda + \beta_0)^{\frac{1}{2}}} \|f\|_2.$$

The case of $L_{\eta'}^*$ can be proved similarly. This completes the proof. \square

We now show that $\sigma(-L_{\eta'}) \cap \{\lambda; |\lambda| < \frac{\beta_0}{2}\}$ consists of a simple eigenvalue whose real part is negative and of order $O(|\eta'|^2)$ as $\eta' \rightarrow 0$.

Theorem 4.7. *There exists a constant $r_0 > 0$ such that if $|\eta'| \leq r_0$, then $\sigma(-L_{\eta'}) \cap \{\lambda; |\lambda| < \frac{\beta_0}{2}\} = \{\lambda_{\eta'}\}$. Here $\lambda_{\eta'}$ is a simple eigenvalue that satisfies*

$$\lambda_{\eta'} = -\frac{\gamma^2}{\nu} \kappa(\eta') + O(|\eta'|^3) \quad (\eta' \rightarrow 0),$$

where

$$\kappa(\eta') = \sum_{j,k=1}^{n-1} a_{jk} \eta_j \eta_k, \quad a_{jk} = \frac{1}{|\Omega_{per}|} (\nabla w_1^{(j)}, \nabla w_1^{(k)}).$$

Here $w_1^{(k)}$ ($k = 1, \dots, n-1$) satisfy the Stokes system (3.4) for some $\phi_1^{(k)}$; and $\kappa(\eta')$ satisfies $\kappa(\eta') \geq \kappa_0 |\eta'|^2$ with some constant $\kappa_0 > 0$. As a result, it holds that $\text{Re } \lambda_{\eta'} \leq -\frac{\kappa_0}{2} \frac{\gamma^2}{\nu} |\eta'|^2$.

Remark 4.8. A similar result holds for $L_{\eta'}^*$ with simple eigenvalue $\lambda_{\eta'}^* = \bar{\lambda}_{\eta'}$.

Remark 4.9. Since $\lambda_{\eta'} \rightarrow 0$ as $\eta' \rightarrow 0$, we see that for any $\beta \in (0, \frac{\beta_0}{2})$, there exists a constant $r = r(\beta) > 0$ such that if $|\eta'| \leq r(\beta)$, then $|\lambda_{\eta'}| < \beta$ and $\{\lambda; \text{Re } \lambda \geq -\frac{3}{4}\beta_0\} \cap \{\lambda; |\lambda| \geq \beta\} \subset \rho(-L_{\eta'})$.

Proof of Theorem 4.7. In view of Proposition 4.3, (4.10) and (4.11), we can apply the analytic perturbation theory to see that $\sigma(-L_{\eta'}) \cap \{\lambda; |\lambda| < \frac{\beta_0}{2}\}$ consists of a simple eigenvalue, say $\lambda_{\eta'}$, for sufficiently small η' , and that $\lambda_{\eta'}$ is expanded as

$$\lambda_{\eta'} = \sum_{j=0}^{n-1} \lambda_j^{(1)} \eta_j + \sum_{j,k=0}^{n-1} \lambda_{jk}^{(2)} \eta_j \eta_k + O(|\eta'|^3)$$

with

$$\lambda_j^{(1)} = -(L_j^{(1)} u^{(0)}, u^{(0)*}),$$

$$\lambda_{jk}^{(2)} = -\frac{1}{2} ((L_{jk}^{(2)} + L_{kj}^{(2)}) u^{(0)}, u^{(0)*}) + \frac{1}{2} ((L_j^{(1)} S L_k^{(1)} + L_k^{(1)} S L_j^{(1)}) u^{(0)}, u^{(0)*}).$$

Here $S = S_\lambda|_{\lambda=0}$. See, e.g., [12, Chap. VII], [21, Chap. XII].

Let us compute $\lambda_j^{(1)}$. Since $(u, u^{(0)*}) = \langle \phi \rangle$ for $u = {}^\top(\phi, w)$ and $L_j^{(1)} u^{(0)} = {}^\top(0, i\gamma e_j)$, we have $\lambda_j^{(1)} = 0$.

As for $\lambda_{jk}^{(2)}$, since $L_{jk}^{(2)} u^{(0)} = 0$, we have $\langle L_{jk}^{(2)} u^{(0)} \rangle = 0$. Furthermore,

$$\frac{1}{2} ((L_j^{(1)} S L_k^{(1)} + L_j^{(1)} S L_k^{(1)}) u^{(0)}, u^{(0)*}) = (L_j^{(1)} S L_k^{(1)} u^{(0)}, u^{(0)*}) = \langle L_j^{(1)} S L_k^{(1)} u^{(0)} \rangle.$$

We compute $\langle L_j^{(1)} S L_k^{(1)} u^{(0)} \rangle$. Set $u_1 = {}^\top(\phi_1, w_1) = S L_k^{(1)} u_0$. Then u_1 is a solution of

$$L_0 u_1 = (I - \Pi^{(0)}) L_k^{(1)} u^{(0)} = L_k^{(1)} u^{(0)}, \quad \langle \phi_1 \rangle = 0,$$

that is,

$$\begin{cases} \gamma \operatorname{div} w_1 = 0, \\ -\nu \Delta w_1 + \gamma \nabla \phi_1 = i\gamma \mathbf{e}_k, \\ w_1|_{\Sigma_{j,+}} = w_1|_{\Sigma_{j,-}}, \quad \phi_1|_{\Sigma_{j,+}} = \phi_1|_{\Sigma_{j,-}}, \quad w_1|_{\Sigma_n} = 0, \quad \langle \phi_1 \rangle = 0. \end{cases}$$

Lemma 4.5 implies that for each $k = 1, \dots, n-1$, there exists a unique solution $\tilde{u}_1^{(k)} = {}^\top(\tilde{\phi}_1^{(k)}, \tilde{w}_1^{(k)})$ of this system. Let $u_1^{(k)} = {}^\top(\phi_1^{(k)}, w_1^{(k)})$ be a unique solution of (3.4) Then $\tilde{\phi}_1^{(k)} = i\phi_1^{(k)}$ and $\tilde{w}_1^{(k)} = \frac{i\gamma}{\nu} w_1^{(k)}$, and hence,

$$\begin{aligned} L_j^{(1)} S L_k^{(1)} u^{(0)} &= i \begin{pmatrix} 0 & \gamma {}^\top \mathbf{e}_j \\ \gamma \mathbf{e}_j & -2\nu \partial_{x_j} - \tilde{\nu} \mathbf{e}_j \operatorname{div} - \tilde{\nu} \nabla ({}^\top \mathbf{e}_j) \end{pmatrix} \begin{pmatrix} \tilde{\phi}_1^{(k)} \\ \tilde{w}_1^{(k)} \end{pmatrix} \\ &= -\frac{\gamma^2}{\nu} \begin{pmatrix} \mathbf{e}_j \cdot w_1^{(k)} \\ * \end{pmatrix}. \end{aligned}$$

It then follows that

$$\begin{aligned} \lambda_{jk}^{(2)} &= \langle L_j^{(1)} S L_k^{(1)} u^{(0)} \rangle = -\frac{\gamma^2}{\nu} \langle \mathbf{e}_j \cdot w_1^{(k)} \rangle \\ &= -\frac{\gamma^2}{\nu} \langle (-\Delta w_1^{(j)} + \nabla \phi_1^{(j)}) \cdot w_1^{(k)} \rangle = -\frac{\gamma^2}{\nu} \frac{1}{|\Omega_{per}|} (\nabla w_1^{(j)}, \nabla w_1^{(k)}). \end{aligned}$$

Let us show that the matrix $((\nabla w_1^{(j)}, \nabla w_1^{(k)}))$ is positive definite. We first observe that $w_1^{(1)}, \dots, w_1^{(n-1)}$ are linearly independent. In fact, suppose that $w_1 = \sum_{j=1}^{n-1} c_j w_1^{(j)} = 0$. Then $\phi_1 = \sum_{j=1}^{n-1} c_j \phi_1^{(j)}$ satisfies $\nabla \phi_1 = \sum_{j=1}^{n-1} c_j \mathbf{e}_j$. Therefore, ϕ_1 is written as $\phi_1 = c + \sum_{j=1}^{n-1} c_j x_j$ with some constant c . Since ϕ_1 is Q -periodic in $x' = (x_1, \dots, x_{n-1})$ and $\langle \phi_1 \rangle = 0$, we see that $c = c_1 = \dots = c_{n-1} = 0$. We thus conclude that $w_1^{(1)}, \dots, w_1^{(n-1)}$ are linearly independent.

Set $V = \operatorname{span}\{w_1^{(1)}, \dots, w_1^{(n-1)}\}$ and take an orthonormal basis $\{f_1, \dots, f_{n-1}\}$ of V as a subspace of $H_{0,per}^1(\Omega_{per})$ with respect to the inner product $(w, v)_{H_{0,per}^1} = (\nabla w, \nabla v)$. Then $w_1^{(m)}$ is written as $w_1^{(m)} = \sum_{k=1}^{n-1} b_{mk} f_k$ for $m = 1, \dots, n-1$, and thus, $(w_1^{(1)}, \dots, w_1^{(n-1)}) = (f_1, \dots, f_{n-1})B$, where $B = (\mathbf{b}_1, \dots, \mathbf{b}_{n-1})$ with $\mathbf{b}_m = {}^\top(b_{m1}, \dots, b_{mn-1})$. It then follows that $\mathbf{b}_1, \dots, \mathbf{b}_{n-1}$ are linearly independent. We have $(\nabla w_1^{(m)}, \nabla w_1^{(l)}) = (w_1^{(m)}, w_1^{(l)})_{H_{0,per}^1} = (BB^*)_{ml}$. Since BB^* is positive definite, so is the matrix $((\nabla w_1^{(m)}, \nabla w_1^{(l)}))$. It then follows that there is a constant $\kappa_0 > 0$ such that

$$\sum_{j,k=1}^{n-1} \lambda_{jk}^{(2)} \eta_j \eta_k = -\sum_{j,k=1}^{n-1} \frac{\gamma^2}{\nu} \frac{1}{|\Omega_{per}|} (\nabla w_1^{(j)}, \nabla w_1^{(k)}) \eta_j \eta_k = -\frac{\gamma^2}{\nu} |B^* \eta'|^2 \leq -\kappa_0 \frac{\gamma^2}{\nu} |\eta'|^2$$

for all $\eta' \in \mathbb{R}^{n-1}$. Therefore, there exists $r_0 > 0$ such that if $|\eta'| \leq r_0$, then $\operatorname{Re} \lambda_{\eta'} \leq -\frac{\kappa_0}{2} \frac{\gamma^2}{\nu} |\eta'|^2$. This completes the proof. \square

Let $\Pi_{\eta'}$ be the eigenprojection for the eigenvalue $\lambda_{\eta'}$. Since $\Pi_{\eta'} e_{\eta'}^{-tL} = e^{\lambda_{\eta'} t} \Pi_{\eta'}$, we have the following estimate.

Theorem 4.10. *If $|\eta'| \leq r_0$, then it holds that*

$$\|e^{-tL_{\eta'}} u_0 - e^{\lambda_{\eta'} t} \Pi_{\eta'} u_0\|_2 \leq C e^{-\frac{\beta_0}{2} t} \|u_0\|_2.$$

Theorem 4.10 follows from Theorems 4.6 and 4.7. See, e.g., [3, Chap. V, Theorem 1.11], [23].

We close this subsection with the estimates for the eigenprojections $\Pi_{\eta'}$ and $\Pi_{\eta'}^*$ for the eigenvalues $\lambda_{\eta'}$ and $\lambda_{\eta'}^* (= \bar{\lambda}_{\eta'})$ of $-L_{\eta'}$ and $-L_{\eta'}^*$, respectively.

Theorem 4.11. *For any nonnegative integer k , there exists a constant $r_k > 0$ such that the following estimates hold uniformly for $|\eta'| \leq r_k$:*

- (i) $\|\Pi_{\eta'} u\|_{H^k} \leq C \|u\|_1.$
- (ii) $\|(\Pi_{\eta'} - \Pi^{(0)}) u\|_{H^k} \leq C |\eta'| \|u\|_1.$

The same conclusion holds with $\Pi_{\eta'}$ replaced by $\Pi_{\eta'}^$.*

Proof. By Theorem 4.6 we have

$$\Pi_{\eta'} = \frac{1}{2\pi i} \int_{|\lambda|=\frac{\beta_0}{2}} (\lambda + L_{\eta'})^{-1} d\lambda, \quad \Pi_{\eta'}^* = \frac{1}{2\pi i} \int_{|\lambda|=\frac{\beta_0}{2}} (\lambda + L_{\eta'}^*)^{-1} d\lambda.$$

Furthermore, $u_{\eta'} = \Pi_{\eta'} u^{(0)}$ and $u_{\eta'}^* = \Pi_{\eta'}^* u^{(0)*}$ are eigenfunctions of $-L_{\eta'}$ and $-L_{\eta'}^*$ for the eigenvalues $\lambda_{\eta'}$ and $\lambda_{\eta'}^* = \bar{\lambda}_{\eta'}$, respectively; and it holds that

$$\Pi_{\eta'} u = \frac{(u, u_{\eta'}^*)}{(u_{\eta'}, u_{\eta'}^*)} u_{\eta'}.$$

Note that $u_{\eta'}|_{\eta'=0} = \Pi^{(0)} u^{(0)} = u^{(0)}$ and $u_{\eta'}^*|_{\eta'=0} = \Pi^{(0)*} u^{(0)*} = u^{(0)*}$.

In view of (4.9), (4.10) and (4.11), we see that $(\lambda + L_{\eta'})^{-1}$ is expanded as

$$(\lambda + L_{\eta'})^{-1} = (\lambda + L_0)^{-1} - (\lambda + L_0)^{-1} \sum_{j=1}^{n-1} \eta_j L_j^{(1)} (\lambda + L_0)^{-1} + R_{\eta'}(\lambda)$$

and $R_{\eta'}(\lambda)$ is estimated as

$$\|R_{\eta'}(\lambda) f\|_2 \leq C |\eta'|^2 \|f\|_2, \quad \|\nabla \tilde{\mathbb{P}} R_{\eta'}(\lambda) f\|_2 \leq C |\eta'|^2 \|f\|_2$$

uniformly for $|\eta'| \leq r_0$ and $|\lambda| = \frac{\beta_0}{2}$. We write $u_{\eta'}$ as

$$\begin{aligned} u_{\eta'} &= u^{(0)} + \frac{1}{2\pi i} \int_{|\lambda|=\frac{\beta_0}{2}} \left(- \sum_{j=1}^{n-1} \eta_j (\lambda + L_0)^{-1} L_j^{(1)} (\lambda + L_0)^{-1} u^{(0)} \right) d\lambda \\ &\quad + \frac{1}{2\pi i} \int_{|\lambda|=\frac{\beta_0}{2}} R_{\eta'}(\lambda) u^{(0)} d\lambda \\ &=: u^{(0)} + \sum_{j=1}^{n-1} \eta_j u_j^{(1)} + u^{(2)}. \end{aligned}$$

Using (4.9), (4.10) and (4.11), we have

$$\|u_j^{(1)}\|_2 + \|\nabla \tilde{\mathbb{P}} u_j^{(1)}\|_2 \leq C, \quad \|u^{(2)}\|_2 + \|\nabla \tilde{\mathbb{P}} u^{(2)}\|_2 \leq C|\eta'|^2.$$

Similarly, we have the expression for $u_{\eta'}^*$:

$$u_{\eta'}^* = u^{(0)*} + \sum_{j=1}^{n-1} \eta_j u_j^{(1)*} + u^{(2)*},$$

with estimates

$$\|u_j^{(1)*}\|_2 + \|\nabla \tilde{\mathbb{P}} u_j^{(1)*}\|_2 \leq C, \quad \|u^{(2)*}\|_2 + \|\nabla \tilde{\mathbb{P}} u^{(2)*}\|_2 \leq C|\eta'|^2.$$

It then follows that

$$(u_{\eta'}, u_{\eta'}^*) = (u^{(0)}, u^{(0)*}) + (u_{\eta'} - u^{(0)}, u_{\eta'}^*) + (u^{(0)}, u_{\eta'}^* - u^{(0)*}) \geq 1 - C|\eta'| \geq \frac{1}{2}$$

for $|\eta'| \leq r_0$ with $r_0 > 0$ replaced by a smaller one if necessary.

If we could have the estimates $\|u_{\eta'}^*\|_\infty \leq C$ and $\|\partial_x^\alpha u_{\eta'}\|_2 \leq C$, then it would hold that $\|\partial_x^\alpha \Pi_{\eta'} u\|_2 \leq C\|u\|_1 \|u_{\eta'}^*\|_\infty \|\partial_x^\alpha u_{\eta'}\|_2 \leq C\|u\|_1$. So we will deduce the estimates for $u_{\eta'}$ and $u_{\eta'}^*$, in other words, for $(\lambda + L_{\eta'})^{-1}u^{(0)}$ and $(\lambda + L_{\eta'}^*)^{-1}u^{(0)}$.

In the remaining we consider only $(\lambda + L_{\eta'})^{-1}u^{(0)}$ since $(\lambda + L_{\eta'}^*)^{-1}u^{(0)}$ can be estimated similarly. We also observe that the integral path of $u_{\eta'} = \frac{1}{2\pi i} \int_{|\lambda|=\frac{\beta_0}{2}} (\lambda + L_{\eta'})^{-1}u^{(0)} d\lambda$ can be deformed into $\{|\lambda| = \beta\} \subset \rho(-L_{\eta'})$.

We claim the following

Proposition 4.12. *Let m be a nonnegative integer. Then there exist constants $r_m > 0$ and $\beta_m > 0$ such that if $|\eta'| \leq r_m$ and $\frac{\beta_m}{2} \leq |\lambda| \leq \beta_m$, then it holds that $(\lambda + L_{\eta'})^{-1}u^{(0)} \in H_{per}^{m+1}(\Omega_{per}) \times (H_{per}^{m+2} \cap H_{0,per}^1)(\Omega_{per})$ and*

$$\|(\lambda + L_{\eta'})^{-1}u^{(0)}\|_{H^{m+1} \times H^{m+2}} \leq C$$

uniformly for $|\eta'| \leq r_m$ and $\frac{\beta_m}{2} \leq |\lambda| \leq \beta_m$.

To prove Proposition 4.12, we employ the following lemma.

Lemma 4.13. *Let m be a nonnegative integer. Then there exists $\tilde{\beta}_m > 0$ such that if $|\lambda| \leq \tilde{\beta}_m$, then it holds that $S_\lambda f \in H_{*,per}^{m+1}(\Omega_{per}) \times (H_{per}^{m+2} \cap H_{0,per}^1)(\Omega_{per})$ for any $f \in H_{*,per}^{m+1}(\Omega_{per}) \times H_{per}^m(\Omega_{per})$ and $S_\lambda f$ satisfies the estimates*

$$\|S_\lambda f\|_{H^{m+1} \times H^{m+2}} \leq C\|f\|_{H^{m+1} \times H^m}.$$

uniformly for λ with $|\lambda| \leq \tilde{\beta}_m$.

The proof of Lemma 4.13 will be given later.

Proof of Proposition 4.12. We prove by induction on m . We set $u := (\lambda + L_{\eta'})^{-1}u^{(0)}$. By Theorems 4.6 and 4.7, we have $\|u\|_{L^2 \times H^1} \leq C$ uniformly for $|\eta'| \leq r_0$ and $\frac{\beta_0}{4} \leq |\lambda| \leq \frac{\beta_0}{2}$ with r_0 replaced by a smaller one if necessary.

We write $(\lambda + L_0)u = u^{(0)} - M_{\eta'}u$ and decompose $u = {}^\top(\phi, w)$ as $u = \langle \phi \rangle u^{(0)} + u_1$, where $\Pi^{(0)}u = \langle \phi \rangle u^{(0)}$ and $u_1 = (I - \Pi^{(0)})u$. Then we have

$$\langle \phi \rangle = \frac{1}{\lambda} \{1 - i\gamma \langle \eta' \cdot w' \rangle\},$$

$$(\lambda + L_0)u_1 = -(M_{\eta'}u_1 + \langle \phi \rangle M_{\eta'}u^{(0)} - \Pi^{(0)}M_{\eta'}u).$$

It then follows that

$$|\langle \phi \rangle| = \frac{1}{|\lambda|} \{1 + \gamma r_0 \|w\|_2\} \leq C \quad (4.12)$$

uniformly for $|\eta'| \leq r_0$ and $\frac{\beta_1}{2} \leq |\lambda| \leq \beta_1$ with $\beta_1 > 0$ to be determined later. On the other hand, we have

$$\|M_{\eta'}u_1 + \langle \phi \rangle M_{\eta'}u^{(0)} - \Pi^{(0)}M_{\eta'}u\|_{H^1 \times L^2} \leq C\|u\|_{L^2 \times H^1} \leq C$$

uniformly for $|\eta'| \leq r_0$ and $\frac{\beta_0}{4} \leq |\lambda| \leq \frac{\beta_0}{2}$. It then follows from Remark 4.9 and Lemma 4.13 that, with a suitable choice of $r_1 > 0$ and $\beta_1 > 0$, the estimate $\|u_1\|_{H^1 \times H^2} \leq C$ holds uniformly for $|\eta'| \leq r_1$ and $\frac{\beta_1}{2} \leq |\lambda| \leq \beta_1$. This, together with (4.12), proves Proposition 4.12 for $m = 0$.

Assume that the proposition holds for $m = k$. We will show that the proposition holds for $m = k + 1$. By the inductive assumption, we have

$$\|M_{\eta'}u_1 + \langle \phi \rangle M_{\eta'}u^{(0)} - \Pi^{(0)}M_{\eta'}u\|_{H^{k+2} \times H^{k+1}} \leq C\|u\|_{H^{k+1} \times H^{k+2}} \leq C$$

uniformly for $|\eta'| \leq r_k$ and $\frac{\beta_k}{2} \leq |\lambda| \leq \beta_k$. It then follows from Remark 4.9 and Lemma 4.13 that the estimate $\|u_1\|_{H^{k+2} \times H^{k+3}} \leq C$ holds uniformly for $|\eta'| \leq r_{k+1}$ and $\frac{\beta_{k+1}}{2} \leq |\lambda| \leq \beta_{k+1}$. Combining this with (4.12), we conclude that the proposition holds for $m = k + 1$. This completes the proof. \square

We continue the proof of Theorem 4.11. Let m be a nonnegative integer. By Proposition 4.12, we see that

$$(\lambda + L_{\eta'})^{-1}u^{(0)} \in H_{per}^{m+1}(\Omega_{per}) \times (H_{per}^{m+2} \cap H_{0,per}^1)(\Omega_{per}),$$

$$\|(\lambda + L_{\eta'})^{-1}u^{(0)}\|_{H^{m+1} \times H^{m+2}} \leq C$$

uniformly for $|\eta'| \leq r_m$ and $|\lambda| = \beta_m$. Deforming the integral path into $\{|\lambda| = \beta_m\}$, we thus deduce that $u_{\eta'} \in H_{per}^{m+1}(\Omega_{per}) \times H_{per}^{m+2}(\Omega_{per})$ and

$$\|u_{\eta'}\|_{H^{m+1} \times H^{m+2}} = \left\| \frac{1}{2\pi i} \int_{|\lambda|=\frac{\beta_m}{2}} (\lambda + L_{\eta'})^{-1}u^{(0)} d\lambda \right\|_{H^{m+1} \times H^{m+2}} \leq C. \quad (4.13)$$

Taking $m = k - 1$, we have $\|\partial_x^\alpha u_{\eta'}\|_{L^2} \leq C$ for $|\alpha| \leq k$ and $|\eta'| \leq r_k$. Similarly we can obtain (4.13) with $u_{\eta'}$ replaced by $u_{\eta'}^*$, and hence, $\|u_{\eta'}^*\|_\infty \leq C\|u_{\eta'}^*\|_{H^{\lfloor \frac{n}{2} \rfloor + 1}} \leq C$. It then follows that

$$\|\partial_x^\alpha \Pi_{\eta'} u\|_2 \leq C\|u_{\eta'}^*\|_\infty \|\partial_x^\alpha u_{\eta'}\|_2 \|u\|_1 \leq C\|u\|_1.$$

This proves (i).

Let us next consider (ii). We write $\Pi_{\eta'} u - \Pi^{(0)} u$ as

$$\begin{aligned}\Pi_{\eta'} u - \Pi^{(0)} u &= \left(\frac{1}{(u_{\eta'}, u_{\eta'}^*)} - 1 \right) (u, u^{(0)*}) u^{(0)} + \frac{1}{(u_{\eta'}, u_{\eta'}^*)} \{ (u, u_{\eta'}^*) u_{\eta'} - (u, u^{(0)*}) u^{(0)} \} \\ &=: I_1 + I_2.\end{aligned}$$

As for I_1 , we have

$$\begin{aligned}|(u_{\eta'}, u_{\eta'}^*) - 1| &= |(u_{\eta'} - u^{(0)}, u_{\eta'}^*) + (u^{(0)}, u_{\eta'}^* - u^{(0)*})| \\ &\leq C \{ \|u_{\eta'} - u^{(0)}\|_2 + \|u_{\eta'}^* - u^{(0)*}\|_2 \}.\end{aligned}$$

Since

$$u_{\eta'} - u^{(0)} = \frac{1}{2\pi i} \int_{|\lambda|=\beta_m} (\lambda + L_0)^{-1} \sum_{N=1}^{\infty} (-1)^N [M_{\eta'}(\lambda + L_0)^{-1}]^N u^{(0)} d\lambda,$$

we have $\|u_{\eta'} - u^{(0)}\|_2 \leq C|\eta'|$, and likewise, $\|u_{\eta'}^* - u^{(0)*}\|_2 \leq C|\eta'|$. We thus obtain

$$\|\partial_x^\alpha I_1\|_2 \leq C|\eta'| |(u, u^{(0)*}) \partial_x^\alpha u^{(0)}| \leq C|\eta'| \|u\|_1 \|u^{(0)*}\|_\infty \|\partial_x^\alpha u^{(0)}\|_\infty \leq C|\eta'| \|u\|_1.$$

As for I_2 , we have

$$\begin{aligned}&\|\partial_x^\alpha \{ (u, u_{\eta'}^*) u_{\eta'} - (u, u^{(0)*}) u^{(0)} \}\|_2 \\ &= \|(u, u_{\eta'}^* - u^{(0)*}) \partial_x^\alpha u_{\eta'} + (u, u^{(0)*}) \partial_x^\alpha (u_{\eta'} - u^{(0)})\|_2 \\ &\leq \|u\|_1 \|u_{\eta'}^* - u^{(0)*}\|_\infty \|\partial_x^\alpha u_{\eta'}\|_2 + \|u\|_1 \|u^{(0)*}\|_\infty \|\partial_x^\alpha (u_{\eta'} - u^{(0)})\|_2 \\ &\leq C \|u\|_1 \{ \|u_{\eta'}^* - u^{(0)*}\|_{H^{\lfloor \frac{n}{2} \rfloor + 1}} \|u_{\eta'}\|_{H^k} + \|u_{\eta'} - u^{(0)}\|_{H^k} \}.\end{aligned}$$

Since $\|M_{\eta'} u\|_{H^k \times H^{k-1}} \leq C|\eta'| \|u\|_{H^{k-1} \times H^k}$, with the aid of Lemma 4.13, we see that

$$\|(\lambda + L_0)^{-1} [M_{\eta'}(\lambda + L_0)^{-1}]^N u^{(0)}\|_{H^k \times H^k} \leq (C|\eta'|)^N$$

uniformly for $|\eta'| \leq r_k$ and $|\lambda| = \beta_k$. Taking $r_k > 0$ smaller if necessary, we obtain $\|u_{\eta'} - u^{(0)}\|_{H^k} \leq C|\eta'|$. Similarly, we can obtain $\|u_{\eta'}^* - u^{(0)*}\|_{H^{\lfloor \frac{n}{2} \rfloor + 1}} \leq C|\eta'|$. It then follows that $\|\partial_x^\alpha I_2\|_2 \leq C|\eta'| \|u\|_1$ for $|\alpha| \leq k$. We thus conclude that

$$\|\partial_x^\alpha (\Pi_{\eta'} - \Pi^{(0)}) u\|_2 \leq C|\eta'| \|u\|_1$$

for $|\alpha| \leq k$. This completes the proof. \square

In the remaining of this subsection we give a proof of Lemma 4.13.

Proof of Lemma 4.13. We set $\tilde{u} = S_\lambda f$. Then

$$L_0 \tilde{u} = f - \lambda \tilde{u},$$

which is regarded as an inhomogeneous Stokes system. This can be solved for \tilde{u} if $|\lambda|$ is suitably small. In fact, let $f \in H_{*,per}^{m+1}(\Omega_{per}) \times H_{per}^m(\Omega_{per})$. Then, for each

$\check{v} \in H_{*,per}^{m+1}(\Omega_{per}) \times (H_{per}^{m+2} \cap H_{0,per}^1)(\Omega_{per})$, there uniquely exists $\check{u} \in H_{*,per}^{m+1}(\Omega_{per}) \times (H_{per}^{m+2} \cap H_{0,per}^1)(\Omega_{per})$ such that $L_0 \check{u} = f - \lambda \check{v}$ and

$$\|\check{u}\|_{H^{m+1} \times H^{m+2}} \leq C|\lambda| \|\check{v}\|_{H^{m+1} \times H^{m+2}} + C\|f\|_{H^{m+1} \times H^m}.$$

See, e.g., [22, Chap. III, Theorem 1.5.3]. This estimate shows that the map $\check{v} \mapsto \check{u}$ is a contraction on $H_{*,per}^{m+1}(\Omega_{per}) \times (H_{per}^{m+2} \cap H_{0,per}^1)(\Omega_{per})$ when $|\lambda| \leq \tilde{\beta}_m$ with suitably small $\tilde{\beta}_m$. This completes the proof. \square

4.3 The case $|\eta'| \geq r_0$

In this subsection we investigate the spectrum of $-L_{\eta'}$ for $\eta' \in Q^*$ satisfying $|\eta'| \geq r_0$. We have already shown in Proposition 4.2 that $-L_{\eta'}$ generates a contraction semi-group $e^{-tL_{\eta'}}$. We will show that $e^{-tL_{\eta'}}$ has an exponential decay estimate uniformly for $\eta' \in Q^*$ with $|\eta'| \geq r_0$.

We first introduce an inner product of $H_{0,per}^1(\Omega_{per})$ in terms of $\nabla_{\eta'}$.

Proposition 4.14. *Let $\eta' \in Q^*$. Then $(\nabla_{\eta'} w, \nabla_{\eta'} v)$ defines an inner product of $H_{0,per}^1(\Omega_{per})$. Furthermore, $\|\nabla_{\eta'} w\|_2$ is equivalent to $\|w\|_{H^1}$ for $w \in H_{0,per}^1(\Omega_{per})$ and there holds the estimate*

$$C^{-1}\|w\|_{H^1} \leq \|\nabla_{\eta'} w\|_2 \leq C\|w\|_{H^1}$$

uniformly for $\eta' \in Q^*$ and $w \in H_{0,per}^1(\Omega_{per})$.

Proof. It suffices to show that $\|\nabla_{\eta'} w\|_2 = (\nabla_{\eta'} w, \nabla_{\eta'} w)$ is equivalent to $\|w\|_{H^1}$ for $w \in H_{0,per}^1(\Omega_{per})$. Let $w \in H_{0,per}^1(\Omega_{per})$. Then by using the Poincaré inequality, we see that

$$\|w\|_{H^1} \leq C\|\nabla w\|_2 \leq C'(\|\nabla'_{\eta'} w\|_2^2 + \|\partial_{x_n} w\|_2^2)^{\frac{1}{2}} = C'\|\nabla_{\eta'} w\|_2 \leq C''\|w\|_{H^1}.$$

This completes the proof. \square

Before going further, we introduce some notations. We define $D_{\eta'}(w)$ by

$$\begin{aligned} D_{\eta'}(w) &:= \nu \|\nabla_{\eta'} w\|_2^2 + \tilde{\nu} \|\nabla_{\eta'} \cdot w\|_2^2 \\ &= \nu \|\nabla'_{\eta'} w\|_2^2 + \nu \|\partial_{x_n} w\|_2^2 + \tilde{\nu} \|\nabla'_{\eta'} \cdot w'\|_2^2 + \tilde{\nu} \|\partial_{x_n} w^n\|_2^2. \end{aligned}$$

In what follows we denote the projection $I - \Pi^{(0)}$ by Π_1 :

$$\Pi_1 := I - \Pi^{(0)}.$$

To study problem (4.5) for $\eta' \in Q^*$ with $|\eta'| \geq r_0$, we decompose u into its $\Pi^{(0)}$ -part and Π_1 -part in X , namely,

$$u = \sigma u^{(0)} + u_1, \tag{4.14}$$

where $\sigma = (u, u^{(0)*}) = \langle \phi \rangle \in \mathbb{C}$, $u^{(0)} = {}^\top(1, 0)$ and $u_1 = {}^\top(\phi_1, w_1) \in X_1$. We note that

$$\langle \phi_1 \rangle = 0. \tag{4.15}$$

It is easy to see that problem (4.5) is reduced to the following system

$$\begin{cases} \lambda\sigma + i\gamma \langle \tilde{\eta}' \cdot w_1 \rangle = \langle f^0 \rangle, \\ \lambda u_1 + L_{\eta'} u_1 - \Pi^{(0)} M_{\eta'} u_1 + M_{\eta'}(\sigma u^{(0)}) = f_1, \end{cases} \quad (4.16)$$

where $\sigma \in \mathbb{C}$, $u_1 = {}^\top(\phi_1, w_1) \in D(L_{\eta'}) \cap X_1$ and $f_1 = \Pi_1 f := {}^\top(f_1^0, \tilde{f}_1) \in X_1$. We observe that

$$\Pi^{(0)} M_{\eta'} u_1 = i\gamma \langle \tilde{\eta}' \cdot w_1 \rangle u^{(0)}, \quad M_{\eta'}(\sigma u^{(0)}) = {}^\top(0, i\gamma \tilde{\eta}' \sigma).$$

We begin with the following

Proposition 4.15. *It holds that*

$$\operatorname{Re} \lambda(|\sigma|^2 + |u_1|^2) + D_{\eta'}(w_1) = \operatorname{Re}\{\langle f^0 \rangle \bar{\sigma} + (f_1, u_1)\}. \quad (4.17)$$

Proof. Multiplying the first equation of (4.16) by $\bar{\sigma}$, we have

$$\lambda|\sigma|^2 + i\gamma \langle \tilde{\eta}' \cdot w_1 \rangle \bar{\sigma} = \langle f^0 \rangle \bar{\sigma}.$$

Taking the inner product of the second equation of (4.16) with u_1 , we have

$$\lambda \|u_1\|_2^2 + (L_{\eta'} u_1, u_1) + (M_{\eta'}(\sigma u^{(0)}), u_1) - (\Pi^{(0)} M_{\eta'} u_1, u_1) = (f_1, u_1).$$

We add these two equations to obtain

$$\begin{aligned} & \lambda(|\sigma|^2 + \|u_1\|_2^2) + (L_{\eta'} u_1, u_1) + i\gamma \langle \tilde{\eta}' \cdot w_1 \rangle \bar{\sigma} \\ & + (M_{\eta'}(\sigma u^{(0)}), u_1) - (\Pi^{(0)} M_{\eta'} u_1, u_1) = \langle f^0 \rangle \bar{\sigma} + (f_1, u_1). \end{aligned}$$

Since $\operatorname{Re}(L_{\eta'} u_1, u_1) = D_{\eta'}(w_1)$,

$$\operatorname{Re}\{i\gamma \langle \tilde{\eta}' \cdot w_1 \rangle \bar{\sigma} + (M_{\eta'}(\sigma u^{(0)}), u_1)\} = \operatorname{Re}\{2i\operatorname{Im}(i\gamma \langle \tilde{\eta}' \cdot w_1 \rangle \bar{\sigma})\} = 0$$

and

$$(\Pi^{(0)} M_{\eta'} u_1, u_1) = (i\gamma \langle \tilde{\eta}' \cdot w_1 \rangle, \phi_1) = i\gamma \langle \tilde{\eta}' \cdot w_1 \rangle \langle \bar{\phi}_1 \rangle = 0,$$

we obtain

$$\operatorname{Re} \lambda(|\sigma|^2 + |u_1|^2) + D_{\eta'}(w_1) = \operatorname{Re}\{\langle f^0 \rangle \bar{\sigma} + (f, u_1)\}.$$

This completes the proof. \square

For later use, we next derive the estimate for λw_1 .

Proposition 4.16. *It holds that*

$$\operatorname{Re} \bar{\lambda} D_{\eta'}(w_1) + |\lambda|^2 \|w_1\|_2^2 \leq C\{\|f_1\|_2^2 + \|w_1\|_2^2 + |\langle f^0 \rangle|^2 + \|\nabla_{\eta'} w_1\|_2^2\}. \quad (4.18)$$

Proof. We write the second equation of (4.16) as

$$\lambda\phi_1 + \gamma \operatorname{div}_{\eta'} w_1 = f_1^0, \quad (4.19)$$

$$\lambda w_1 - \nu \Delta_{\eta'} w_1 - \tilde{\nu} \nabla_{\eta'} (\nabla_{\eta'} \cdot w_1) + \gamma \nabla_{\eta'} \phi_1 + i\gamma \sigma \tilde{\eta}' = \tilde{f}_1. \quad (4.20)$$

We take the inner product of (4.20) with λw_1 . Then the real part of the resulting equation yields

$$|\lambda|^2 \|w_1\|_2^2 + \operatorname{Re} \bar{\lambda} D_{\eta'}(w_1) = \operatorname{Re} \{ \gamma \bar{\lambda} (\phi_1, \nabla_{\eta'} \cdot w_1) - i \bar{\lambda} \gamma \sigma \langle \tilde{\eta}' \cdot \bar{w}_1 \rangle + \bar{\lambda} (\tilde{f}_1, w_1) \}.$$

The equation (4.19) gives that $\phi = \frac{1}{\lambda} f_1^0 - \frac{\gamma}{\lambda} \operatorname{div}_{\eta'} w_1$, and hence,

$$\begin{aligned} \operatorname{Re} |\lambda|^2 \|w_1\|_2^2 + \operatorname{Re} \bar{\lambda} D_{\eta'}(w_1) &= \operatorname{Re} \left\{ \frac{\gamma \bar{\lambda}}{\lambda} (f_1^0, \nabla_{\eta'} \cdot w_1) - \frac{\bar{\lambda}}{\lambda} \gamma^2 (\operatorname{div}_{\eta'} w_1, \nabla_{\eta'} \cdot w_1) \right. \\ &\quad \left. - i \bar{\lambda} \gamma \sigma \langle \tilde{\eta}' \cdot \bar{w}_1 \rangle + \bar{\lambda} (\tilde{f}_1, w_1) \right\}. \end{aligned} \quad (4.21)$$

By the first equation of (4.16), we have $\sigma = \frac{1}{\lambda} \langle f^0 \rangle - \frac{i\gamma}{\lambda} \langle \tilde{\eta}', w_1 \rangle$. Therefore, the right-hand side of (4.21) is estimated as

$$\begin{aligned} &\text{R.H.S. of (4.21)} \\ &= \operatorname{Re} \left\{ \frac{\gamma \bar{\lambda}}{\lambda} (f_1^0, \nabla_{\eta'} \cdot w_1) - \frac{\bar{\lambda}}{\lambda} \gamma^2 (\operatorname{div}_{\eta'} w_1, \nabla_{\eta'} \cdot w_1) - i \gamma \frac{\bar{\lambda}}{\lambda} \langle f^0 \rangle \langle \tilde{\eta}' \cdot \bar{w}_1 \rangle \right. \\ &\quad \left. - \frac{\bar{\lambda}}{\lambda} \gamma^2 |\langle \tilde{\eta}' \cdot w_1 \rangle|^2 + \bar{\lambda} (\tilde{f}_1, w_1) \right\} \\ &\leq \epsilon |\lambda|^2 \|w_1\|_2^2 + C \{ \|\nabla_{\eta'} w_1\|_2^2 + |\langle f^0 \rangle|^2 + \|f_1^0\|_2^2 + \frac{1}{\epsilon} \|\tilde{f}_1\|_2^2 \} \end{aligned}$$

for any $\epsilon > 0$, where C is a positive constant independent of ϵ . Taking ϵ suitably small, we see that

$$|\lambda|^2 \|w_1\|_2^2 + \operatorname{Re} \bar{\lambda} D_{\eta'}(w_1) \leq C_1 \{ \|f_1\|_2^2 + |\langle f^0 \rangle|^2 + \|\nabla_{\eta'} w_1\|_2^2 \}.$$

This completes the proof. \square

We next derive the coercive estimate for σ .

Proposition 4.17. *It holds that*

$$\operatorname{Re} \lambda |\sigma|^2 + \frac{c_2 \gamma^2}{2\nu} |\eta'|^2 |\sigma|^2 \leq C \left\{ \left(1 + \frac{1}{|\eta'|^2}\right) |\langle f^0 \rangle|^2 + \|\tilde{f}_1\|_2^2 + |\lambda|^2 \|w_1\|_2^2 + D_{\eta'}(w_1) \right\}, \quad (4.22)$$

where c_2 is a positive constant independent of γ , ν and $\eta' \in Q^*$.

To prove Proposition 4.17, we prepare several lemmas.

Lemma 4.18. *Let $f^0 \in L_{*,per}^2(\Omega_{per})$ and let $\tilde{f} \in H_{per}^{-1}(\Omega_{per})$. Then there uniquely exists ${}^\top(\phi, w) \in L_{*,per}^2(\Omega_{per}) \times H_{0,per}^1(\Omega_{per})$ satisfying*

$$\begin{cases} \operatorname{div}_{\eta'} w = f^0, \\ -\Delta_{\eta'} w + \nabla_{\eta'} \phi = \tilde{f}, \\ \phi|_{\Sigma_{j,+}} = \phi|_{\Sigma_{j,-}}, \quad w|_{\Sigma_{j,+}} = w|_{\Sigma_{j,-}}, \quad w|_{\Sigma_n} = 0. \end{cases} \quad (4.23)$$

Furthermore, it holds that

$$\|\phi\|_2 + \|\nabla_{\eta'} w\|_2 \leq C\{\|f^0\|_2 + \|\tilde{f}\|_{H_{per}^{-1}(\Omega_{per})}\}.$$

Lemma 4.18 can be proved in a similar manner to the proof of [22, Chap. III, Theorem 1.4.1]. An outline of the proof of Lemma 4.18 will be given in section 6.

Setting $f^0 = 0$ and $\tilde{f} = \mathbf{e}_k$ in Lemma 4.18, we have the following

Lemma 4.19. *Let ${}^\top(\phi_{1,k,\eta'}^{(1)}, w_{1,k,\eta'}^{(1)}) \in L_{*,per}^2(\Omega_{per}) \times H_{0,per}^1(\Omega_{per})$ be the pair of functions satisfying*

$$\begin{cases} \operatorname{div}_{\eta'} w_{1,k,\eta'}^{(1)} = 0, \\ -\Delta_{\eta'} w_{1,k,\eta'}^{(1)} + \nabla_{\eta'} \phi_{1,k,\eta'}^{(1)} = \mathbf{e}_k, \\ \phi_{1,k,\eta'}|_{\Sigma_{j,+}} = \phi_{1,k,\eta'}|_{\Sigma_{j,-}}, \quad w_{1,k,\eta'}^{(1)}|_{\Sigma_{j,+}} = w_{1,k,\eta'}^{(1)}|_{\Sigma_{j,-}}, \quad w_{1,k,\eta'}|_{\Sigma_n} = 0. \end{cases} \quad (4.24)$$

Then there exists a constant $C > 0$ such that the following inequalities

$$\|w_{1,k,\eta'}^{(1)}\|_{H^1} + \|\phi_{1,k,\eta'}^{(1)}\|_2 \leq C \quad (k = 1, \dots, n-1)$$

hold uniformly for $\eta' \in Q^*$ with $|\eta'| \geq r_0$.

Lemma 4.20. *For each $k = 1, \dots, n-1$, let ${}^\top(\phi_{1,k,\eta'}^{(1)}, w_{1,k,\eta'}^{(1)}) \in L_{*,per}^2(\Omega_{per}) \times H_{0,per}^1(\Omega_{per})$ be the pair of functions satisfying (4.24). Then $w_{1,1,\eta'}^{(1)}, \dots, w_{1,n-1,\eta'}^{(1)}$ are linearly independent.*

Proof. Let $w := c_1 w_{1,1,\eta'}^{(1)} + \dots + c_{n-1} w_{1,n-1,\eta'}^{(1)} = 0$. It then follows from (4.24) that $\nabla_{\eta'} \tilde{\phi} = \sum_{j=1}^{n-1} c_j \mathbf{e}_j$. Here $\tilde{\phi} = c_1 \phi_{1,1,\eta'}^{(1)} + \dots + c_{n-1} \phi_{1,n-1,\eta'}^{(1)}$. Since $|\eta'| \geq r_0$, there exists j such that $\eta_j \neq 0$. For this η_j , since $(\partial_{x_j} + i\eta_j) \tilde{\phi} = c_j$, we have $\partial_{x_j} (e^{i\eta_j x_j} \tilde{\phi}) = c_j e^{i\eta_j x_j}$. This implies that there exists a function $a(\tilde{x}_j)$ ($\tilde{x}_j = (x_1, \dots, x_{j-1}, x_{j+1}, \dots, x_{n-1})$) such that

$$e^{i\eta_j x_j} \tilde{\phi} = a(\tilde{x}_j) + \frac{c_j}{i\eta_j} e^{i\eta_j x_j},$$

namely,

$$\tilde{\phi} = a(\tilde{x}_j) e^{-i\eta_j x_j} + \frac{c_j}{i\eta_j}.$$

Since $\tilde{\phi}$ is Q -periodic in x' , we see that $a(\tilde{x}_j)$ is Q -periodic and $a(\tilde{x}_j) e^{-i\frac{\pi}{\alpha_j} \eta_j} = a(\tilde{x}_j) e^{i\frac{\pi}{\alpha_j} \eta_j}$. Since $0 < |\eta_j| \leq \frac{\alpha_j}{2}$, we have $a(\tilde{x}_j) = 0$, and hence, $\tilde{\phi} = \frac{c_j}{i\eta_j}$. But, since $\langle \tilde{\phi} \rangle = 0$, we have $c_j = 0$, and so $\tilde{\phi} = 0$. This implies that $\sum_{j=1}^{n-1} c_j \mathbf{e}_j = 0$. We thus conclude that $c_j = 0$ ($j = 1, \dots, n-1$). This completes the proof. \square

We are now in a position to prove Proposition 4.17.

Proof of Proposition 4.17. We multiply the first equation of (4.16) by $\bar{\sigma}$ and take the real part of the resulting equation to obtain

$$\operatorname{Re} \lambda |\sigma|^2 + \operatorname{Re} \{i\gamma \langle \tilde{\eta}' \cdot w_1 \rangle \bar{\sigma}\} = \operatorname{Re} \{ \langle f^0 \rangle \bar{\sigma} \}. \quad (4.25)$$

Let us estimate $\text{Re}\{i\gamma < \tilde{\eta}' \cdot w_1 > \bar{\sigma}\}$ on the left-hand side of (4.25). To do so, we decompose w_1 in the following way. In (4.19) and (4.20) we decompose ϕ_1 and w_1 as follows:

$$\phi_1 = \phi_1^{(1)} + \phi_1^{(2)}, \quad w_1 = w_1^{(1)} + w_1^{(2)},$$

where ${}^\top(\phi_1^{(1)}, w_1^{(1)})$ and ${}^\top(\phi_1^{(2)}, w_1^{(2)})$ satisfy the following systems, respectively:

$$\begin{cases} \text{div}_{\eta'} w_1^{(1)} = 0, \\ -\Delta_{\eta'} w_1^{(1)} + \frac{\gamma}{\nu} \nabla_{\eta'} \phi_1^{(1)} = -\frac{i\gamma\sigma}{\nu} \tilde{\eta}', \\ \phi_1^{(1)}|_{\Sigma_{j,+}} = \phi_1^{(1)}|_{\Sigma_{j,-}}, \quad w_1^{(1)}|_{\Sigma_{j,+}} = w_1^{(1)}|_{\Sigma_{j,-}}, \quad w_1^{(1)}|_{\Sigma_n} = 0, \\ \langle \phi_1^{(1)} \rangle = 0 \end{cases} \quad (4.26)$$

and

$$\begin{cases} \gamma \text{div}_{\eta'} w_1^{(2)} = f_1^0 - \lambda \phi_1, \\ -\nu \Delta_{\eta'} w_1^{(2)} + \gamma \nabla_{\eta'} \phi_1^{(2)} = \tilde{f}_1 - \lambda w_1 + \tilde{\nu} \nabla_{\eta'} (\nabla_{\eta'} \cdot w_1), \\ \phi_1^{(2)}|_{\Sigma_{j,+}} = \phi_1^{(2)}|_{\Sigma_{j,-}}, \quad w_1^{(2)}|_{\Sigma_{j,+}} = w_1^{(2)}|_{\Sigma_{j,-}}, \quad w_1^{(1)}|_{\Sigma_n} = 0, \\ \langle \phi_1^{(2)} \rangle = 0. \end{cases} \quad (4.27)$$

Let us estimate $i < \tilde{\eta}' \cdot w_1^{(1)} >$. We see from (4.26) that ${}^\top(\phi_1^{(1)}, w_1^{(1)})$ is written as

$$\begin{pmatrix} \phi_1^{(1)} \\ w_1^{(1)} \end{pmatrix} = -\frac{i\gamma}{\nu} \sigma \sum_{k=1}^{n-1} \eta_k \begin{pmatrix} \frac{\nu}{\gamma} \phi_{1,k,\eta'}^{(1)} \\ w_{1,k,\eta'}^{(1)} \end{pmatrix}.$$

Since $\langle \phi_{1,k,\eta'}^{(1)} \rangle = 0$, we see that $0 = i < \tilde{\eta}' \cdot w_{1,k,\eta'}^{(1)} > \overline{\langle \phi_{1,k,\eta'}^{(1)} \rangle} = (i < \tilde{\eta}' \cdot w_{1,k,\eta'}^{(1)} >, \phi_{1,k,\eta'}^{(1)})$, which implies that $(w_{1,k,\eta'}^{(1)}, \nabla_{\eta'} \phi_{1,k,\eta'}^{(1)}) = -(\text{div}_{\eta'} w_{1,k,\eta'}^{(1)}, \phi_{1,k,\eta'}^{(1)}) = 0$. Taking this into account, we have

$$\begin{aligned} i < \tilde{\eta}' \cdot w_1^{(1)} > &= \frac{\gamma}{\nu} \sigma \sum_{j,k=1}^{n-1} \eta_j \eta_k (w_{1,j,\eta'}^{(1)}, \mathbf{e}_k) = \frac{\gamma}{\nu} \sigma \sum_{j,k=1}^{n-1} \eta_j \eta_k (w_{1,j,\eta'}^{(1)}, -\Delta_{\eta'} w_{1,k,\eta'}^{(1)} + \nabla_{\eta'} \phi_{1,k,\eta'}^{(1)}) \\ &= \frac{\gamma}{\nu} \sigma \sum_{j,k=1}^{n-1} \eta_j \eta_k (\nabla_{\eta'} w_{1,k,\eta'}^{(1)}, \nabla_{\eta'} w_{1,j,\eta'}^{(1)}). \end{aligned}$$

Let $\{f_{1,\eta'}, \dots, f_{n-1,\eta'}\}$ be an orthonormal basis of $\text{span}\{w_{1,k,\eta'}^{(1)}\}_{k=1}^{n-1}$ in $H_{0,per}^1(\Omega_{per})$. Then $w_{1,k,\eta'}^{(1)}$ is written as $w_{1,k,\eta'}^{(1)} = \sum_{m=1}^{n-1} b_{km,\eta'} f_{m,\eta'}$, and therefore,

$$(\nabla_{\eta'} w_{1,j,\eta'}^{(1)}, \nabla_{\eta'} w_{1,k,\eta'}^{(1)}) = \sum_{l,m} b_{jm,\eta'} \overline{b_{kl,\eta'}} (\nabla_{\eta'} f_{m,\eta'}, \nabla_{\eta'} f_{l,\eta'}) = (B_{\eta'} B_{\eta'}^*)_{jk}.$$

Here $B_{\eta'}$ is the $(n-1) \times (n-1)$ matrix given by $B_{\eta'} = (b_{jk,\eta'})_{j,k=1}^{n-1}$. Note that $B_{\eta'}$ is nonsingular since $\{w_{1,k,\eta'}^{(1)}\}_{k=1}^{n-1}$ are linearly independent by Lemma 4.20. We thus find that $B_{\eta'} B_{\eta'}^*$ is positive definite for each η' , and

$$i < \tilde{\eta}' \cdot w_1^{(1)} > = \frac{\gamma}{\nu} \sigma |B_{\eta'}^* \eta'|^2, \quad (4.28)$$

$$|B_{\eta'}^* \eta'|^2 \geq c_2 |\eta'|^2 \quad (4.29)$$

uniformly for $\eta' \in Q^*$ with $|\eta'| \geq r_0$. Here c_2 is the number given by

$$c_2 = \inf_{\eta' \in Q^*, |\eta'| \geq r_0} c_{2, \eta'}$$

with

$$c_{2, \eta'} = \min\{\lambda; \lambda \text{ is an eigenvalue of } B_{\eta'} B_{\eta'}^*\} > 0.$$

Let us show that $c_2 > 0$. To do so, we first show that, for each $j, k = 1, \dots, n-1$, $b_{jk, \eta'} = (\nabla_{\eta'} w_{1,j, \eta'}^{(1)}, \nabla_{\eta'} w_{1,k, \eta'}^{(1)})$ is continuous in η' . If this could be shown, then, by the continuity of the eigenvalue with respect to the components of matrix, we would have $c_2 > 0$.

Set $u_{1,k, \eta'}^{(1)} = {}^\top(\phi_{1,k, \eta'}^{(1)}, w_{1,k, \eta'}^{(1)})$. Then $u_{1,k, \eta'}^{(1)}$ satisfies $u_{1,k, \eta'}^{(1)} \in D(L_0) \cap X_1$ and $(L_0 + \Pi_1 M_{\eta'}) u_{1,k, \eta'}^{(1)} = f_{1,k}$ with $f_{1,k} = {}^\top(0, \mathbf{e}_k)$. By Lemma 4.18, we see that $L_0 + \Pi_1 M_{\eta'}$ has a bounded inverse $(L_0 + \Pi_1 M_{\eta'})^{-1}$ on X_1 and it holds that

$$\|(L_0 + \Pi_1 M_{\eta'})^{-1} f\|_{L^2 \times H^1} \leq C \|f\|_2 \quad (4.30)$$

uniformly for $\eta' \in Q^*$ and $f \in X_1$. On the other hand, we see from (4.10) and (4.11) that

$$\|\Pi_1(M_{\eta'+h'} - M_{\eta'})u\|_2 \leq C|h'| \|(L_0 + \Pi_1 M_{\eta'})u\|_2$$

for $u \in D(L_0) \cap X_1$ and $h' \in \mathbb{R}^{n-1}$ with $|h'| \leq 1$. This, together with (4.30), implies that for each fixed $f \in X_1$, $(L_0 + \Pi_1 M_{\eta'})^{-1} f$ is analytic in $\eta' \in Q^*$ in $L_{*,per}^2(\Omega_{per}) \times H_{0,per}^1(\Omega_{per})$. Since $u_{1,k, \eta'}^{(1)} = (L_0 + \Pi_1 M_{\eta'})^{-1} f_{1,k}$, we find that $w_{1,k, \eta'}^{(1)}$ is analytic in $\eta' \in Q^*$ in $H_{0,per}^1(\Omega_{per})$. We thus see that $b_{jk, \eta'} = (\nabla_{\eta'} w_{1,j, \eta'}^{(1)}, \nabla_{\eta'} w_{1,k, \eta'}^{(1)})$ is continuous in η' , and hence, the eigenvalues of $B_{\eta'} B_{\eta'}^*$ are continuous in η' . Since $c_{2, \eta'}$ is positive for each η' and is continuous in η' , we deduce that

$$c_2 = \inf_{\eta' \in Q^*, |\eta'| \geq r_0} c_{2, \eta'} > 0.$$

By (4.28) and (4.29), we have

$$\operatorname{Re}\{i\gamma < \tilde{\eta}' \cdot w_1^{(1)} > \bar{\sigma}\} = \operatorname{Re}\left\{\frac{\gamma^2}{\nu} |B_{\eta'}^* \eta'|^2 \sigma \bar{\sigma}\right\} \geq c_2 \frac{\gamma^2}{\nu} |\eta'|^2 |\sigma|^2.$$

As for $\operatorname{Re}\{i\gamma < \tilde{\eta}' \cdot w_1^{(2)} > \bar{\sigma}\}$, by Proposition 4.14, we have

$$\operatorname{Re}\{i\gamma < \tilde{\eta}' \cdot w_1^{(2)} > \bar{\sigma}\} \leq \epsilon \frac{\gamma^2}{\nu} |\eta'|^2 |\sigma|^2 + \frac{C\nu}{\epsilon} \|\nabla_{\eta'} w_1^{(2)}\|_2^2$$

for all $\epsilon > 0$ with $C > 0$ independent of ϵ . On the other hand, using Lemma 4.18, we see from (4.27) that

$$\begin{aligned} \|\nabla_{\eta'} w_1^{(2)}\|_2 &\leq C\{\|\operatorname{div}_{\eta'} w_1\|_2 + \|\tilde{f}_1 - \lambda w_1 + \tilde{\nu} \nabla_{\eta'} (\nabla_{\eta'} \cdot w_1)\|_{H_{per}^{-1}(\Omega_{per})}\} \\ &\leq C\{D_{\eta'}(w_1) + \|\lambda w_1\|_2 + \|\tilde{f}_1\|_2\}, \end{aligned}$$

and hence,

$$\operatorname{Re}\{i\gamma < \tilde{\eta}' \cdot w_1 > \bar{\sigma}\} \leq \epsilon \frac{\gamma^2}{\epsilon} |\sigma|^2 + \frac{C}{\epsilon} \{D_{\eta'}(w_1) + \|\lambda w_1\|_2 + \|\tilde{f}_1\|_2\}.$$

Taking $\epsilon = \frac{1}{4}c_2$, we arrive at

$$\operatorname{Re} |\sigma|^2 + \frac{3}{4} \frac{\gamma^2 c_2}{\nu} |\sigma|^2 |\eta'|^2 \leq C_3 \{ | < f^0 > | \|\sigma\| + \|\tilde{f}_1\|_2^2 + |\lambda|^2 \|w_1\|_2^2 + D_{\eta'}(w_1) \},$$

which yields

$$\operatorname{Re} \lambda |\sigma|^2 + \frac{c_0 \gamma^2}{2\nu} |\eta'|^2 |\sigma|^2 \leq C_2 \{ (1 + \frac{1}{|\eta'|^2}) | < f^0 > |^2 + \|\tilde{f}_1\|_2^2 + |\lambda|^2 \|w_1\|_2^2 + D_{\eta'}(w_1) \}.$$

This completes the proof. \square

We now establish the resolvent estimate for $-L_{\eta'}$ with $|\eta'| \geq r_0$.

Theorem 4.21. *Let $\eta' \in Q^*$ satisfy $|\eta'| \geq r_0$. Then there exists a constant $\beta_1 > 0$ such that $\{\lambda; \operatorname{Re} \lambda > -\beta_1\} \subset \rho(-L_{\eta'})$ and if $\operatorname{Re} \lambda > -\beta_1$, then*

$$\|(\lambda + L_{\eta'})^{-1} f\|_2 + \|\nabla \tilde{\mathbb{P}}(\lambda + L_{\eta'})^{-1} f\|_2 \leq \frac{C}{(\operatorname{Re} \lambda + \beta_1)^{\frac{1}{2}}} \|f\|_2.$$

The same conclusion holds with $L_{\eta'}$ replaced by $L_{\eta'}^$.*

Proof. Set $E[u] = (1 + b_2)|\sigma|^2 + \|u_1\|_2^2 + b_1 D_{\eta'}(w_1)$ with constants $b_1, b_2 > 0$ to be determined later. It suffices to show that

$$E[u] \leq \frac{C}{\operatorname{Re} \lambda + \beta_1} \{ | < f^0 > |^2 + \|f_1^0\|_2^2 + \|\tilde{f}_1\|_2^2 \}.$$

Consider (4.17) + (4.18) $\times b_1$. Then taking $b_1 > 0$ suitably small, we have

$$\begin{aligned} \operatorname{Re} \lambda (|\sigma|^2 + \|u_1\|_2^2 + b_1 D_{\eta'}(w_1)) + \frac{1}{4} D_{\eta'}(w_1) + b_1 |\lambda|^2 \|w_1\|_2^2 \\ \leq C \{ | < f^0 > | \|\sigma\| + |(f_1^0, \phi_1)| + \|f_1\|_2^2 + | < f^0 > |^2 \}. \end{aligned} \quad (4.31)$$

We next consider (4.22) $\times b_2$ + (4.31). Then with a suitably small $b_2 > 0$, we have

$$\begin{aligned} \operatorname{Re} \lambda E[u] + \frac{1}{4} D_{\eta'}(w_1) + \frac{b_1}{2} |\lambda|^2 \|w_1\|_2^2 + \frac{b_2}{4} \frac{c_0 \gamma^2}{\nu} |\eta'|^2 |\sigma|^2 \\ \leq C \{ (1 + \frac{1}{|\eta'|^2}) | < f^0 > |^2 + |(f_1^0, \phi_1)| + \|f_1\|_2^2 \}. \end{aligned} \quad (4.32)$$

Since ${}^\top(\phi_1, w_1)$ satisfies ${}^\top(\phi_1, w_1) \in L_{*,per}^2(\Omega_{per}) \times H_{0,per}^1(\Omega_{per})$ and

$$\begin{cases} -\Delta_{\eta'} w_1 + \nabla_{\eta'}(\frac{\gamma}{\nu} \phi_1) = \frac{1}{\nu} \tilde{f}_1 - \frac{1}{\nu} \{ \lambda w_1 - \tilde{\nu} \nabla_{\eta'}(\nabla_{\eta'} \cdot w_1) + i\gamma \sigma \tilde{\eta}' \}, \\ \operatorname{div}_{\eta'} w_1 = \frac{1}{\gamma} \{ f_1^0 - \lambda \phi_1 \}, \end{cases}$$

we see from Lemma 4.18 that

$$\begin{aligned}
\|\phi_1\|_2^2 &\leq C \frac{\nu^2}{\gamma^2} \{ \|\operatorname{div}_{\eta'} w_1\|_2^2 + \frac{1}{\nu^2} \|\tilde{f}_1\|_2^2 + \frac{1}{\nu^2} |\lambda|^2 \|w_1\|_2^2 \\
&\quad + \frac{\tilde{\nu}^2}{\nu^2} \|\nabla_{\eta'}(\nabla_{\eta'} \cdot w_1)\|_{H_{per}^{-1}(\Omega_{per})}^2 + \frac{\gamma^2}{\nu^2} |\eta'|^2 |\sigma|^2 \} \\
&\leq C \{ \frac{(\nu+\tilde{\nu})^2}{\gamma^2} \|\nabla_{\eta'} w_1\|_2^2 + \frac{1}{\gamma^2} \|\tilde{f}_1\|_2^2 + \frac{1}{\gamma^2} |\lambda|^2 \|w_1\|_2^2 + |\eta'|^2 |\sigma|^2 \}.
\end{aligned} \tag{4.33}$$

We consider (4.33) $\times b_3$ + (4.32). Taking $b_3 > 0$ suitably small, we have

$$\begin{aligned}
&\operatorname{Re} \lambda E[u] + \frac{1}{8} D_{\eta'}(w_1) + \frac{b_1}{4} |\lambda| \|w_1\|_2^2 + \frac{b_3}{2} \|\phi_1\|_2^2 + \frac{b_2 c_2}{8} \frac{\gamma^2}{\nu} |\eta'|^2 |\sigma|^2 \\
&\leq C \{ (1 + \frac{1}{|\eta'|^2}) | \langle f^0 \rangle |^2 + \|f_1^0\|_2^2 + \|\tilde{f}_1\|_2^2 + |(f_1^0, \phi_1)| \} \\
&\leq \frac{b_3}{4} \|\phi_1\|_2^2 + C \{ (1 + \frac{1}{|\eta'|^2}) | \langle f^0 \rangle |^2 + \|f_1^0\|_2^2 + \|\tilde{f}_1\|_2^2 \},
\end{aligned}$$

and hence,

$$\begin{aligned}
&\operatorname{Re} \lambda E[u] + \frac{1}{8} D_{\eta'}(w_1) + \frac{b_1}{4} |\lambda| \|w_1\|_2^2 + \frac{b_3}{4} \|\phi_1\|_2^2 + \frac{b_2 c_2}{8} \frac{\gamma^2}{\nu} |\eta'|^2 |\sigma|^2 \\
&\leq C \{ (1 + \frac{1}{|\eta'|^2}) | \langle f^0 \rangle |^2 + \|f_1^0\|_2^2 + \|\tilde{f}_1\|_2^2 \}.
\end{aligned}$$

Using the Poincaré inequality, we have $\frac{1}{16} D_{\eta'}(w_1) + \frac{b_3}{4} \|\phi_1\|_2^2 + \frac{b_2 c_2}{8} |\eta'|^2 |\sigma|^2 \geq \beta_1 E[u]$ for some constant $\beta_1 = \beta_1(r_0) > 0$. We thus obtain

$$(\operatorname{Re} \lambda + \beta_1) E[u] + \frac{1}{16} D_{\eta'}(w_1) \leq C \{ | \langle f^0 \rangle |^2 + \|f_1^0\|_2^2 + \|\tilde{f}_1\|_2^2 \}$$

for η' with $|\eta'| \geq r_0$. This completes the proof. \square

We have already shown in Proposition 4.2 that $-L_{\eta'}$ generates a contraction semigroup $e^{-tL_{\eta'}}$. Theorem 4.21 implies that $e^{-tL_{\eta'}}$ decays exponentially for $\eta' \in Q^*$ with $|\eta'| \geq r_0$.

Theorem 4.22. *There holds the estimate*

$$\|e^{-tL_{\eta'}} u_0\|_2 \leq C e^{-\frac{\beta_1}{2} t} \|u_0\|_2$$

uniformly for $\eta' \in Q^$ satisfying $|\eta'| \geq r_0$.*

Theorem 4.22 follows from Theorem 4.21 and [3, Chap. V, Theorem 1.11].

5 Proof of Theorems 3.1 and 3.2

In this section we give proofs of Theorem 3.1 and 3.2.

Proof of Theorem 3.1. As in the proof of Proposition 4.1, one can show that $\{\lambda; \operatorname{Re} \lambda > 0\} \subset \rho(-L)$ and if $\operatorname{Re} \lambda > 0$, then

$$\begin{aligned}
\|(\lambda + L)^{-1} f\|_{L^2(\Omega)} &\leq \frac{1}{\operatorname{Re} \lambda} \|f\|_{L^2(\Omega)}, \\
\|\nabla \tilde{\mathbb{P}} |(\lambda + L)^{-1} f\|_{L^2(\Omega)} &\leq \frac{1}{(\nu \operatorname{Re} \lambda)^{\frac{1}{2}}} \|f\|_{L^2(\Omega)}.
\end{aligned} \tag{5.1}$$

Therefore, $-L$ generates a contraction semigroup e^{-tL} on $L^2(\Omega)$. This completes the proof. \square

Proof of Theorem 3.2. We set

$$\Pi := U\chi_0\Pi_{\eta'}T, \quad \chi_0(\eta') = \begin{cases} 1 & |\eta'| \leq r_0, \\ 0 & |\eta'| \geq r_0. \end{cases}$$

It then follows from Proposition 2.2 that $\Pi^2 = \Pi$. Furthermore, by Theorem 4.7, we have

$$e^{-tL}\Pi u_0 = U\chi_0 e^{-tL_{\eta'}}\Pi_{\eta'}T u_0 = U\chi_0 e^{\lambda_{\eta'}t}\Pi_{\eta'}T u_0.$$

Since

$$\sup_{\eta' \in Q^*} \|T u_0\|_1 \leq C \|u_0\|_{L^1(\Omega)},$$

we see from Theorems 4.7 and 4.11 that

$$\begin{aligned} \|e^{-tL}\Pi u_0\|_{L^2(\Omega)}^2 &\leq C \int_{\eta' \in Q^*} \|\chi_0 e^{-tL_{\eta'}}\Pi_{\eta'}T u_0\|_2^2 d\eta' \leq C \int_{|\eta'| \leq r_0} e^{-\frac{\kappa_0}{2} \frac{\gamma^2}{\nu} |\eta'|^2 t} \|\Pi_{\eta'}T u_0\|_2^2 d\eta' \\ &\leq C \int_{|\eta'| \leq r_0} e^{-\frac{\kappa_0}{2} \frac{\gamma^2}{\nu} |\eta'|^2 t} \|T u_0\|_1^2 d\eta' \leq C t^{-\frac{n-1}{2}} \|u_0\|_{L^1(\Omega)}^2. \end{aligned}$$

On the other hand, we have

$$\|e^{-tL}\Pi u_0\|_{L^2(\Omega)}^2 \leq C \int_{|\eta'| \leq r_0} d\eta' \|u_0\|_{L^1(\Omega)}^2 \leq C \|u_0\|_{L^1(\Omega)}^2.$$

We thus obtain

$$\|e^{-tL}\Pi u_0\|_{L^2(\Omega)} \leq C(1+t)^{-\frac{n-1}{4}} \|u_0\|_{L^1(\Omega)}.$$

This proves (i) of Theorem 3.2.

As for the estimate for $e^{-tL}(I - \Pi)u_0$, we write it as

$$\begin{aligned} e^{-tL}(I - \Pi)u_0 &= U\chi_0 e^{-tL_{\eta'}}(I - \Pi_{\eta'})T + U(1 - \chi_0)e^{-tL_{\eta'}}T \\ &= U\chi_0(e^{-tL_{\eta'}} - e^{\lambda_{\eta'}t}\Pi_{\eta'})T + U(1 - \chi_0)e^{-tL_{\eta'}}T. \end{aligned}$$

It follows from Theorems 4.10 and 4.22 that

$$\|e^{-tL}(I - \Pi)u_0\|_{L^2(\Omega)} \leq C e^{-\beta t} \|u\|_{L^2(\Omega)},$$

where $\beta = \frac{1}{2} \min\{\beta_0, \beta_1\}$. This proves (ii) of Theorem 3.2.

Let us prove (iii) of Theorem 3.2. We write $e^{-tL}\Pi u_0$ as

$$e^{-tL}\Pi u_0 = U\chi_0 e^{\lambda_{\eta'}t}\Pi^{(0)}T u_0 + U\chi_0 e^{\lambda_{\eta'}t}(\Pi_{\eta'} - \Pi^{(0)})T u_0 =: J_1 + J_2.$$

For $\ell' = (\ell_1, \dots, \ell_{n-1}) \in \mathbb{Z}^{n-1}$, we denote by $\Omega_{per, \ell'}$ the set $\{(x' + \sum_{j=1}^{n-1} \frac{2\pi}{\alpha_j} e'_j, x_n, ; (x', x_n) \in \Omega_{per}\}$. By the definition of T , we have

$$\begin{aligned} \Pi^{(0)} T u_0 &= \left[\int_{\Omega_{per}} \langle (T\phi_0)(x', \cdot) \rangle dx \right] u^{(0)} \\ &= \left[\frac{1}{|\Omega_{per}| |Q^*|^{\frac{1}{2}}} \sum_{\ell' \in \mathbb{Z}} \int_{\Omega_{per, \ell'}} \phi_0(x) e^{-i\eta' \cdot x'} dx \right] u^{(0)} \\ &= \left[\frac{1}{|\Omega_{per}| |Q^*|^{\frac{1}{2}}} \int_{\Omega} \phi_0(x) e^{-i\eta' \cdot x'} dx \right] u^{(0)} = \frac{1}{(2\pi)^{\frac{n-1}{2}} |Q|^{\frac{1}{2}}} \hat{\sigma}_0(\eta') u^{(0)}, \end{aligned}$$

where

$$\sigma_0(x') = \frac{|Q|}{|\Omega_{per}|} \int_{\omega_1(x')}^{\omega_2(x')} \phi_0(x', x_n) dx_n.$$

It then follows that J_1 is written as

$$J_1 = \left[\frac{1}{(2\pi)^{n-1}} \int_{Q^*} \chi_0 e^{\lambda_{\eta'} t} \hat{\sigma}_0(\eta') e^{i\eta' \cdot x'} d\eta' \right] u^{(0)} = [e^{-tH} \sigma_0(x')] u^{(0)} + J_1^{(1)} + J_1^{(2)}.$$

Here

$$\begin{aligned} J_1^{(1)} &= \mathcal{F}^{-1} \left[(\chi_0 - 1) e^{-\frac{\gamma^2}{\nu} \kappa(\eta') t} \hat{\sigma}_0(\eta') \right] u^{(0)}, \\ J_1^{(2)} &= \left[\frac{1}{(2\pi)^{n-1}} \int_{Q^*} \chi_0 (e^{\lambda_{\eta'} t} - e^{-\frac{\gamma^2}{\nu} \kappa(\eta') t}) \hat{\sigma}_0(\eta') e^{i\eta' \cdot x'} d\eta' \right] u^{(0)}. \end{aligned}$$

By the Plancherel Theorem, $J_1^{(1)}$ is estimated as

$$\begin{aligned} \|J_1^{(1)}\|_{L^2(\Omega)}^2 &\leq \bar{d} \|\mathcal{F}^{-1}[(\chi_0 - 1) e^{-\frac{\gamma^2}{\nu} \kappa(\eta') t} \hat{\sigma}_0] u^{(0)}\|_{L^2(\mathbb{R}^{n-1})}^2 \\ &= (2\pi)^{-(n-1)} \bar{d} \|(\chi_0 - 1) e^{-\frac{\gamma^2}{\nu} \kappa(\eta') t} \hat{\sigma}_0\|_{L^2(\mathbb{R}^{n-1})}^2 \end{aligned}$$

with $\bar{d} > 0$ given by $\bar{d} = \max_{x' \in \mathbb{R}^{n-1}} \{\omega_2(x') - \omega_1(x')\}$. Since $\text{supp}(\chi_0 - 1) = \{|\eta'| \geq r_0\}$, we see that

$$\|(\chi_0 - 1) e^{-\frac{\gamma^2}{\nu} \kappa(\eta') t} \hat{\sigma}_0\|_{L^2(\mathbb{R}^{n-1})}^2 \leq C t^{-\frac{n-1}{2}} e^{-\frac{\gamma^2}{\nu} r_0^2 t} \|\phi_0\|_{L^1(\Omega)}^2,$$

and hence,

$$\|J_1^{(1)}\|_2 \leq C t^{-\frac{n-1}{4}} e^{-\frac{\gamma^2}{2\nu} r_0^2 t} \|\phi_0\|_{L^1(\Omega)}.$$

As for $J_1^{(2)}$, we have

$$e^{\lambda_{\eta'} t} - e^{-\kappa(\eta') t} = (\lambda_{\eta'} + \kappa(\eta')) t \int_0^1 e^{-\kappa(\eta') t + \theta(\lambda_{\eta'} + \kappa(\eta')) t} d\theta.$$

Since $\lambda_{\eta'} = -\frac{\gamma^2}{\nu}\kappa(\eta') + O(|\eta'|^3)$, we obtain

$$|e^{\lambda_{\eta'} t} - e^{-\frac{\gamma^2}{\nu}\kappa(\eta')t}| \leq C|\eta'|^3 t e^{-\frac{\kappa_0}{2}\frac{\gamma^2}{\nu}|\eta'|^2 t} \leq C|\eta'| e^{-\frac{\kappa_0}{4}\frac{\gamma^2}{\nu}|\eta'|^2 t},$$

and hence,

$$\|J_1^{(2)}\|_{L^2(\Omega)}^2 \leq C \int_{|\eta'| \leq r_0} |\eta'|^2 e^{-\frac{\kappa_0}{2}\frac{\gamma^2}{\nu}|\eta'|^2 t} d\eta' \left(\sup_{\eta' \in \mathbb{R}^{n-1}} |\widehat{\sigma}_0(\eta')| \right)^2 \leq C t^{-\frac{n-1}{2}-1} \|\phi_0\|_{L^1(\Omega)}^2.$$

Concerning J_2 , we see from Theorem 4.11 that

$$\begin{aligned} \|J_2\|_{L^2(\Omega)} &\leq C \|\chi_0 e^{\lambda_{\eta'} t} (\Pi_{\eta'} - \Pi^{(0)}) T u_0\|_{L^2(Q^*; L^2(\Omega_{per}))} \\ &\leq C \|\chi_0 |\eta'| e^{\lambda_{\eta'} t} T u_0\|_1 \|1\|_{L^2(Q^*)} \leq C(1+t)^{-\frac{n-1}{4}-\frac{1}{2}} \|u_0\|_{L^1(\Omega)}. \end{aligned}$$

We thus obtain the desired estimate. This completes the proof. \square

6 Outline of proof of Lemma 4.18

In this section we give an outline of the proof of Lemma 4.18. We here only give several lemmas necessary for the proof since Lemma 4.18 can be proved by an argument similar to the proof of [22, Chap. III, Theorem 1.4.1], where the proof for the Stokes system (i.e., $\eta' = 0$) is given.

We begin by

Lemma 6.1. *There holds the estimate*

$$\|u\|_2 \leq C \{ \|\nabla_{\eta'} u\|_{H_{per}^{-1}(\Omega_{per})} + \|u\|_{H_{per}^{-1}(\Omega_{per})} \}$$

for $u \in L^2(\Omega_{per})$.

Lemma 6.1 can be proved in a similar manner to that of [19, Chap. 3, Lemma 7.1]. (Cf., [22, Chap. II, Lemma 1.1.3].)

Lemma 6.2. *There hold the following inequalities for $u \in L_{*,per}^2(\Omega_{per})$:*

$$\|u\|_2 \leq C_1 \|\nabla_{\eta'} u\|_{H_{per}^{-1}(\Omega_{per})} \leq C_1 C_2 \|u\|_2.$$

Lemma 6.2 follows from Lemma 6.1 as in the proof of [22, Chap. II, Lemma 1.5.4].

Lemma 6.3. (i) *For every $g \in L_{*,per}^2(\Omega_{per})$, there exists $w \in H_{0,per}^1(\Omega_{per})$ satisfying*

$$\operatorname{div}_{\eta'} w = g, \quad \|\nabla_{\eta'} w\|_2 \leq C \|g\|_2.$$

(ii) *For every $f \in H_{per}^{-1}(\Omega_{per})$ satisfying*

$$[f, w] = 0, \quad \text{for all } w \in H_{0,per}^1(\Omega_{per}) \text{ with } \operatorname{div}_{\eta'} w = 0,$$

there uniquely exists $p \in L_{,per}^2(\Omega_{per})$ such that*

$$\nabla_{\eta'} p = f, \quad \|p\|_2 \leq C \|f\|_{H_{per}^{-1}(\Omega_{per})}.$$

One can prove Lemma 6.3 by using Lemma 6.2 as in the proof of [22, Chap. II, Lemma 2.1.1].

We define $H_{0,\sigma}^1(\Omega_{per})$ by

$$H_{0,\sigma}^1(\Omega_{per}) = \{w \in H_{0,per}^1(\Omega_{per}) ; \operatorname{div}_{\eta'} w = 0\}.$$

Lemma 6.4. *For every $f \in H_{0,per}^{-1}(\Omega_{per})$, there uniquely exists $w \in H_{0,\sigma}^1(\Omega_{per})$ satisfying*

$$(\nabla_{\eta'} w, \nabla_{\eta'} v) = [f, v]$$

for all $v \in H_{0,\sigma}^1(\Omega_{per})$, and

$$\|\nabla_{\eta'} w\|_2 \leq C \|f\|_{H_{per}^{-1}(\Omega_{per})}.$$

Furthermore, there uniquely exists $\phi \in L_{,per}^2(\Omega_{per})$ such that*

$$-\Delta_{\eta'} w - f = -\nabla_{\eta'} \phi$$

and

$$\|\phi\|_2 \leq C \|f\|_{H_{per}^{-1}(\Omega_{per})}.$$

Lemma 6.4 can be proved in a similar manner to the proof of [22, Chap. III, Theorem 1.3.1].

Lemma 4.18 follows from Lemmas 6.3 and 6.4 as in the proof of [22, Chap. III, Theorem 1.4.1]. \square

Acknowledgements. The work of Y. Kagei was partly supported by JSPS KAK-ENHI Grant Number 24340028, 22244009, 24224003.

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