

The Number of Orbits of Periodic Box-Ball Systems

Mikoda, Akihiro
Graduate School of Mathematics, Kyushu University

Inokuchi, Shuichi
Faculty of Mathematics, Kyushu University

Mizoguchi, Yoshihiro
Faculty of Mathematics, Kyushu University

Fujio, Mitsuhiro
Department of Systems Innovation and Informatics, Kyushu Institute of Technology

<https://hdl.handle.net/2324/11867>

出版情報 : Lecture notes in computer science. 4135, pp.181-194, 2006-08-22. Springer
バージョン :
権利関係 :



MHF Preprint Series

Kyushu University
21st Century COE Program
Development of Dynamic Mathematics with
High Functionality

The number of orbits of box-ball systems

A. Mikoda, S. Inokuchi
Y. Mizoguchi, M. Fujio

MHF 2006-24

(Received June 12, 2006)

Faculty of Mathematics
Kyushu University
Fukuoka, JAPAN

The number of orbits of periodic box-ball systems

Akihiro Mikoda¹, Shuichi Inokuchi², Yoshihiro Mizoguchi², and
Mitsuhiko Fujio³

¹ Graduate School of Mathematics, Kyushu University, JAPAN,

² Faculty of Mathematics, Kyushu University, JAPAN,
ym@math.kyushu-u.ac.jp,

³ Department of Systems Innovation and Informatics,
Kyushu Institute of Technology, JAPAN

Abstract. A box-ball system is a kind of cellular automata obtained by the ultradiscrete Lotka-Volterra equation. Similarities and differences between behaviours of discrete systems (cellular automata) and continuous systems (differential equations) are investigated using techniques of ultradiscretizations. Our motivation is to take advantage of behaviours of box-ball systems for new kinds of computations. Especially, we tried to find out useful periodic box-ball systems (pBBS) for random number generators. Applicable pBBS systems should have long fundamental cycles. We focus on pBBS with at most two kinds of solitons and investigate their behaviours, especially, the length of cycles and the number of orbits. We showed some relational equations of soliton sizes, a box size and the number of orbits. Varying a box size, we also found out some simulation results of the periodicity of orbits of pBBS with same kinds of solitons.

1 Introduction

In 1990, Takahashi and Satsuma introduced a soliton cellular automaton (SCA) [7]. The SCA is now called a box and ball system (BBS) because they explained transitions of the system using an infinite array of boxes and a finite number of balls. BBS has a property of solitons because its transition is obtained by the ultradiscrete Lotka-Volterra equation [6, 8].

In 1997, a new soliton cellular automaton is proposed by Takahashi et al [6]. That system is called box and ball system with a carrier (BBSC). BBSC can be considered as a kind of abstract model of Hyper-Threading (HT) Technology. HT Technology is a recent attractive CPU hardware technology. The main aim of HT Technology brings out the parallel efficiency of CPUs and improves the performance of a system. We hope that we could make a connection between a study of BBSC and the HT Technology in the future.

Recently, the research areas using ultradiscretizations are extending and it contains crystal formulations, combinatorics, stochastic cellular automata and algorithms [1, 2, 4, 5].

In 2003, the notion of periodic box-ball system (pBBS) is introduced by Yoshihara et al.[9]. They have shown a formula to determine the fundamental cycle of a pBBS for a given initial state. In the same year, Habu et al.[3] investigated properties about randomness and autocorrelations of configurations of pBBS and compared with Gold sequences. They showed some experimental results about their properties for a fixed system size varying the number of balls and the size of solitons.

In this paper, we focus on pBBS with at most two kinds of solitons. We re-formulate the pBBS and define sets of configurations precisely. A set of configurations with a same type is divided into some disjoint same size of orbits. We investigate the size of the configuration set and the number of orbits for designing a pBBS with a longer fundamental cycle. According to the result of Yoshihara et al.[9], we reformulate the equation of the fundamental cycles. Further, we induce the equation of the number of orbits and prove that its upperbound is not depended on the size of boxes. Finally, we show some experimental results between a size of boxes and the number of orbits.

2 periodic box-ball systems (pBBS)

Let $Q = \{0, 1\}$, N a natural number, $\bar{N} = \{1, 2, \dots, N\}$ and $\overline{2N} = \{1, 2, \dots, 2N\}$. We define three functions $dbl : Q^{\bar{N}} \rightarrow Q^{\overline{2N}}$, $snd : Q^{\overline{2N}} \rightarrow Q^{\bar{N}}$ and $trs : Q^{\overline{2N}} \rightarrow Q^{\overline{2N}}$ by $dbl(c)_j = c_{((j-1) \bmod N)+1}$, $snd(c)_j = c_{N+j}$ and $trs(c)_j = \min \left(1 - c_j, \sum_{i=1}^{j-1} (c_i - trs(c)_i) \right)$. The shift function $sft_\alpha : Q^{\bar{N}} \rightarrow Q^{\bar{N}}$ is defined by $sft_\alpha(c)_j = c_{((j-1+\alpha) \bmod N)+1}$ ($\alpha = 0, \dots, N-1$).

Definition 1 (N-pBBS). *The periodic box-ball system with the size N (N-pBBS) is the dynamical system (C, f) , where $C = \{c \in Q^{\bar{N}} \mid \sum_{j=1}^N c_j < \frac{N}{2}\}$ and the transition function $f : C \rightarrow C$ is defined by $f = snd \circ trs \circ dbl$.*

The definition of the N -pBBS is well-defined. It is guaranteed by the next proposition.

Proposition 1. *Assume $\#\{i \in \bar{N} \mid c_i = 1\} \leq \frac{N}{2}$ for $c \in Q^{\bar{N}}$.*

- (1) $\#\{i \in \bar{N} \mid c_i = 1\} = \#\{i \in \bar{N} \mid (snd \circ trs \circ dbl(c))_i = 1\}$, where $\#S$ is the size of the set S .
- (2) $(snd \circ trs \circ dbl) \circ sft_\alpha(c) = sft_\alpha \circ (snd \circ trs \circ dbl)(c)$ ($\alpha = 0, 1, \dots, N-1$).

The proposition is proved using the following lemma.

Lemma 1. *For $c \in Q^{\bar{N}}$, we put $\delta_j = \sum_{i=1}^j (dbl(c)_i - trs(dbl(c))_i)$,*

$\Delta_j = \sum_{i=1}^j \left(\text{dbl}(c)_i - \overline{\text{dbl}(c)}_i \right)$, where \overline{x} denote the complement $1 - x$ for $x \in Q$. Then we have

- (1) $\delta_j = \Delta_j + \max_{1 \leq i \leq j} \{ \text{dbl}(c)_i - \Delta_i \}$ ($j = 1, 2, \dots, 2N$).
- (2) $\Delta_{N+j} = \Delta_N + \Delta_j$ ($j = 1, 2, \dots, N$).
- (3) $\delta_{N+j} = \max\{\delta_N + \Delta_j, \delta_j\}$ ($j = 1, 2, \dots, 2N$).

The proof of Lemma 1 and Proposition 1 is listed in an appendix.

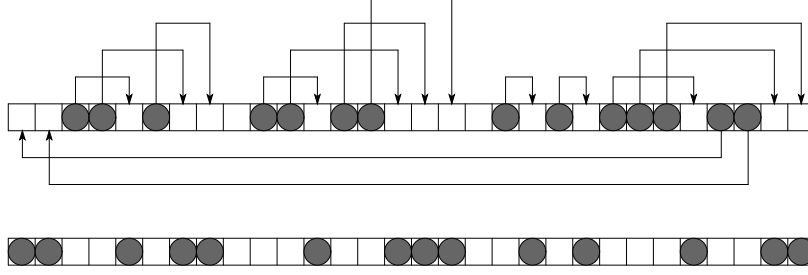


Fig. 1. Transition of pBBS

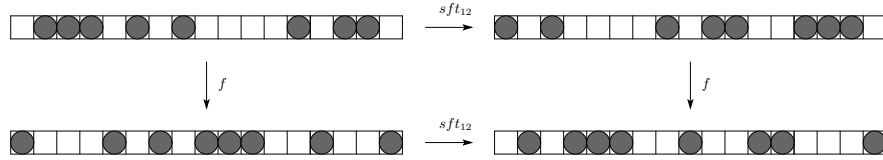


Fig. 2. commutative diagram

Example 1. Fig. 1 is an example of a transition of pBBS with size 30. Fig. 2 is an example transition ($f \circ \text{sft}_{12} = \text{sft}_{12} \circ f$) to confirm Proposition 1(2).

Definition 2 (Fundamental cycle of a pBBS). Let (C, f) be a pBBS with size N . The fundamental cycle of a configuration $c \in C$ is defined by $l(c) = \min \{t | f^t(c) = c, t > 0\}$.

Yoshihara et al. classified configurations of pBBS using size of solitons L_1, \dots, L_s and introduced an equation to compute the fundamental cycle of it.

Theorem 1 (Yoshihara 2003[9]). *Let (C, f) be a pBBS with size N . If a configuration $c \in C$ has a type (L_1, L_2, \dots, L_s) , then the fundamental cycle T of the configuration c is*

$$T = L.C.M \left(\frac{N_s N_{s-1}}{l_s l_0}, \frac{N_{s-1} N_{s-2}}{l_{s-1} l_0}, \dots, \frac{N_1 N_0}{l_1 l_0}, 1 \right),$$

where $l_j = L_j - L_{j+1}$ ($j = 1, 2, \dots, s-1$) and $N_j = l_0 + 2 \sum_{i=1}^j n_i (L_i - L_{j+1})$. \square

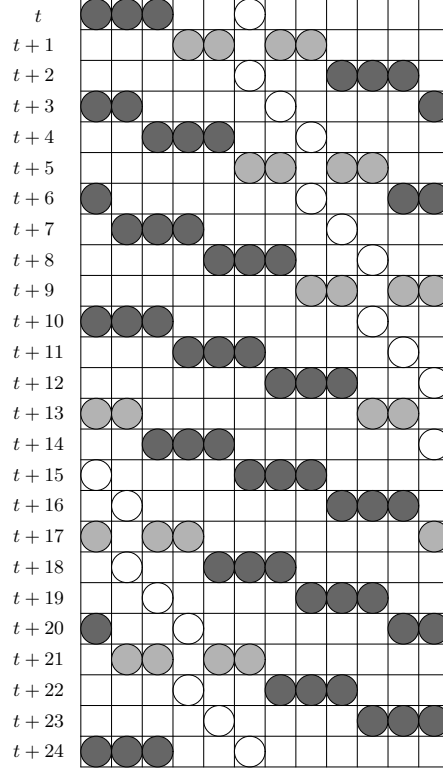


Fig. 3. Time evolution rule of pBBS

3 The number of orbits of a pBBS

In this section, we restrict the number of solitons up to 2. We re-formulate the class of configurations and imply a simple equation of the fundamental cycle.

We also introduce an equation of the total number of all configurations and the number of orbits.

Definition 3. Let (C, f) be pBBS with size N . All configurations with two solitons is defined by



$$C_2 = \{c \in C \mid c = 0^{x_1} 1^{l_1} 0^{x_2} 1^{l_2} 0^{x_3}, 0 \leq x_1, x_3, 1 \leq l_1, l_2, x_2, x_1 + l_1 + x_2 + l_2 + x_3 = N\}.$$

For numbers L_1 and L_2 ($L_1 + L_2 < \frac{N}{2}$, $L_1 \geq L_2$), we define a set $C_{(L_1, L_2, N)}$ of configurations with a type (L_1, L_2, N) as follows:

- (a) If $c = 0^{x_1} 1^{l_1} 0^{x_2} 1^{l_2} 0^{x_3}$ and $(l_1 \geq l_2, l_2 \leq x_2)$ then $c \in C_{(l_1, l_2, N)}$ and $sft_\alpha(c) \in C_{(l_1, l_2, N)}$ for $\alpha = 0, 1, \dots, N-1$.
- (b) If $c = 0^{x_1} 1^{l_1} 0^{x_2} 1^{l_2} 0^{x_3}$ and $(l_1 \geq l_2, x_2 < l_2)$ then $c \in C_{(l_1 + l_2 - x_2, x_2, N)}$, and $sft_\alpha(c) \in C_{(l_1 + l_2 - x_2, x_2, N)}$ for $\alpha = 0, 1, \dots, N-1$.

We note that we can find some number L_1 and L_2 for a configuration $c = 0^{x_1} 1^{l_1} 0^{x_2} 1^{l_2} 0^{x_3}$ ($l_1 < l_2$) to belong in $C_{(L_1, L_2, N)}$ using above Definition and sft_α .

Example 2. Let $N = 16$.

- (a) $c = 0^3 1^3 0^2 1^2 0^6 \in C_{(3, 2, N)}$

 $\dots L_1 = 3, L_2 = 2.$
- (b) $c = 0^5 1^2 0^1 1^2 0^6 \in C_{(3, 1, N)}$

 $\dots L_1 = 3, L_2 = 1.$

Definition 4 $((L_1, L_2, N)$ -pBBS). We define a subsystem (L_1, L_2, N) -pBBS of pBBS (C, f) with size N by a dynamical system $(C_{(L_1, L_2, N)}, f)$. The fundamental cycles for all $c \in C_{(L_1, L_2, N)}$ are the same number T . We call T as the fundamental cycle of $C_{(L_1, L_2, N)}$.

The definition of the (L_1, L_2, N) -pBBS is well-defined. It is guaranteed by the next proposition.

- Proposition 2.** (1) $f(c) \in C_{(L_1, L_2, N)}$ for any $c \in C_{(L_1, L_2, N)}$.
 (2) If $c_0, c_1 \in C_{(L_1, L_2, N)}$ then $l(c_0) = l(c_1)$.
 (3) Let $\alpha = L_1 + L_2, \beta = L_1 - L_2, N = 2(L_1 + L_2) + n$. The number of configurations of (L_1, L_2, N) -pBBS is $(2\alpha + n)(2\beta + n)$.

□

We denote the number $S = (2\alpha + n)(2\beta + n)$ in Proposition 2(3) by S .

Definition 5 (Orbits of pBBS). Configuration c and d are on the same orbit if and only if $d = f^i(c)$ for some i . (cf. Fig. 4)

$C_{(L_1, L_2, N)}$ is covered by several disjoint orbits like $\{f^i(c) \mid i \geq 0\}$. By Proposition 2(2), each orbit contains T elements, where T is the fundamental cycle of $C_{(L_1, L_2, N)}$.

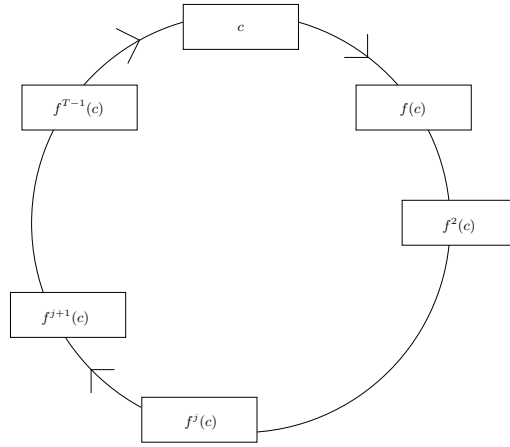


Fig. 4. The orbits of (L_1, L_2, N) -pBBS

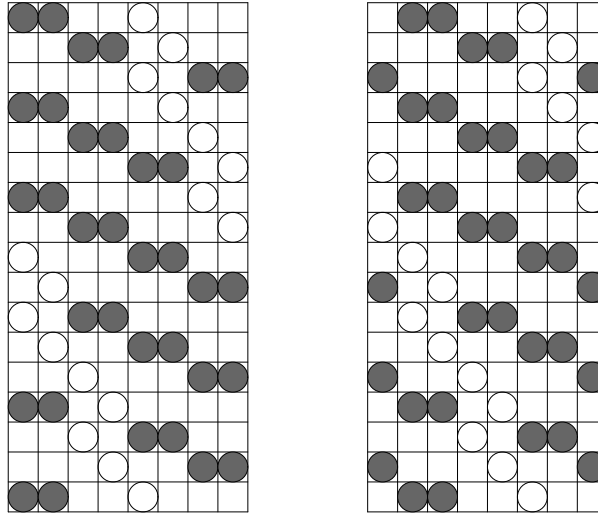


Fig. 5. The orbits of $(2, 1, 8)$ -pBBS

Example 3. Fig. 5 is an example of orbit sets. $C_{(2,1,8)}$ is covered by two orbit sets where each set contains 17 elements. The fundamental cycle of c is 17 for any $c \in C_{(2,1,8)}$. $S = 34$, $T = 17$ and $K = 2$.

Theorem 2 (The number of orbits). *Let $\alpha = L_1 + L_2, \beta = L_1 - L_2, N = 2(L_1 + L_2) + n$.*

(1) *The fundamental cycle T of $(L_1, L_2, 2(L_1 + L_2) + n)$ -pBBS is*

$$T = L.C.M \left(\frac{(2\alpha+n)(2\beta+n)}{G.C.D((2\alpha+n)(2\beta+n), \frac{\alpha-\beta}{2}n)}, \frac{2\beta+n}{G.C.D(2\beta+n, \beta)} \right),$$

(2) *The number of orbits K of $(L_1, L_2, 2(L_1 + L_2) + n)$ -pBBS is*

$$K = G.C.D \left((2\alpha + n)(2\beta + n), (2\alpha + n)\beta, \frac{\alpha - \beta}{2}n \right).$$

Proof. (1) is immediately induced by Theorem 1. Since $T = L.C.M \left(\frac{N_2 N_1}{l_2 l_0}, \frac{N_1 N_0}{l_1 l_0}, 1 \right)$, $N_2 = N$, $N_1 = N - 4L_2$, $N_0 = l_0$, $l_1 = L_1 - L_2$ and $l_2 = L_2$, we have $T = L.C.M \left(\frac{N(N-4L_2)}{L_2(N-2L_1-2L_2)}, \frac{N-4L_2}{L_1-L_2}, 1 \right)$. Since $\alpha = L_1 + L_2, \beta = L_1 - L_2$, we have $T = L.C.M \left(\frac{(2\alpha+n)(2\beta+n)}{G.C.D((2\alpha+n)(2\beta+n), \frac{\alpha-\beta}{2}n)}, \frac{2\beta+n}{G.C.D(2\beta+n, \beta)} \right)$.

(2) By Proposition 2(3) and above results, we have

$$\begin{aligned} K &= \frac{(2\alpha + n)(2\beta + n)}{L.C.M \left(\frac{(2\alpha+n)(2\beta+n)}{G.C.D((2\alpha+n)(2\beta+n), \frac{\alpha-\beta}{2}n)}, \frac{2\beta+n}{G.C.D(2\beta+n, \beta)} \right)} \\ &= \frac{2\alpha + n}{L.C.M \left(\frac{2\alpha+n}{G.C.D((2\alpha+n)(2\beta+n), \frac{\alpha-\beta}{2}n)}, \frac{1}{G.C.D(2\beta+n, \beta)} \right)} \\ &= \frac{(2\alpha + n)G.C.D(2\beta + n, \beta)}{L.C.M \left(\frac{(2\alpha+n)G.C.D(2\beta+n, \beta)}{G.C.D((2\alpha+n)(2\beta+n), \frac{\alpha-\beta}{2}n)}, 1 \right)} \\ &= \frac{(2\alpha + n)G.C.D(2\beta + n, \beta)}{\frac{(2\alpha+n)G.C.D(2\beta+n, \beta)}{G.C.D((2\alpha+n)(2\beta+n), (2\alpha+n)\beta, \frac{\alpha-\beta}{2}n)}} \\ &= G.C.D \left((2\alpha + n)(2\beta + n), (2\alpha + n)\beta, \frac{\alpha - \beta}{2}n \right) \end{aligned}$$

□

The next theorem shows some relations between the box-size n and the number of orbits K , especially the upper bound of the number of orbits K .

Theorem 3. *Let $\alpha = L_1 + L_2, \beta = L_1 - L_2$.*

- (1) $\gcd(L_1 - L_2, n) | K$,
- (2) $\gcd(2, n) | K$,
- (3) $\gcd(L_1 + L_2, n) | K$, and

$$(4) \quad K \mid \frac{\alpha\beta(\alpha-\beta)}{G.C.D(L_1, L_2)}.$$

Proof. (1) Let $L_1 - L_2 = ka, n = ma$. We have

$$\begin{aligned} K &= G.C.D \left((2\alpha + n)(2\beta + n), (2\alpha + n)\beta, \frac{\alpha - \beta}{2}n \right) \\ &= G.C.D \left((2\alpha + ma)(2ka + ma), (2\alpha + ma)ka, \frac{\alpha - ka}{2}ma \right) \\ &= a \times G.C.D \left((2\alpha + ma)(2k + m), (2\alpha + ma)k, \frac{\alpha - ka}{2}m \right). \end{aligned}$$

(2) Let $n = 2k$. We have

$$\begin{aligned} K &= G.C.D \left((2\alpha + n)(2\beta + n), (2\alpha + n)\beta, \frac{\alpha - \beta}{2}n \right) \\ &= G.C.D \left((2\alpha + 2k)(2\beta + 2k), (2\alpha + 2k)\beta, \frac{\alpha - \beta}{2}2k \right) \\ &= 2 \times G.C.D \left(2(\alpha + k)(\beta + k), (\alpha + k)\beta, \frac{\alpha - \beta}{2}k \right). \end{aligned}$$

(3) Let $L_1 + L_2 = ka, n = ma$. We have

$$\begin{aligned} K &= G.C.D \left((2\alpha + n)(2\beta + n), (2\alpha + n)\beta, \frac{\alpha - \beta}{2}n \right) \\ &= G.C.D \left((2ka + ma)(2\beta + ma), (2ka + ma)\beta, \frac{ka - \beta}{2}ma \right) \\ &= a \times G.C.D \left((2k + m)(2\beta + ma), (2k + m)\beta, \frac{ka - \beta}{2}m \right). \end{aligned}$$

(4) Let $g = G.C.D(an + b, cn)$. Since $g \mid cn$ and $cn = G.C.D(a, c) \times \frac{c}{G.C.D(a, c)} \times n$, we can set $g = g_a g_c g_n$ where $g_a \mid G.C.D(a, c)$, $g_c \mid \frac{c}{G.C.D(a, c)}$ and $g_n \mid n$.

Since $g_a g_n \mid an$ and $g_a g_n \mid (an + b)$, we have $g_a g_n \mid b$. So we can induce $g_a g_c g_n \mid \frac{bc}{G.C.D(a, c)}$.

Let $a = \beta$, $b = 2\alpha\beta$ and $c = \frac{\alpha - \beta}{2}$.

Then we have $g = G.C.D \left(2\alpha\beta + \beta n, \frac{\alpha - \beta}{2}n \right) \mid \frac{2\alpha\beta \cdot \frac{\alpha - \beta}{2}}{G.C.D(\beta, \frac{\alpha - \beta}{2})}$.

$$\begin{aligned}
K &= G.C.D \left((2\alpha + n)(2\beta + n), (2\alpha + n)\beta, \frac{\alpha - \beta}{2}n \right) \\
&= G.C.D \left((2\alpha + n)(2\beta + n), G.C.D \left(2\alpha\beta + \beta n, \frac{\alpha - \beta}{2}n \right) \right) \\
&\mid G.C.D \left((2\alpha + n)(2\beta + n), \frac{\alpha\beta(\alpha - \beta)}{G.C.D \left(\beta, \frac{\alpha - \beta}{2} \right)} \right) \\
&\mid \frac{\alpha\beta(\alpha - \beta)}{G.C.D(L_1 - L_2, L_2)} \\
&= \frac{\alpha\beta(\alpha - \beta)}{G.C.D(L_1, L_2)}.
\end{aligned}$$

□

4 Simulations

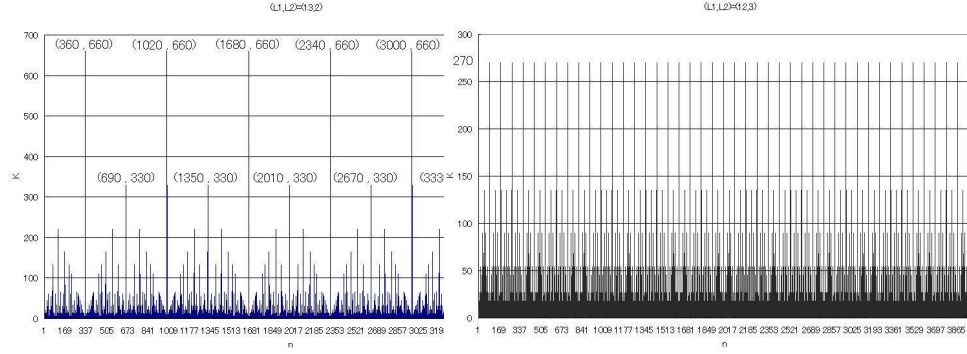


Fig. 6. Simulation results

The lefthand side of Fig. 6 is a graph of n and K for $C_{(13,2,2(13+2)+n)}$. A peak of K is 660 and $\frac{\alpha\beta(\alpha - \beta)}{G.C.D(L_1, L_2)} = \frac{15 \cdot 13 \cdot 2 \cdot 2}{G.C.D(13, 2)} = 660$. The righthand side of Fig. 6 is a graph of n and K for $C_{(12,3,2(12+3)+n)}$. A peak of K is 270 and $\frac{\alpha\beta(\alpha - \beta)}{G.C.D(L_1, L_2)} = \frac{15 \cdot 9 \cdot 2 \cdot 3}{G.C.D(12, 3)} = 270$.

In Theorem 3 we showed an upperbound of K . By the simulation results $\frac{\alpha\beta(\alpha - \beta)}{G.C.D(L_1, L_2)}$ may not only be an upper bound but also the maximum value of K .

Finally we have another conjecture from experimental results. pBBS with the number of orbits $K = 1$ must have a longer fundamental cycle, so the next conjecture may be useful to design a pBBS with a longer fundamental cycle.

Conjecture 1. Let K be the number of orbits for $C_{(L_1, L_2, 2(L_1+L_2)+n)}$. If $\gcd(L_1 - L_2, n) = 1$, $\gcd(2, n) = 1$ and $\gcd(L_1 + L_2, n) = 1$ then $K = 1$.

5 Concluding remarks

We re-formulate the pBBS with up to 2 kinds of solitons using precise equations. We showed the formula for the fundamental cycle and the number of orbits for pBBS. Further we proved the number of orbits is bounded some constant defined by the type of solitons. This means that we can design pBBS with longer fundamental cycle if we can choose larger box size pBBS. Future works contain to investigate an expression of orbits and behaviour of orbits when we increase sorts and number of solitons.

Acknowledgments

The author thanks the anonymous reviewers for their valuable comments. This work is partially supported by Grand-in-Aid for Scientific Research, Ministry of Education, Culture, Sports, Science and Technology, Japan.

References

1. K.Fukuda. *Box-ball systems and Robinson-Schensted-Knuth correspondence*. J.Algebraic combinatorics19(2004).
2. K.Fukuda, M.Okado and Y.Yamada. *Energy Functions in Box Ball Systems*. Int.J.Mod.Phys.A15(2000)1379-1392.
3. N.Habu, K.Imazu, M.Fujio, private communications, 2003.
4. N.Konno, T.Kunimatsu and X.Ma. *From stochastic partial difference equations to stochastic cellular automata through the ultra-discretization*. Applied Mathematics and Computation.155(2004)727-735.
5. A.Nagai, D.Takahashi and T.Tokihiro. *Soliton cellular automaton, Toda molecule equation and sorting algorithm*. Phys.Lett.A255(1999)265-271.
6. D.Takahashi and J.Matsukidaira. *Box and ball system with a carrier and ultradiscrete modified KdV equation*. J.Phys.A:Math.Gen.30(1997)L733-L739.
7. D.Takahashi and J.Satsuma. *A Soliton Cellular Automaton*. J.Phys.Soc.Jpn.59(1990)3514-3519.
8. T.Tokihiro, D.Takahashi, J.Matsukidaira and J.Satsuma. *From Soliton Equations to Integrable Cellular Automata through a Limiting Procedure*. Phys.Rev. Lett.76(1996)3247-3250.
9. D.Yoshihara, F.Yura and T.Tokihiro. *Fundamental cycle of a periodic box-ball system*. J.Phys.A:Math.Gen.36(2003)99-121.

Appendix

Proof. (of Lemma 1)

- (1) Since for any c , $trs(c)_1 = 0$, we have for $j = 1$, (left hand side) $= c_1 - trs(c)_1 = c_1 =$ (right hand side). Now suppose that the equality holds for some j . Then it follows that

$$\begin{aligned}
 \delta_{j+1} &= \delta_j + dbl(c)_{j+1} - trs(dbl(c))_{j+1} \\
 &= \delta_j + dbl(c)_{j+1} - \min \left\{ \overline{dbl(c)_{j+1}}, \delta_j \right\} \\
 &= \delta_j + \max \left\{ dbl(c)_{j+1} - \overline{dbl(c)_{j+1}}, dbl(c)_{j+1} - \delta_j \right\} \\
 &= \delta_j + \max \{ \Delta_{j+1} - \Delta_j, dbl(c)_{j+1} - \delta_j \} \\
 &= \max \{ \delta_j - \Delta_j + \Delta_{j+1}, dbl(c)_{j+1} \} \\
 &= \max \left\{ \max_{1 \leq i \leq j} \{ dbl(c)_i - \Delta_i \} + \Delta_{j+1}, dbl(c)_{j+1} \right\} \\
 &= \Delta_{j+1} + \max \left\{ \max_{1 \leq i \leq j} \{ dbl(c)_i - \Delta_i \}, dbl(c)_{j+1} - \Delta_{j+1} \right\} \\
 &= \Delta_{j+1} + \max_{1 \leq i \leq j+1} \{ dbl(c)_i - \Delta_i \},
 \end{aligned}$$

which establishes the equality for $j + 1$.

- (2) It follows from the fact that $dbl(c)_{N+i} = dbl(c)_i$ ($i = 1, 2, \dots, N$).
 (3) By virtue of (1) and (2),

$$\begin{aligned}
 \delta_{N+j} &= \Delta_{N+j} + \max_{1 \leq i \leq N+j} \{ dbl(c)_i - \Delta_i \} \\
 &= \Delta_{N+j} + \max \left\{ \max_{1 \leq i \leq N} \{ dbl(c)_i - \Delta_i \}, \max_{1 \leq i \leq j} \{ dbl(c)_{N+i} - \Delta_{N+i} \} \right\} \\
 &= \Delta_N + \Delta_j \\
 &\quad + \max \left\{ \max_{1 \leq i \leq N} \{ dbl(c)_i - \Delta_i \}, \max_{1 \leq i \leq j} \{ dbl(c)_i - \Delta_N - \Delta_i \} \right\} \\
 &= \max \left\{ \Delta_j + \Delta_N + \max_{1 \leq i \leq N} \{ dbl(c)_i - \Delta_i \}, \right. \\
 &\quad \left. \Delta_j + \max_{1 \leq i \leq j} \{ dbl(c)_i - \Delta_i \} \right\} \\
 &= \max \{ \Delta_j + \delta_N, \delta_j \}.
 \end{aligned}$$

□

Proof. (of Proposition 1).

- (1) By (3) of Lemma 1, $\delta_{2N} = \max \{ \delta_N + \Delta_N, \delta_N \}$. On the other hand, by the assumption, $\Delta_N = \sum_{i=1}^N \left(dbl(c)_i - \overline{dbl(c)_i} \right) = \left(\sum_{i=1}^N 2dbl(c)_i \right) - N = 2 \left(\sum_{i=1}^N c_i \right) -$

$N \leq 0$. Hence we have $\delta_{2N} = \delta_N$. This implies that

$$\begin{aligned}
& \#\{i \in \overline{N} \mid c_i = 1\} - \#\{i \in \overline{N} \mid (snd \circ trs \circ dbl(c))_i = 1\} \\
&= \left(\sum_{i=1}^N c_i \right) - \left(\sum_{i=1}^N (snd \circ trs \circ dbl(c))_i \right) \\
&= \left(\sum_{i=N+1}^{2N} d_i \right) - \left(\sum_{i=N+1}^{2N} trs(d)_i \right) \\
&= \sum_{i=N+1}^{2N} (d_i - trs(d)_i) \\
&= s_{2N} - s_N \\
&= 0.
\end{aligned}$$

(2) Since $sft_\alpha = \underbrace{sft_1 \circ \cdots \circ sft_1}_\alpha$, it suffices to show this for $\alpha = 1$. For the sake of simplicity, we put $\bar{d} = dbl(c)$ and $e = dbl(sft_1(c))$. Then the equations are rewritten as

$$sft_1(snd(trs(d)))_j = snd(trs(e))_j \quad (j = 1, 2, \dots, N) \quad (1)$$

Furthermore, to describe the effect of shift, we put $\delta_j = \sum_{i=1}^j (d_i - trs(d)_i)$, $\varepsilon_j = \sum_{i=1}^j (e_i - trs(e)_i)$, $\Delta_j = \sum_{i=1}^j (d_i - \bar{d}_i)$, $E_j = \sum_{i=1}^j (e_i - \bar{e}_i)$.

These variables are related as $\Delta_{j+1} = E_j + (c_1 - \bar{c}_1)$, $\delta_{j+1} = \max\{\varepsilon_j, E_j + c_1\}$ for $j = 1, 2, \dots, 2N - 1$. In fact, $\Delta_{j+1} = \sum_{i=1}^{j+1} (d_i - \bar{d}_i) = (d_1 - \bar{d}_1) + \sum_{i=1}^j (e_i - \bar{e}_i) = E_j + (c_1 - \bar{c}_1)$.

For the second one, by Lemma 1 (1),

$$\begin{aligned}
\delta_{j+1} &= \Delta_{j+1} + \max_{1 \leq i \leq j+1} \{d_i - \Delta_i\} \\
&= \Delta_{j+1} + \max \left\{ d_1 - \Delta_1, \max_{2 \leq i \leq j+1} \{d_i - \Delta_i\} \right\} \\
&= E_j + (c_1 - \bar{c}_1) + \max \left\{ \bar{c}_1, \max_{1 \leq i \leq j} \{d_{i+1} - \Delta_{i+1}\} \right\} \\
&= E_j + \max \left\{ c_1, \max_{1 \leq i \leq j} \{d_{i+1} - \Delta_{i+1} + (c_1 - \bar{c}_1)\} \right\} \\
&= E_j + \max \left\{ c_1, \max_{1 \leq i \leq j} \{e_i - E_i\} \right\} \\
&= \max \left\{ E_j + c_1, E_j + \max_{1 \leq i \leq j} \{e_i - E_i\} \right\} \\
&= \max \{E_j + c_1, \varepsilon_j\}.
\end{aligned}$$

Next, we claim that $sft_1(snd(trs(d)))_j$ and $snd(trs(e))_j$ are related by

$$sft_1(snd(trs(d)))_j = \max \{snd(trs(e))_j, \min\{\bar{e}_j, E_{N+j-1} + c_1\}\} \quad (2)$$

for $j = 1, 2, \dots, N$. In fact, if $j < N$,

$$\begin{aligned}
sft_1(snd(trs(d)))_j &= snd(trs(d))_{j+1} \\
&= trs(d)_{N+j+1} \\
&= \min \{\bar{d}_{N+j+1}, \delta_{N+j}\} \\
&= \min \{\bar{e}_{N+j}, \max\{\varepsilon_{N+j-1}, E_{N+j-1} + c_1\}\} \\
&= \max \{\min\{\bar{e}_{N+j}, \varepsilon_{N+j-1}\}, \min\{\bar{e}_{N+j}, E_{N+j-1} + c_1\}\} \\
&= \max \{trs(e)_{N+j}, \min\{\bar{e}_{N+j}, E_{N+j-1} + c_1\}\} \\
&= \max \{snd(trs(e))_j, \min\{\bar{e}_j, E_{N+j-1} + c_1\}\}.
\end{aligned}$$

For $j = N$,

$$\begin{aligned}
sft_1(snd(trs(d)))_N &= snd(trs(d))_1 \\
&= trs(d)_{N+1} \\
&= \min \{\bar{d}_{N+1}, \delta_N\} \\
&= \min \{\bar{e}_N, \max\{\varepsilon_{N-1}, E_{N-1} + c_1\}\} \\
&= \max \{\min\{\bar{e}_N, \varepsilon_{N-1}\}, \min\{\bar{e}_N, E_{N-1} + c_1\}\} \\
&= \max \{trs(e)_N, \min\{\bar{e}_N, E_{N-1} + c_1\}\} \\
&= \max \{snd(trs(e))_N, \min\{\bar{e}_N, E_{N-1} + c_1\}\}.
\end{aligned}$$

Now all we have to show is that

$$snd(trs(e))_j \geq \min\{\bar{e}_j, E_{N+j-1} + c_1\} \quad (j = 1, 2, \dots, N). \quad (3)$$

In fact, by combining this with the relation (2), we obtain (1).
 To show (3), we apply similar argument about δ_j 's and Δ_j 's to ε_j 's and E_j 's. Recall that, from the assumption of c , it follows that $E_N \leq 0$. By Lemma 1 (3), we have $\varepsilon_N = \max\{\varepsilon_N + E_N, \varepsilon_N\} = \varepsilon_{2N}$. On the other hand, by Lemma 1 (1),

$$\varepsilon_{2N} = E_{2N} + \max_{1 \leq i \leq 2N} \{e_i - E_i\} \geq E_{2N} + e_{2N} - E_{2N} = e_{2N} = c_1.$$

Thus we have $\varepsilon_N \geq c_1$. From this it follows that

$$\begin{aligned} \varepsilon_{N+j-1} &= \max\{\varepsilon_{j-1}, \varepsilon_N + E_{j-1}\} \\ &\geq \varepsilon_N + E_{j-1} \\ &\geq c_1 + E_{j-1} \\ &\geq c_1 + E_{j-1} + E_N \\ &= E_{N+j-1} + c_1. \end{aligned}$$

Consequently, we have

$$\begin{aligned} \text{snd}(\text{trs}(e))_j &= \text{trs}(e)_{N+j} \\ &= \min\{\bar{e}_{N+j}, \varepsilon_{N+j-1}\} \\ &\geq \min\{\bar{e}_{N+j}, E_{N+j-1} + c_1\} \\ &= \min\{\bar{e}_j, E_{N+j-1} + c_1\}, \end{aligned}$$

that is, the inequality (3). □

List of MHF Preprint Series, Kyushu University

21st Century COE Program

Development of Dynamic Mathematics with High Functionality

- MHF2005-1 Hideki KOSAKI
Matrix trace inequalities related to uncertainty principle
- MHF2005-2 Masahisa TABATA
Discrepancy between theory and real computation on the stability of some finite element schemes
- MHF2005-3 Yuko ARAKI & Sadanori KONISHI
Functional regression modeling via regularized basis expansions and model selection
- MHF2005-4 Yuko ARAKI & Sadanori KONISHI
Functional discriminant analysis via regularized basis expansions
- MHF2005-5 Kenji KAJIWARA, Tetsu MASUDA, Masatoshi NOUMI, Yasuhiro OHTA & Yasuhiko YAMADA
Point configurations, Cremona transformations and the elliptic difference Painlevé equations
- MHF2005-6 Kenji KAJIWARA, Tetsu MASUDA, Masatoshi NOUMI, Yasuhiro OHTA & Yasuhiko YAMADA
Construction of hypergeometric solutions to the q -Painlevé equations
- MHF2005-7 Hiroki MASUDA
Simple estimators for non-linear Markovian trend from sampled data:
I. ergodic cases
- MHF2005-8 Hiroki MASUDA & Nakahiro YOSHIDA
Edgeworth expansion for a class of Ornstein-Uhlenbeck-based models
- MHF2005-9 Masayuki UCHIDA
Approximate martingale estimating functions under small perturbations of dynamical systems
- MHF2005-10 Ryo MATSUZAKI & Masayuki UCHIDA
One-step estimators for diffusion processes with small dispersion parameters from discrete observations
- MHF2005-11 Junichi MATSUKUBO, Ryo MATSUZAKI & Masayuki UCHIDA
Estimation for a discretely observed small diffusion process with a linear drift
- MHF2005-12 Masayuki UCHIDA & Nakahiro YOSHIDA
AIC for ergodic diffusion processes from discrete observations

- MHF2005-13 Hiromichi GOTO & Kenji KAJIWARA
Generating function related to the Okamoto polynomials for the Painlevé IV equation
- MHF2005-14 Masato KIMURA & Shin-ichi NAGATA
Precise asymptotic behaviour of the first eigenvalue of Sturm-Liouville problems with large drift
- MHF2005-15 Daisuke TAGAMI & Masahisa TABATA
Numerical computations of a melting glass convection in the furnace
- MHF2005-16 Raimundas VIDŪNAS
Normalized Leonard pairs and Askey-Wilson relations
- MHF2005-17 Raimundas VIDŪNAS
Askey-Wilson relations and Leonard pairs
- MHF2005-18 Kenji KAJIWARA & Atsushi MUKAIHIRA
Soliton solutions for the non-autonomous discrete-time Toda lattice equation
- MHF2005-19 Yuu HARIYA
Construction of Gibbs measures for 1-dimensional continuum fields
- MHF2005-20 Yuu HARIYA
Integration by parts formulae for the Wiener measure restricted to subsets in \mathbb{R}^d
- MHF2005-21 Yuu HARIYA
A time-change approach to Kotani's extension of Yor's formula
- MHF2005-22 Tadahisa FUNAKI, Yuu HARIYA & Mark YOR
Wiener integrals for centered powers of Bessel processes, I
- MHF2005-23 Masahisa TABATA & Satoshi KAIZU
Finite element schemes for two-fluids flow problems
- MHF2005-24 Ken-ichi MARUNO & Yasuhiro OHTA
Determinant form of dark soliton solutions of the discrete nonlinear Schrödinger equation
- MHF2005-25 Alexander V. KITAEV & Raimundas VIDŪNAS
Quadratic transformations of the sixth Painlevé equation
- MHF2005-26 Toru FUJII & Sadanori KONISHI
Nonlinear regression modeling via regularized wavelets and smoothing parameter selection
- MHF2005-27 Shuichi INOKUCHI, Kazumasa HONDA, Hyen Yeal LEE, Tatsuro SATO, Yoshihiro MIZOGUCHI & Yasuo KAWAHARA
On reversible cellular automata with finite cell array

- MHF2005-28 Toru KOMATSU
Cyclic cubic field with explicit Artin symbols
- MHF2005-29 Mitsuhiro T. NAKAO, Kouji HASHIMOTO & Kaori NAGATOU
A computational approach to constructive a priori and a posteriori error estimates for finite element approximations of bi-harmonic problems
- MHF2005-30 Kaori NAGATOU, Kouji HASHIMOTO & Mitsuhiro T. NAKAO
Numerical verification of stationary solutions for Navier-Stokes problems
- MHF2005-31 Hidefumi KAWASAKI
A duality theorem for a three-phase partition problem
- MHF2005-32 Hidefumi KAWASAKI
A duality theorem based on triangles separating three convex sets
- MHF2005-33 Takeaki FUCHIKAMI & Hidefumi KAWASAKI
An explicit formula of the Shapley value for a cooperative game induced from the conjugate point
- MHF2005-34 Hideki MURAKAWA
A regularization of a reaction-diffusion system approximation to the two-phase Stefan problem
- MHF2006-1 Masahisa TABATA
Numerical simulation of Rayleigh-Taylor problems by an energy-stable finite element scheme
- MHF2006-2 Ken-ichi MARUNO & G R W QUISPEL
Construction of integrals of higher-order mappings
- MHF2006-3 Setsuo TANIGUCHI
On the Jacobi field approach to stochastic oscillatory integrals with quadratic phase function
- MHF2006-4 Kouji HASHIMOTO, Kaori NAGATOU & Mitsuhiro T. NAKAO
A computational approach to constructive a priori error estimate for finite element approximations of bi-harmonic problems in nonconvex polygonal domains
- MHF2006-5 Hidefumi KAWASAKI
A duality theory based on triangular cylinders separating three convex sets in R^n
- MHF2006-6 Raimundas VIDŪNAS
Uniform convergence of hypergeometric series
- MHF2006-7 Yuji KODAMA & Ken-ichi MARUNO
N-Soliton solutions to the DKP equation and Weyl group actions

- MHF2006-8 Toru KOMATSU
Potentially generic polynomial
- MHF2006-9 Toru KOMATSU
Generic sextic polynomial related to the subfield problem of a cubic polynomial
- MHF2006-10 Shu TEZUKA & Anargyros PAPAGEORGIOU
Exact cubature for a class of functions of maximum effective dimension
- MHF2006-11 Shu TEZUKA
On high-discrepancy sequences
- MHF2006-12 Raimundas VIDŪNAS
Detecting persistent regimes in the North Atlantic Oscillation time series
- MHF2006-13 Toru KOMATSU
Tamely Eisenstein field with prime power discriminant
- MHF2006-14 Nalini JOSHI, Kenji KAJIWARA & Marta MAZZOCCO
Generating function associated with the Hankel determinant formula for the solutions of the Painlevé IV equation
- MHF2006-15 Raimundas VIDŪNAS
Darboux evaluations of algebraic Gauss hypergeometric functions
- MHF2006-16 Masato KIMURA & Isao WAKANO
New mathematical approach to the energy release rate in crack extension
- MHF2006-17 Toru KOMATSU
Arithmetic of the splitting field of Alexander polynomial
- MHF2006-18 Hiroki MASUDA
Likelihood estimation of stable Lévy processes from discrete data
- MHF2006-19 Hiroshi KAWABI & Michael RÖCKNER
Essential self-adjointness of Dirichlet operators on a path space with Gibbs measures via an SPDE approach
- MHF2006-20 Masahisa TABATA
Energy stable finite element schemes and their applications to two-fluid flow problems
- MHF2006-21 Yuzuru INAHAMA & Hiroshi KAWABI
Asymptotic expansions for the Laplace approximations for Itô functionals of Brownian rough paths
- MHF2006-22 Yoshiyuki KAGEI
Resolvent estimates for the linearized compressible Navier-Stokes equation in an infinite layer

MHF2006-23 Yoshiyuki KAGEI

Asymptotic behavior of the semigroup associated with the linearized
compressible Navier-Stokes equation in an infinite layer

MHF2006-24 Akihiro MIKODA, Shuichi INOKUCHI, Yoshihiro MIZOGUCHI & Mitsuhiko
FUJIO

The number of orbits of box-ball systems